

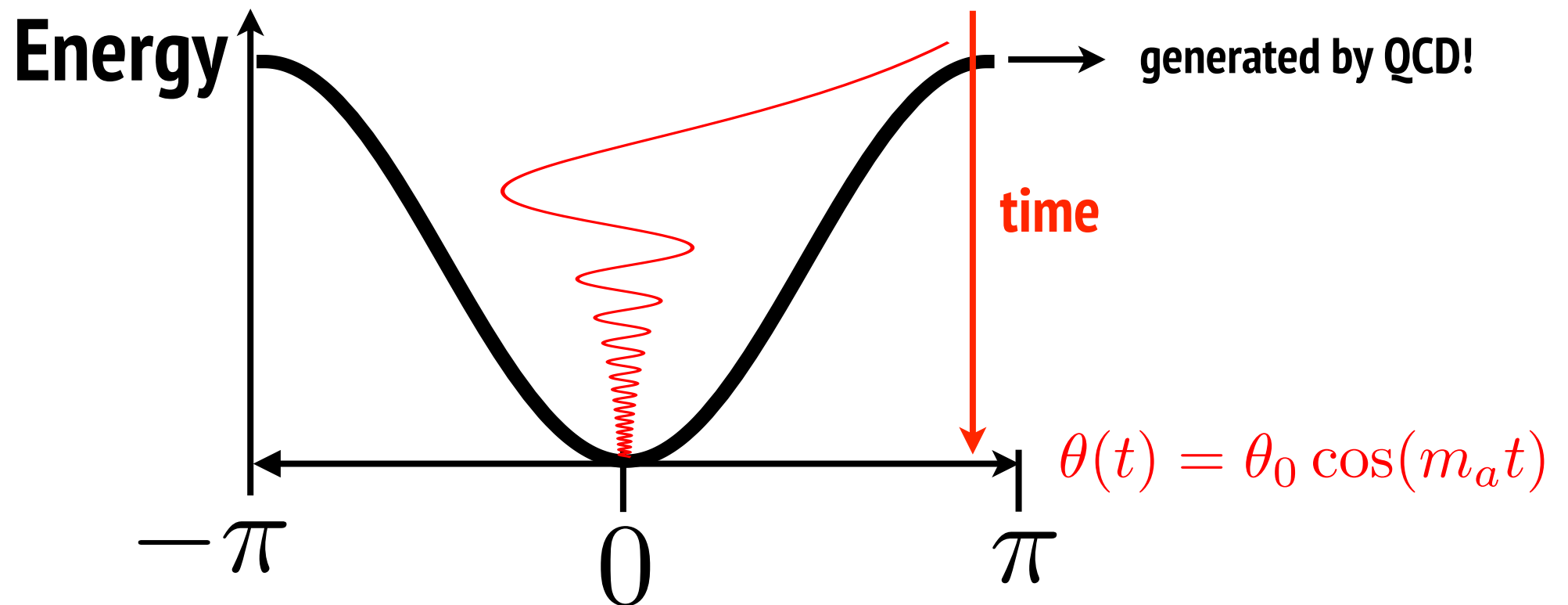
MAD MAX

A black and white photograph of a woman in a patterned dress taking a mirror selfie in a jewelry store. The image is used as a background for the text. The woman is holding a camera up to her face, and her reflection is visible in the mirror. The store has many mirrors and jewelry displays.

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Universidad de Zaragoza (Spain)
Max Planck Institute für Physik

Axions are necessarily dark matter

- is it a dynamical field? $\theta(t, \mathbf{x})$

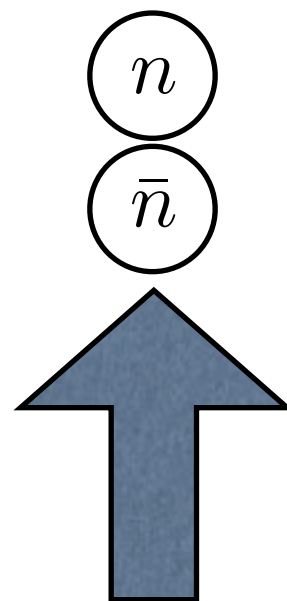


~ One parameter theory

$$\theta(t, x) = a(t, x)/f_a$$

axion mass

$$m_a = 6 \text{ meV} \frac{10^9 \text{ GeV}}{f_a}$$



Coherent oscillations

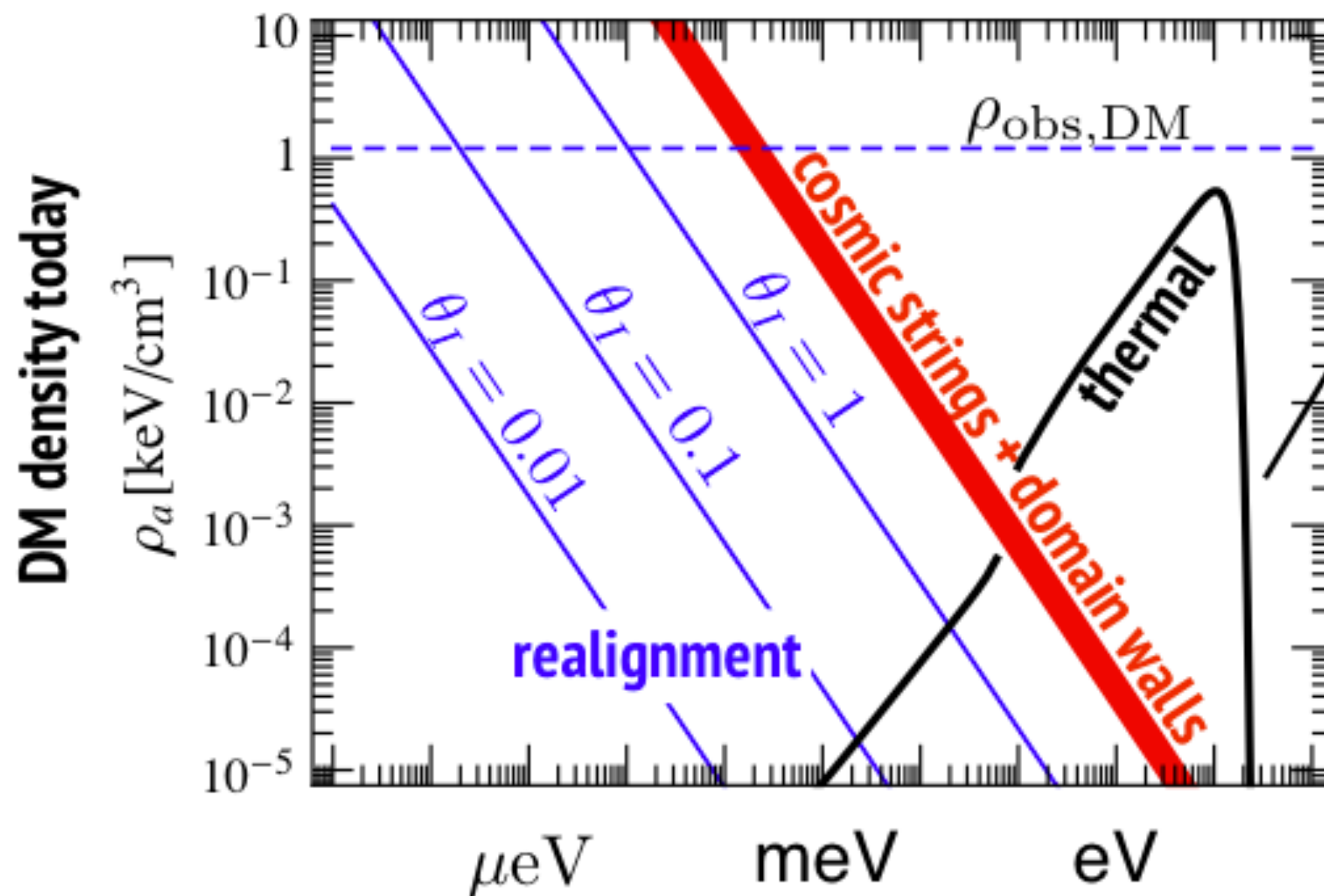
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Dark Matter Axions

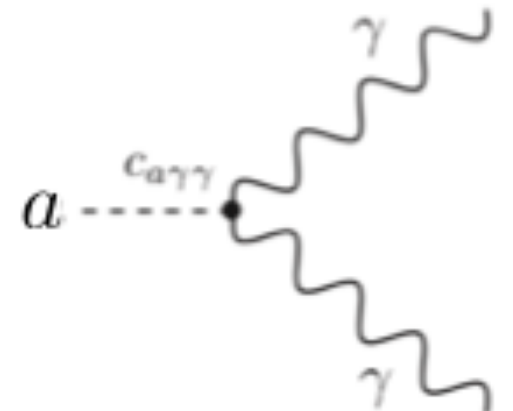
Measured today $|\theta| < 10^{-10}$ (strong CP problem)

Relic density as function of mass

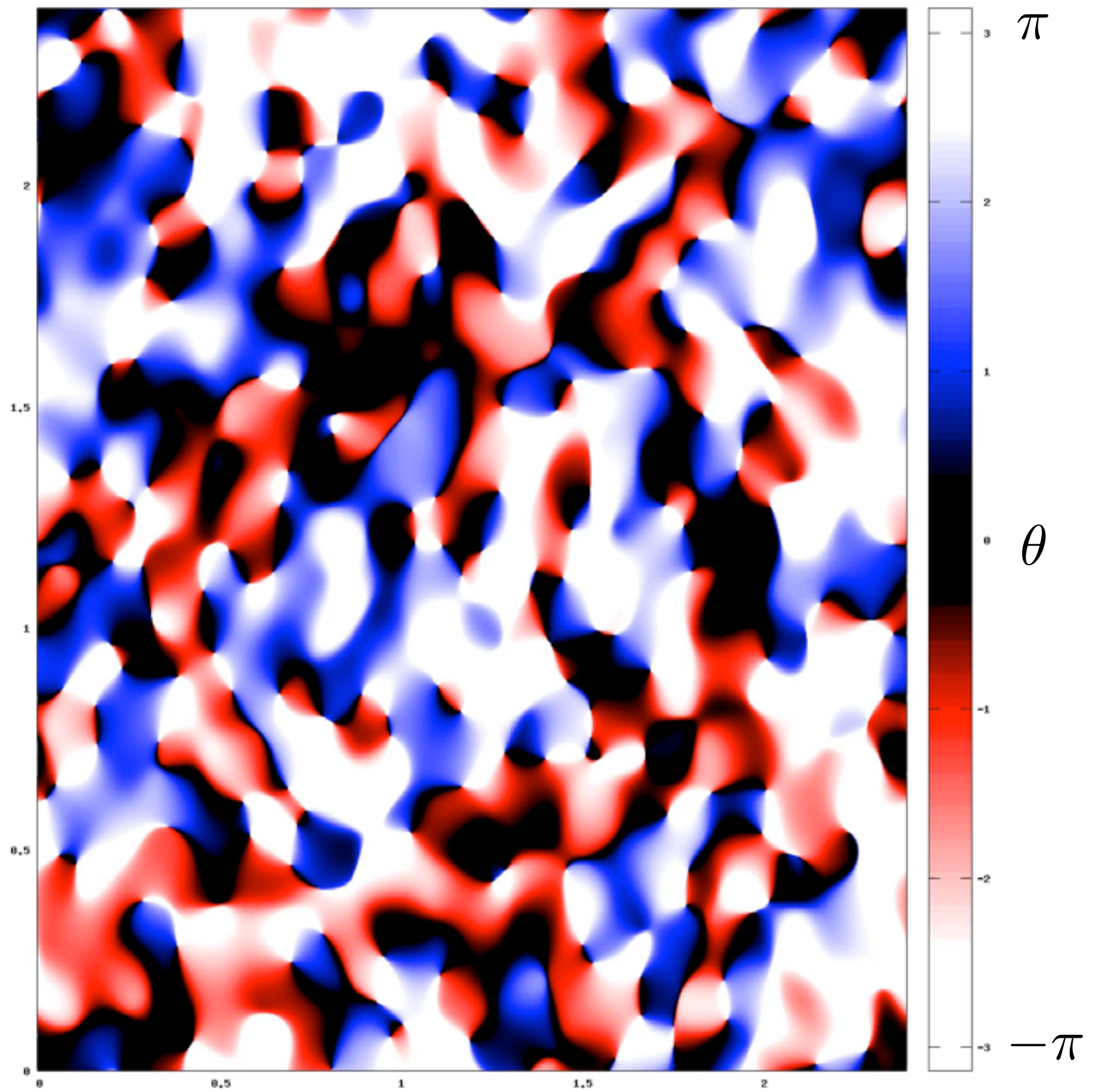
- Random IC (after a phase transition)
- One Universal IC (inflation smoothens the axion field)



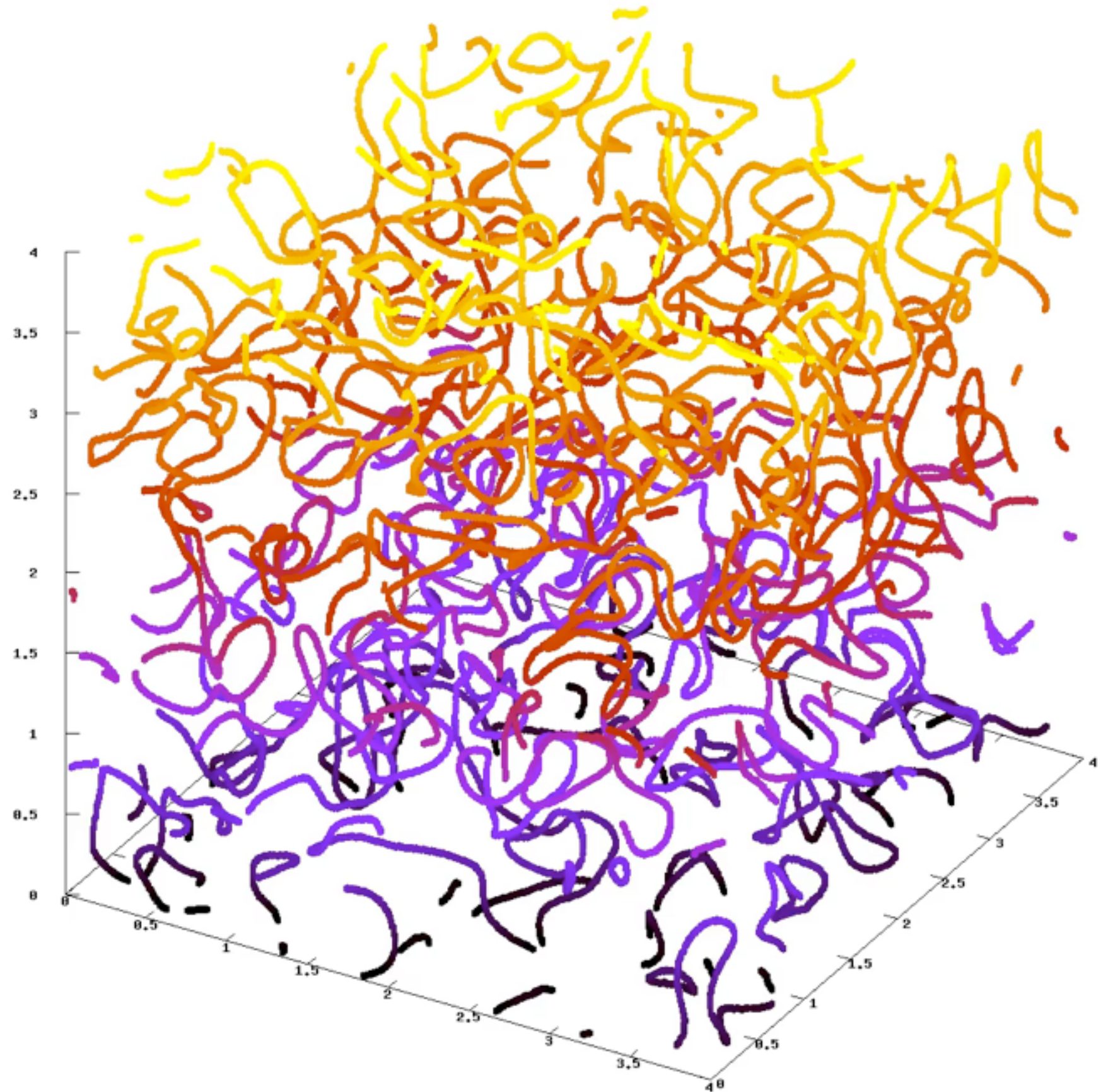
Axion decay



Theta evolution, Averaged SCENARIO I

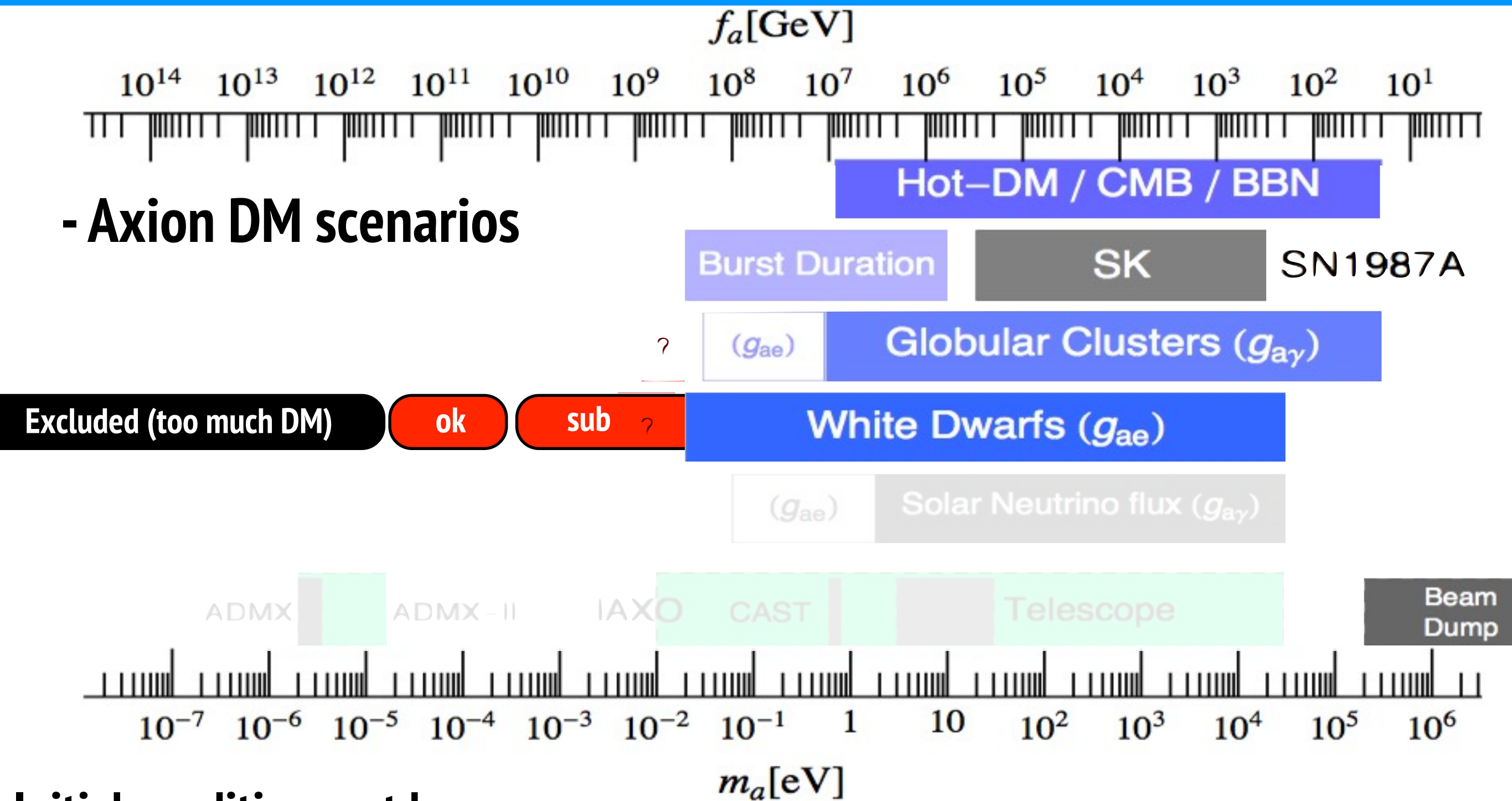


DM dominated by string radiation



Axion dark matter

- Axion DM scenarios

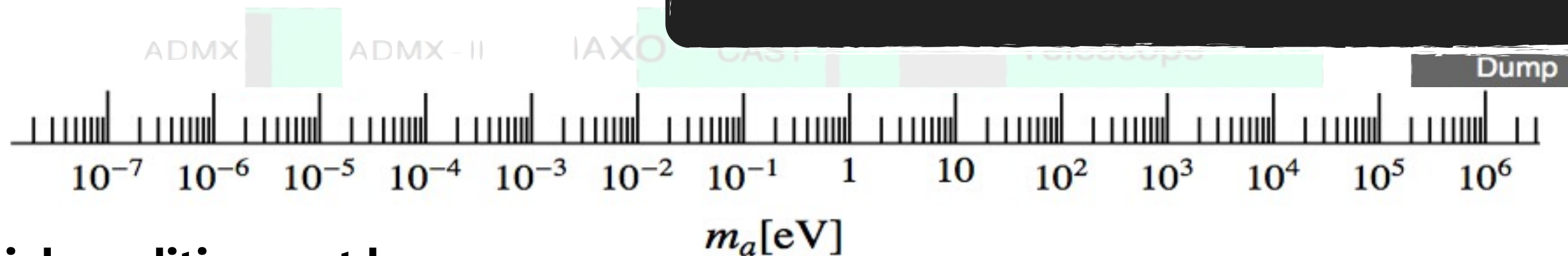
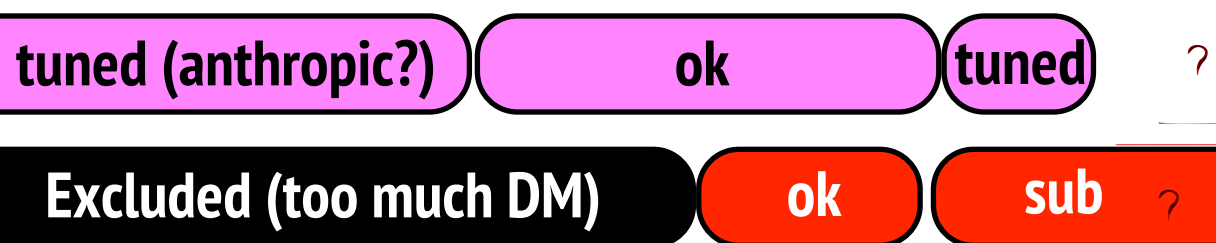


Initial conditions set by :

**Phase transition (N=1)
strings+unstable DW's**

Axion dark matter

- Axion DM scenarios



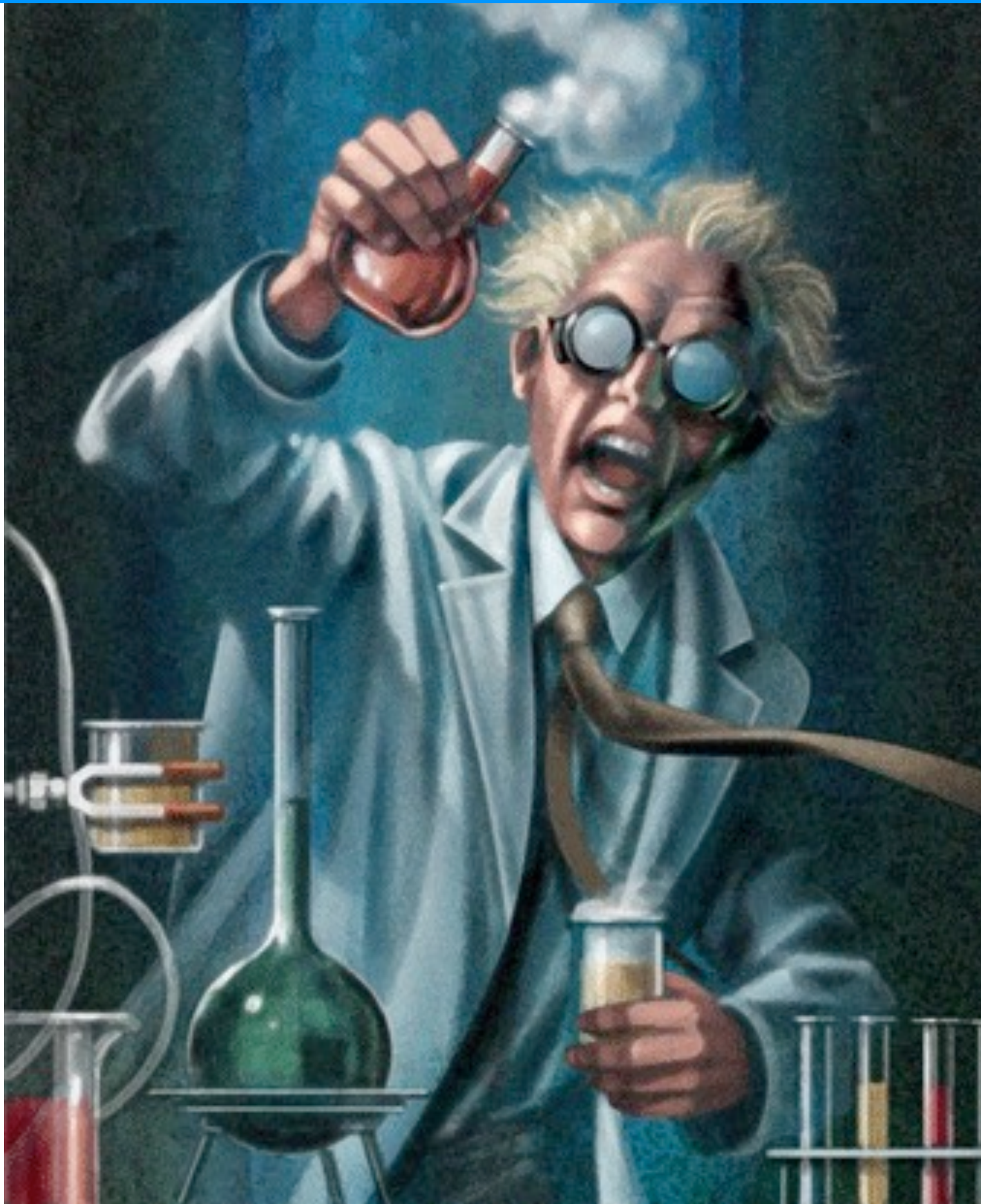
Initial conditions set by :

Inflation smooth

$$\Omega_{\text{aDM}} h^2 \simeq \theta_I^2 \left(\frac{80 \mu\text{eV}}{m_a} \right)^{1.19}$$

**Phase transition (N=1)
strings+unstable DW's**

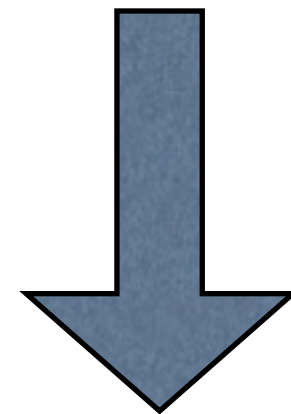
Detecting Axions



$$\rho_{\text{aDM}} = 0.3 \frac{\text{GeV}}{\text{cm}^3}$$

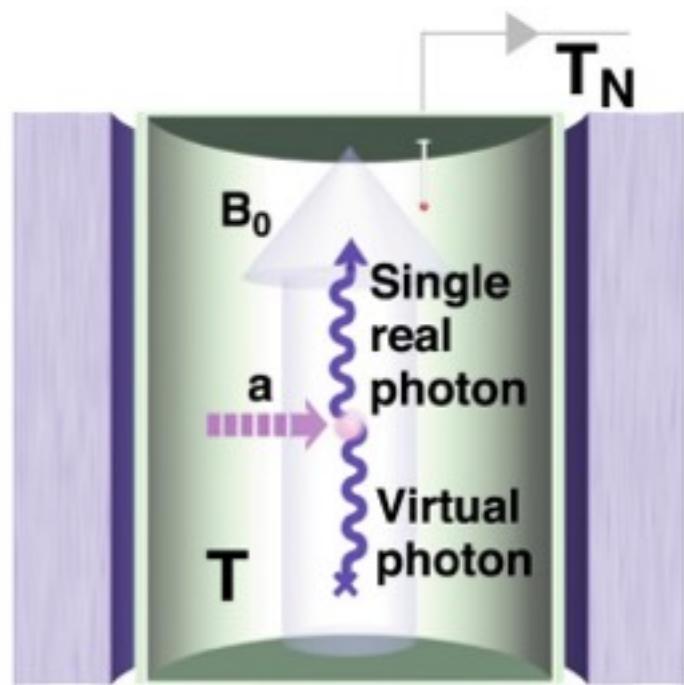
$$\rho_a = \frac{1}{2}(\dot{a})^2 + \frac{1}{2}m_a^2 a^2 = \frac{1}{2}m_a^2 f_a^2 \theta_0^2$$

$V_{\theta\theta}$

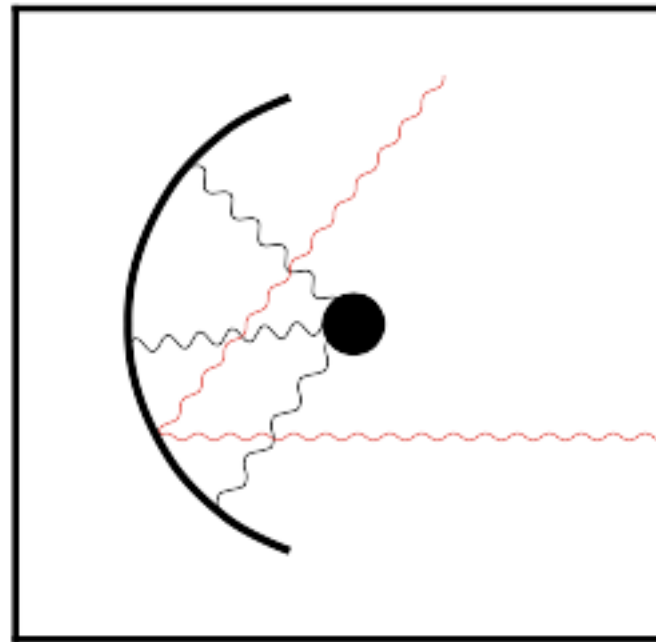


$$\theta_0 = 3.6 \times 10^{-19}$$

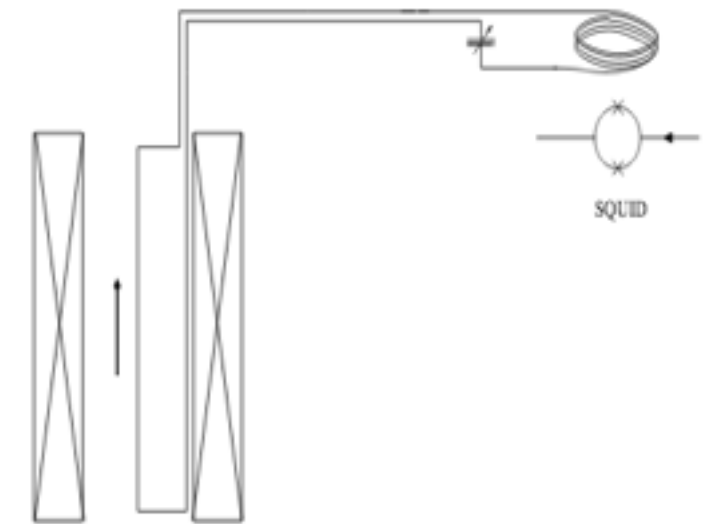
Cavities



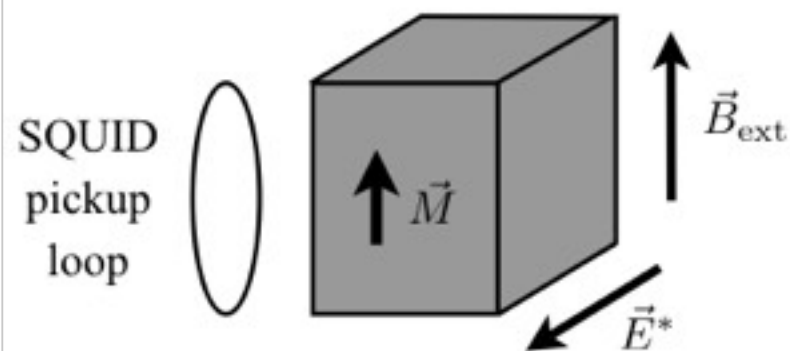
Mirrors



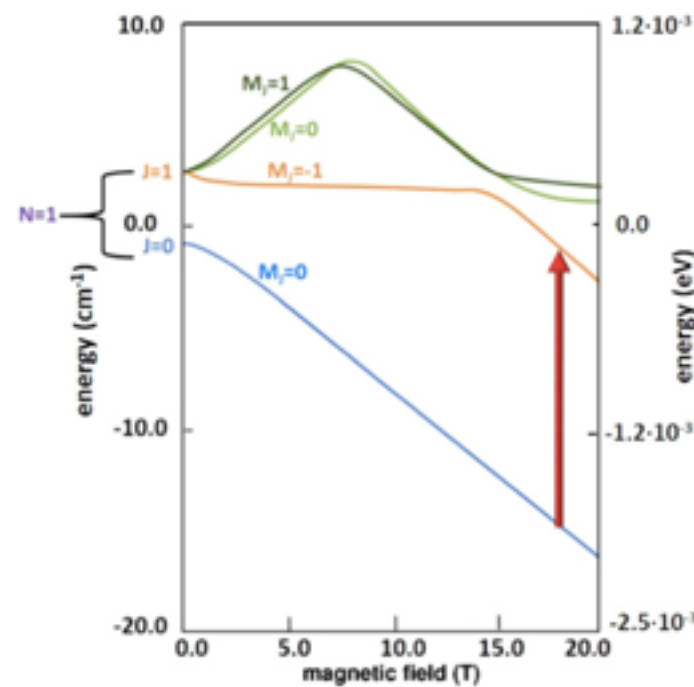
LC-circuit



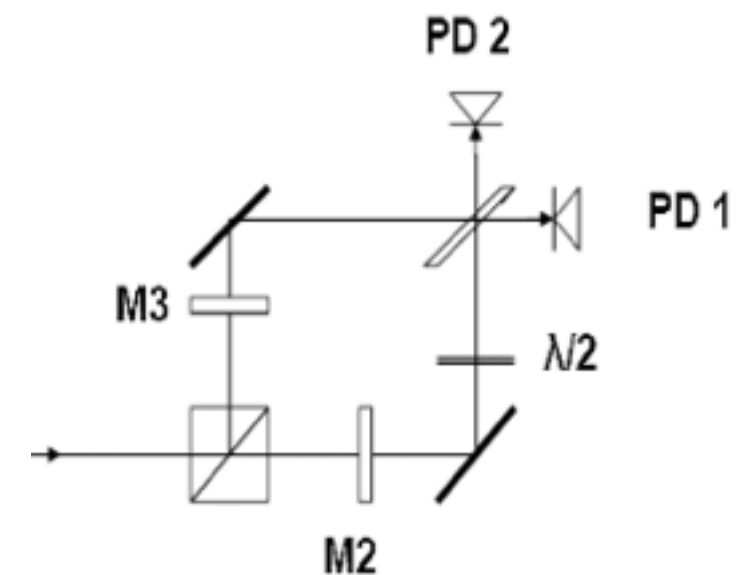
Spin precession



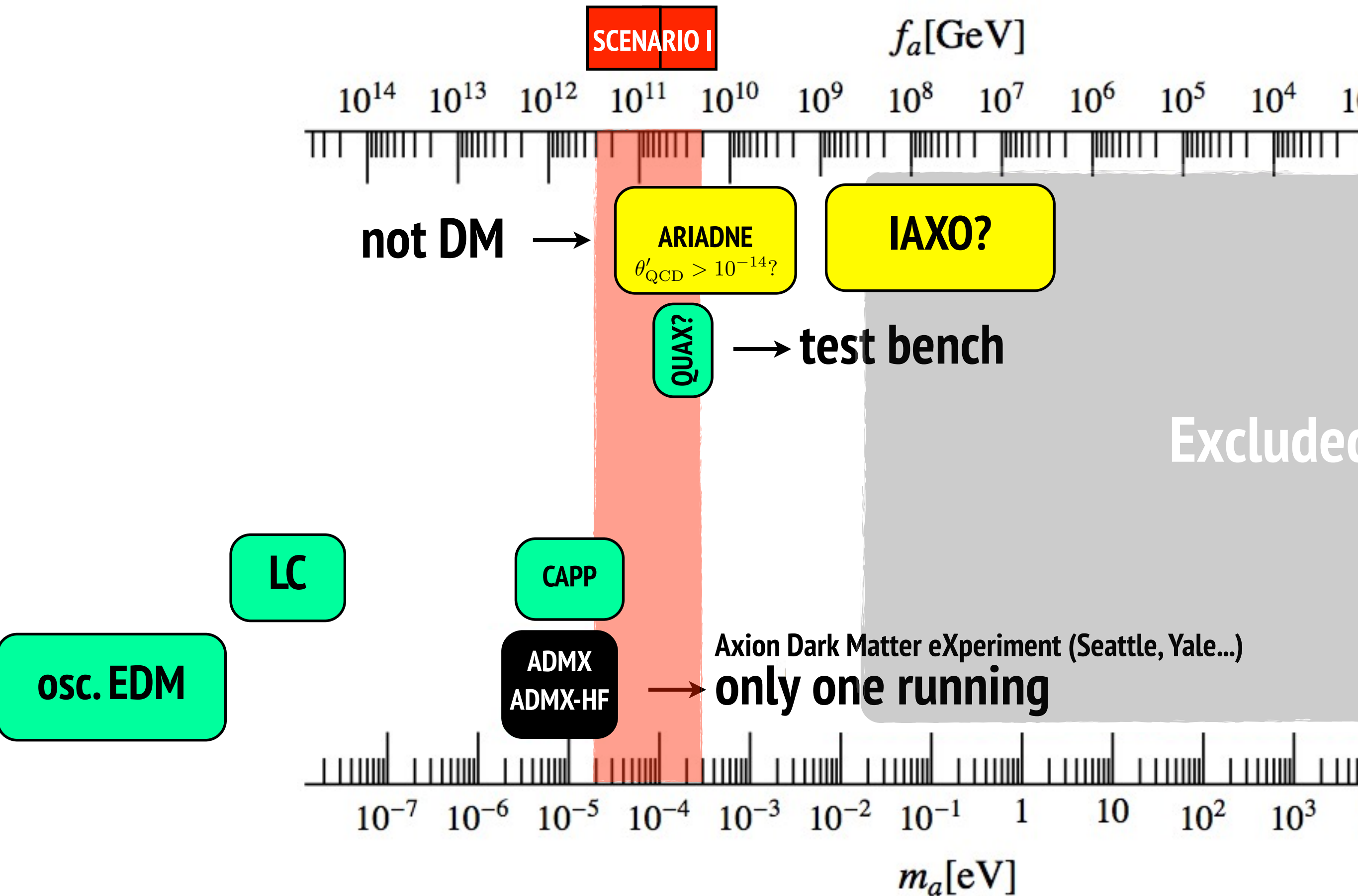
Atomic transitions



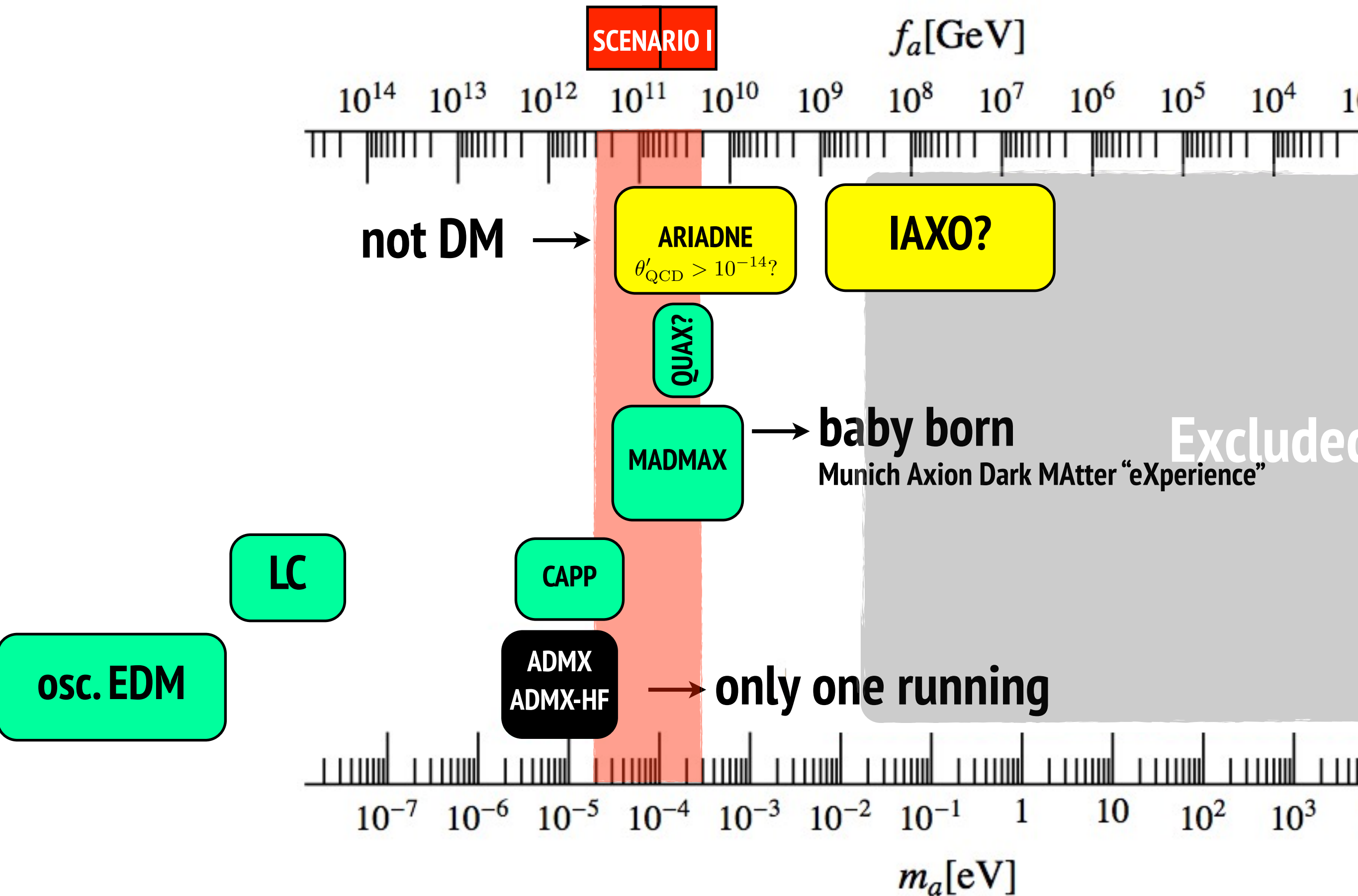
Optical



Axion experiments (target areas)



Axion experiments (target areas)

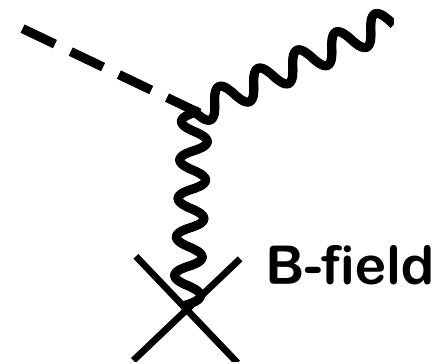


Axion DM in a B-field

$$\mathcal{L}_I = -C_{a\gamma} \frac{\alpha}{2\pi} \frac{a}{f_a} \mathbf{B} \cdot \mathbf{E}$$

- In a static magnetic field, the oscillating axion field generates EM-fields

$$\mathcal{L}_I = - \underbrace{C_{a\gamma} \frac{\alpha}{2\pi} \theta(t)}_{\text{source}} \mathbf{B}_{\text{ext}} \cdot \mathbf{E}$$

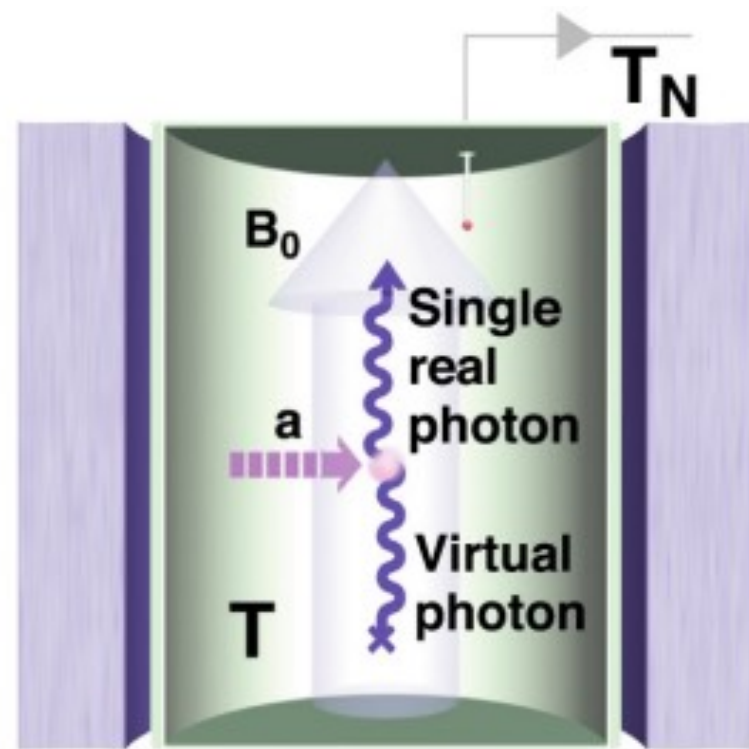


- Electric fields $\mathbf{E}_a = C_{a\gamma} \frac{\alpha \mathbf{B}_{\text{ext}}}{2\pi} \theta_0 \cos(m_a t)$ (amp independent of mass!)
- Oscillating at a frequency $\omega \simeq m_a$

Resonant cavity experiments

- Haloscope (Sikivie 83)

“Amplify resonantly the EM field in a resonant cavity”



Output EM power

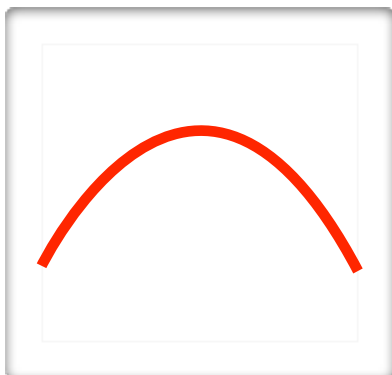
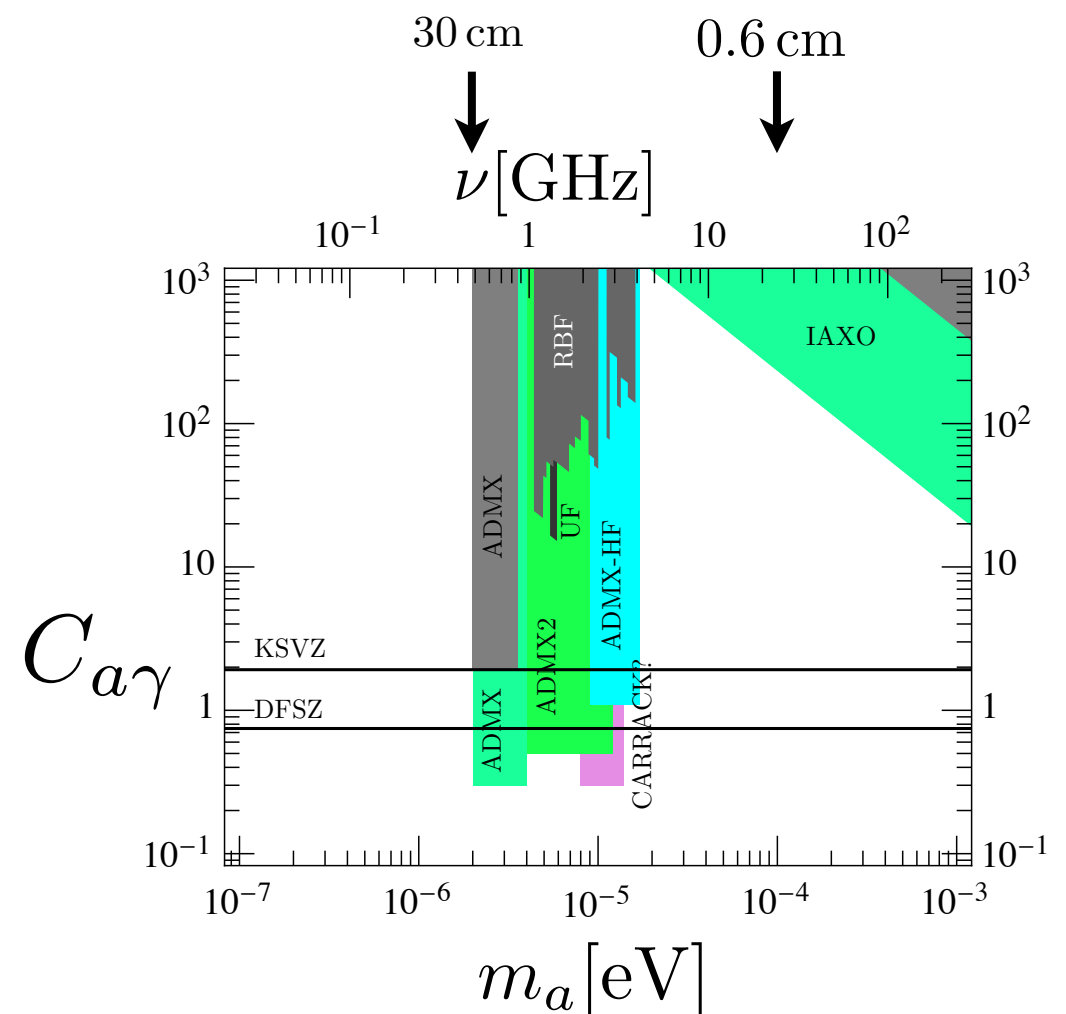
$$P \sim Q |\mathbf{E}_a|^2 (V m_a) \mathcal{G} \kappa$$

on resonance

$$\nu_{\text{res}} = \frac{m_a}{2\pi}$$

- Noise $P_{\text{noise}} = T_{\text{sys}} \Delta \nu_a \propto m_a^2$

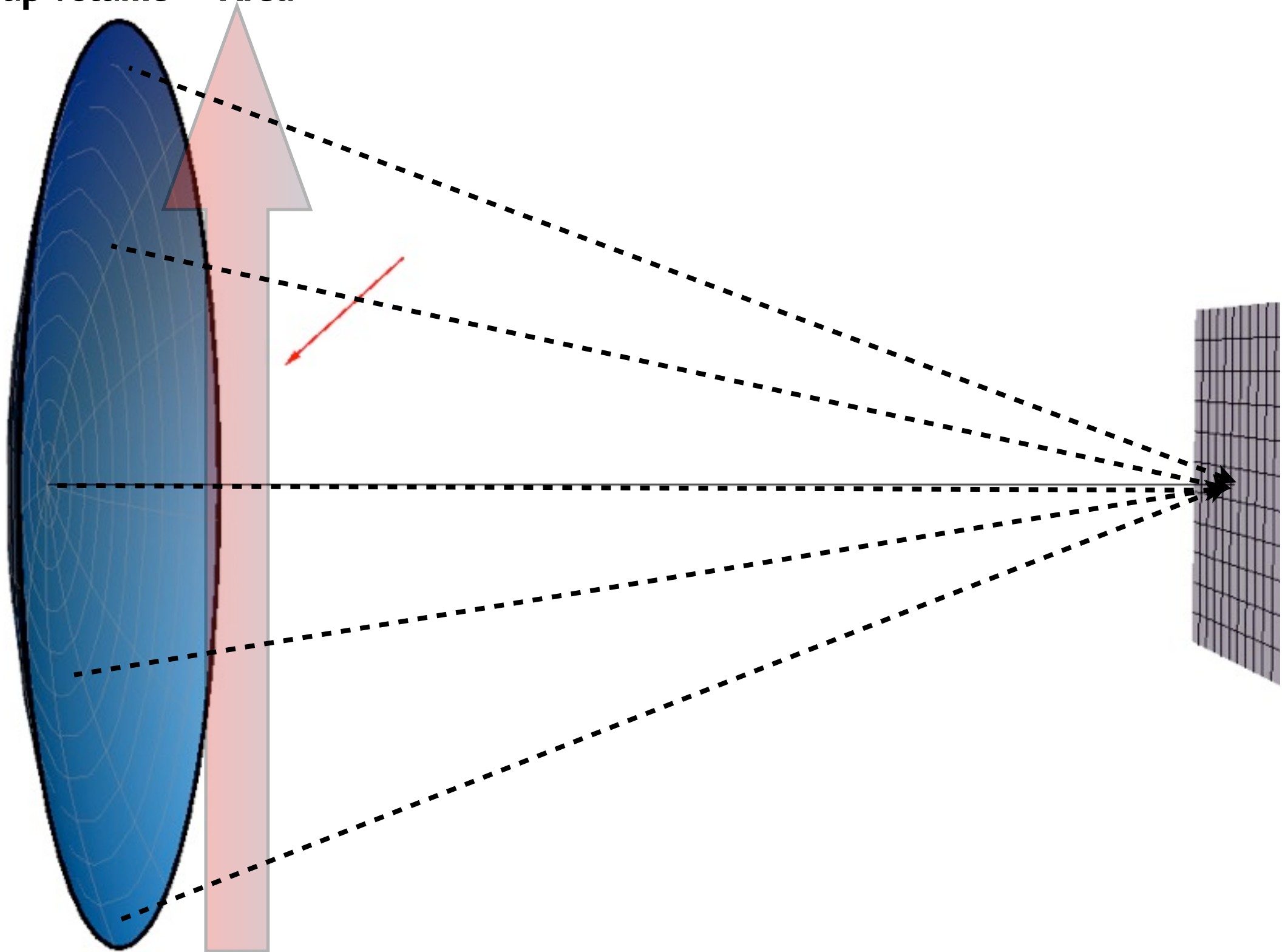
- Signal $(V \propto m_a^{-3}) \quad P_{\text{out}} \propto V m_a \sim \frac{1}{m_a^2}$



pre MADMAX (dish antenna)

- Do not give up volume -> Area

$$D \gg \frac{1}{m_a}$$

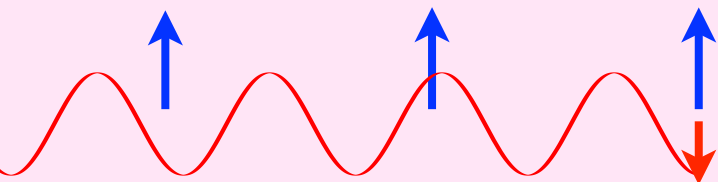


spherical reflecting dish

$$P \sim |\mathbf{E}_a|^2 A$$

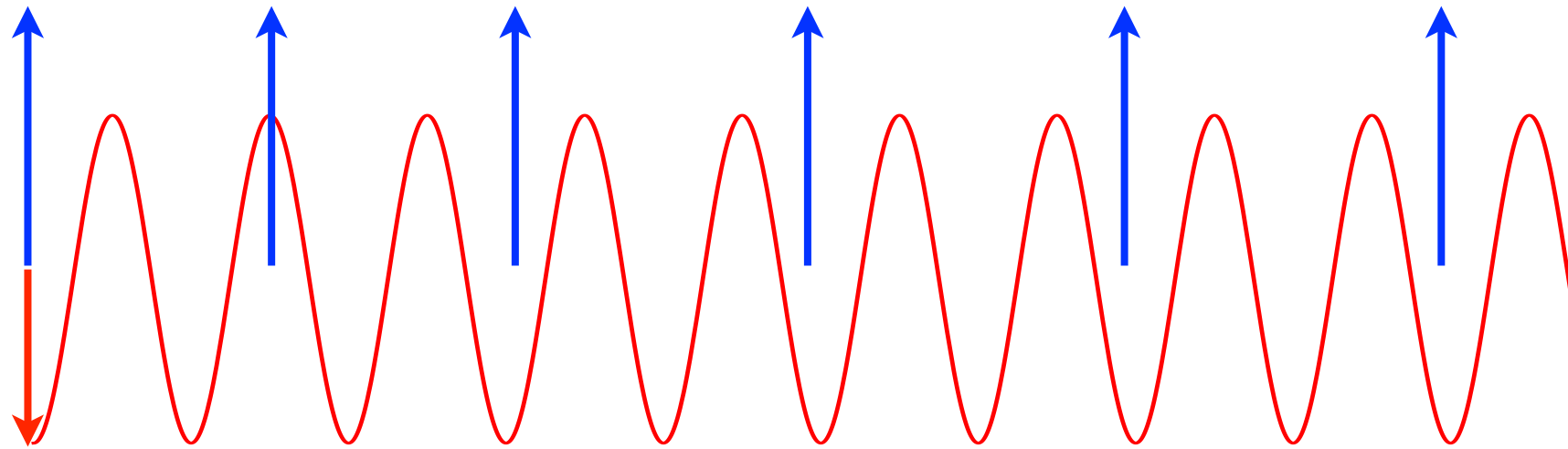
Radiation from a dielectric interface ...

$$\mathbf{E}(t) = \frac{c_\gamma \alpha \theta_0 \mathbf{B}}{2\pi\epsilon} \cos(m_a t)$$



Emitted EM-wave

$$\mathbf{E}(t) = \frac{c_\gamma \alpha \theta_0 \mathbf{B}}{2\pi} \cos(m_a t)$$



← **Boundary conditions!**

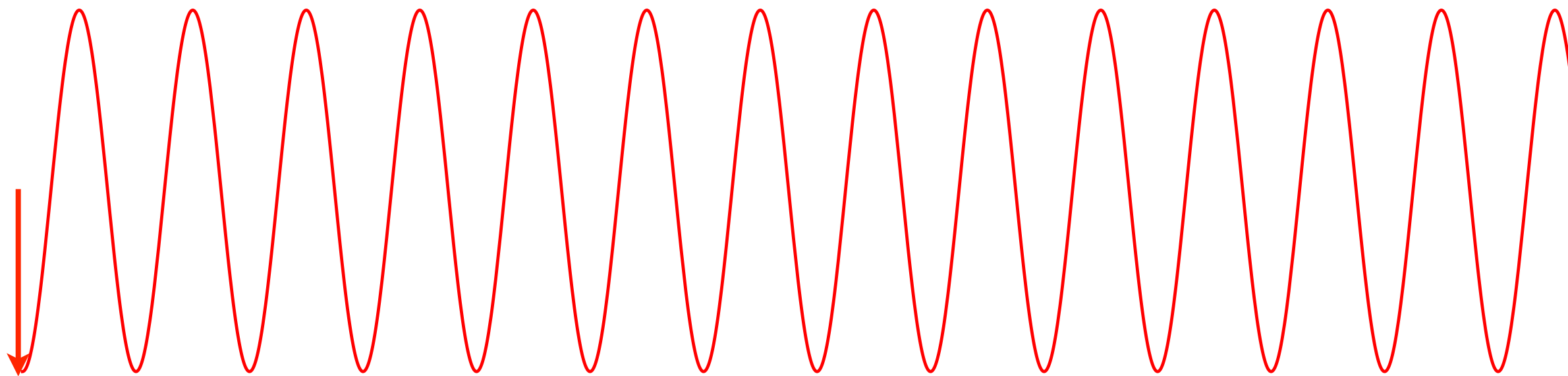
$$\mathbf{E}_{||1} = \mathbf{E}_{||2}$$



Emitted EM-wave

Radiation from a magnetised mirror : Power

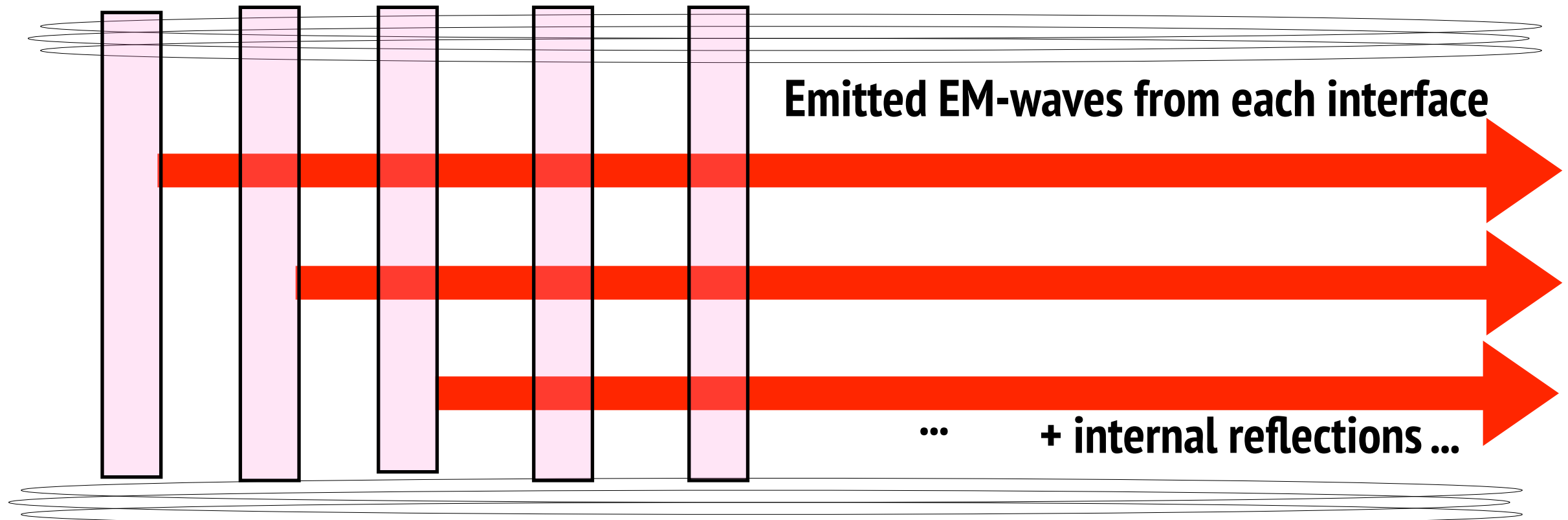
MIRROR



Emitted EM-wave

$$\frac{\text{Power}}{\text{Area}} \sim 2.2 \times 10^{-27} \frac{\text{Watt}}{\text{m}^2} \left(\frac{C_{a\gamma}}{2} \frac{B_{||}}{5\text{T}} \right)^2$$

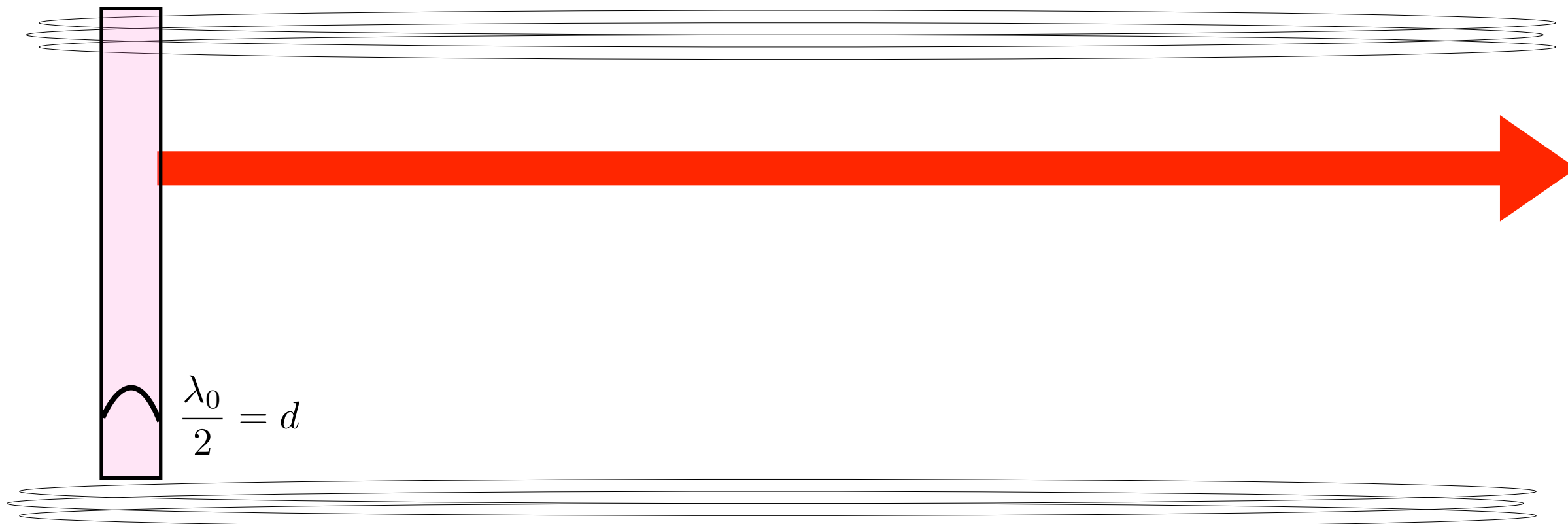
Many dielectrics : MADMAX at MPP Munich



- Emission has large spatial coherence; adjusting plate separation -> coherence

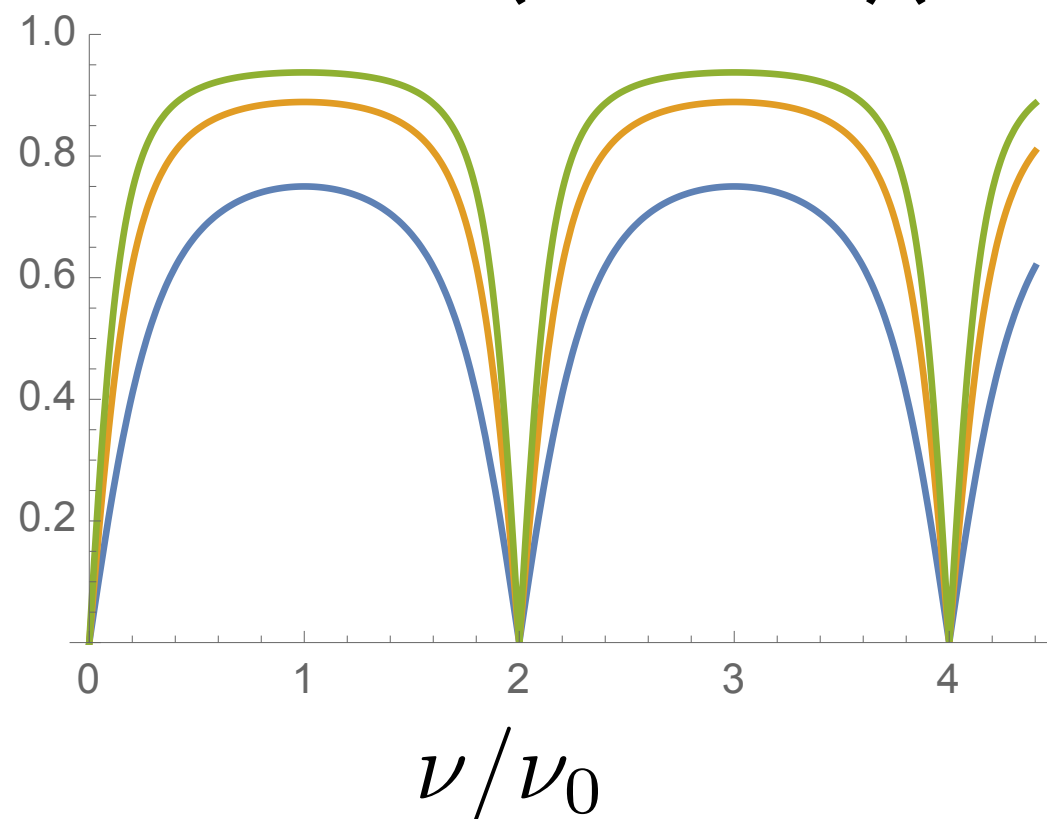
$$\frac{P}{Area} \sim 2.2 \times 10^{-27} \frac{W}{m^2} \left(C_{a\gamma} \frac{B_{||}}{10T} \right)^2 \times \beta^2(\omega) \quad \text{boost factor}$$

One dielectric

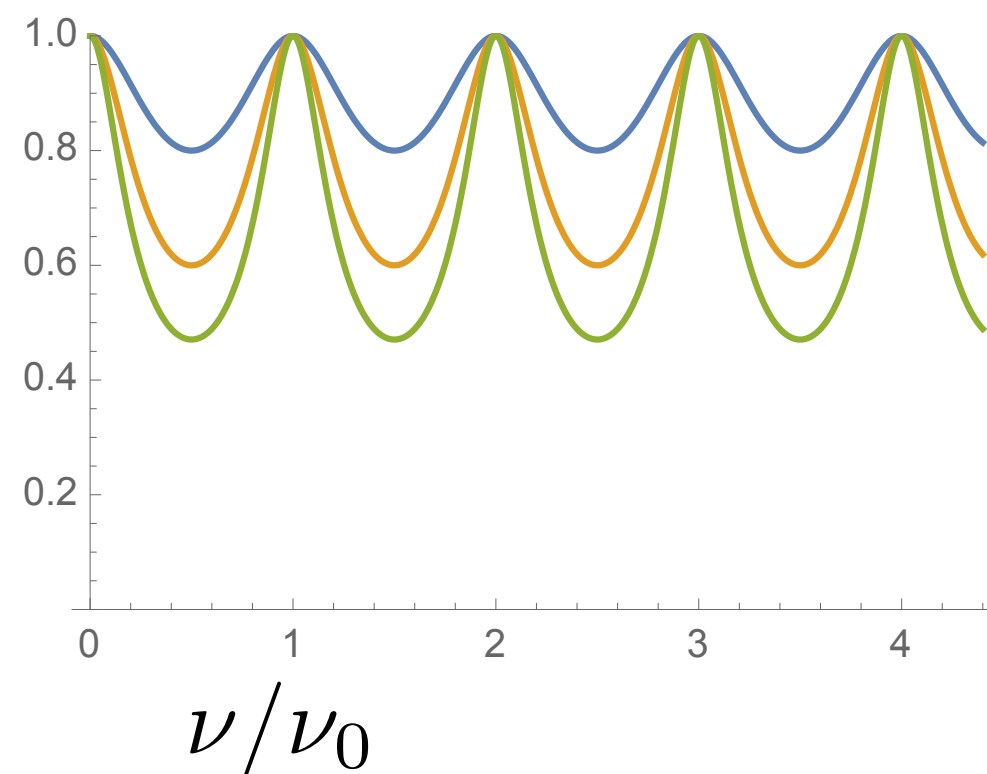


- Frequency response

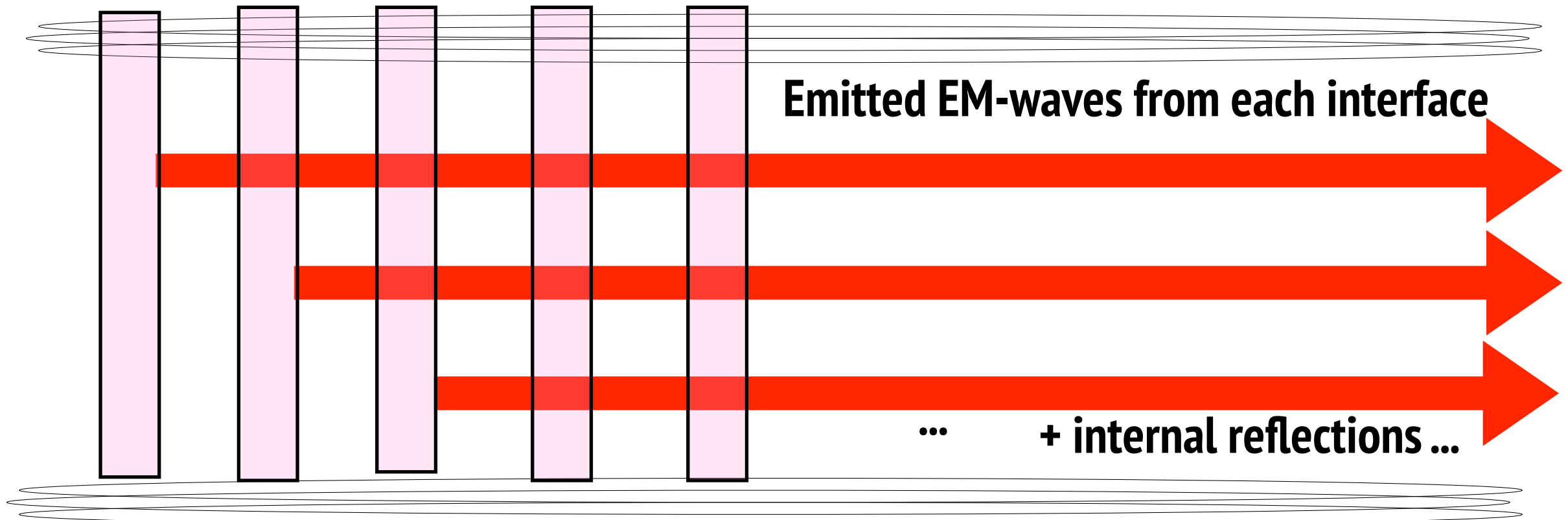
Axion emission (boost factor) ($n=2,3,4$)



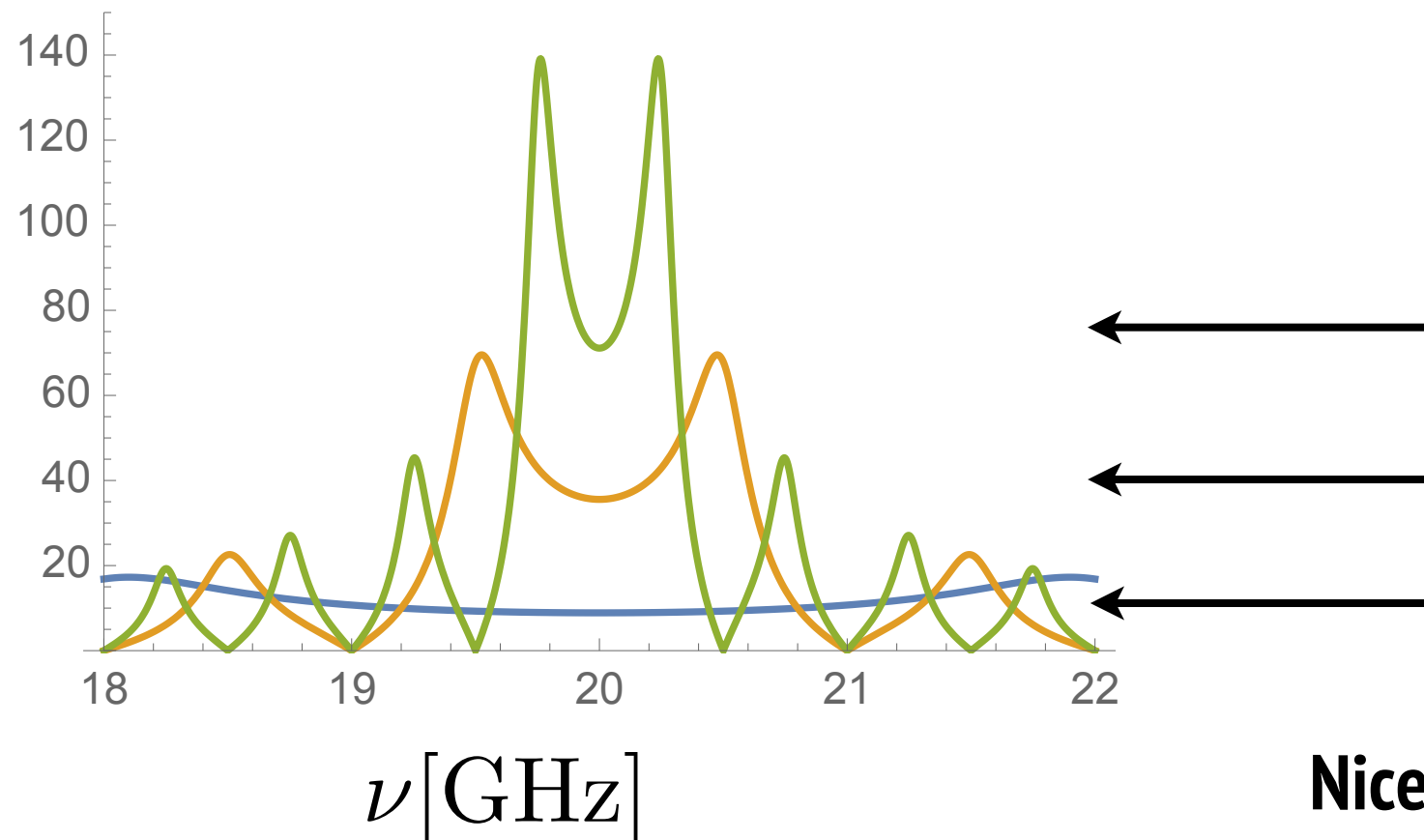
transmission coefficient



Close to ν_0 , many layers



boost factor ($N=10,40,80$; $n=3, \nu_0=20$ GHz)

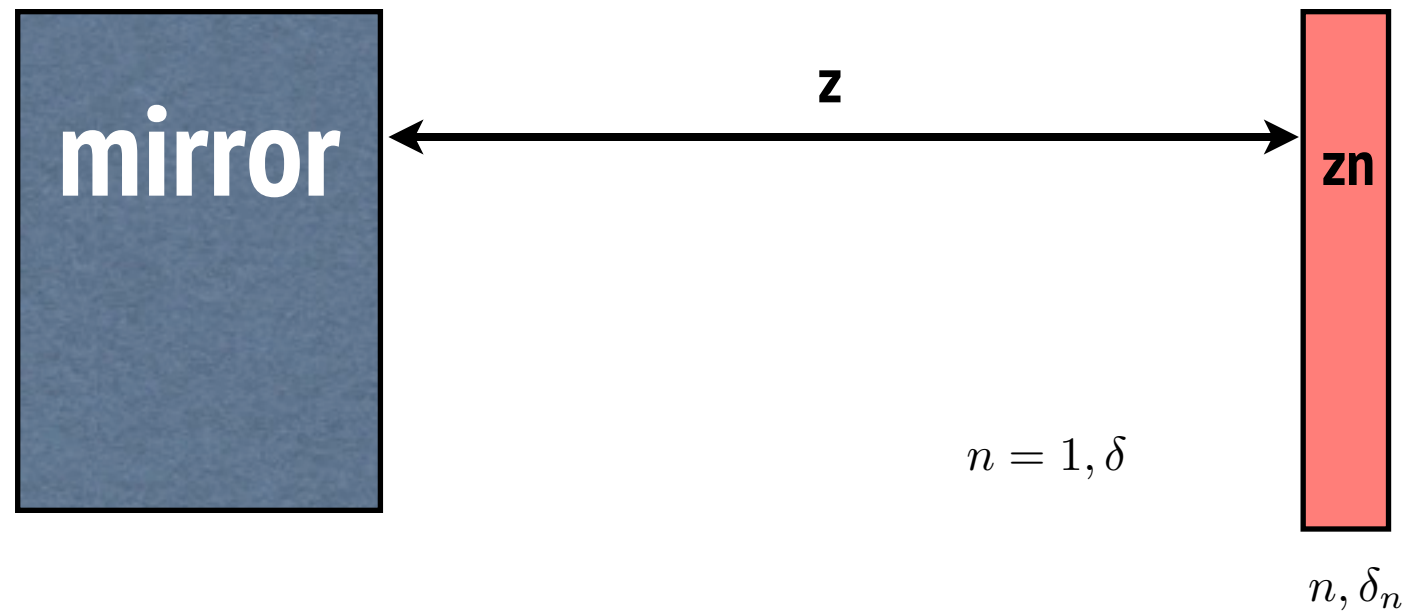


$$\beta_N \sim N \times \beta_1$$

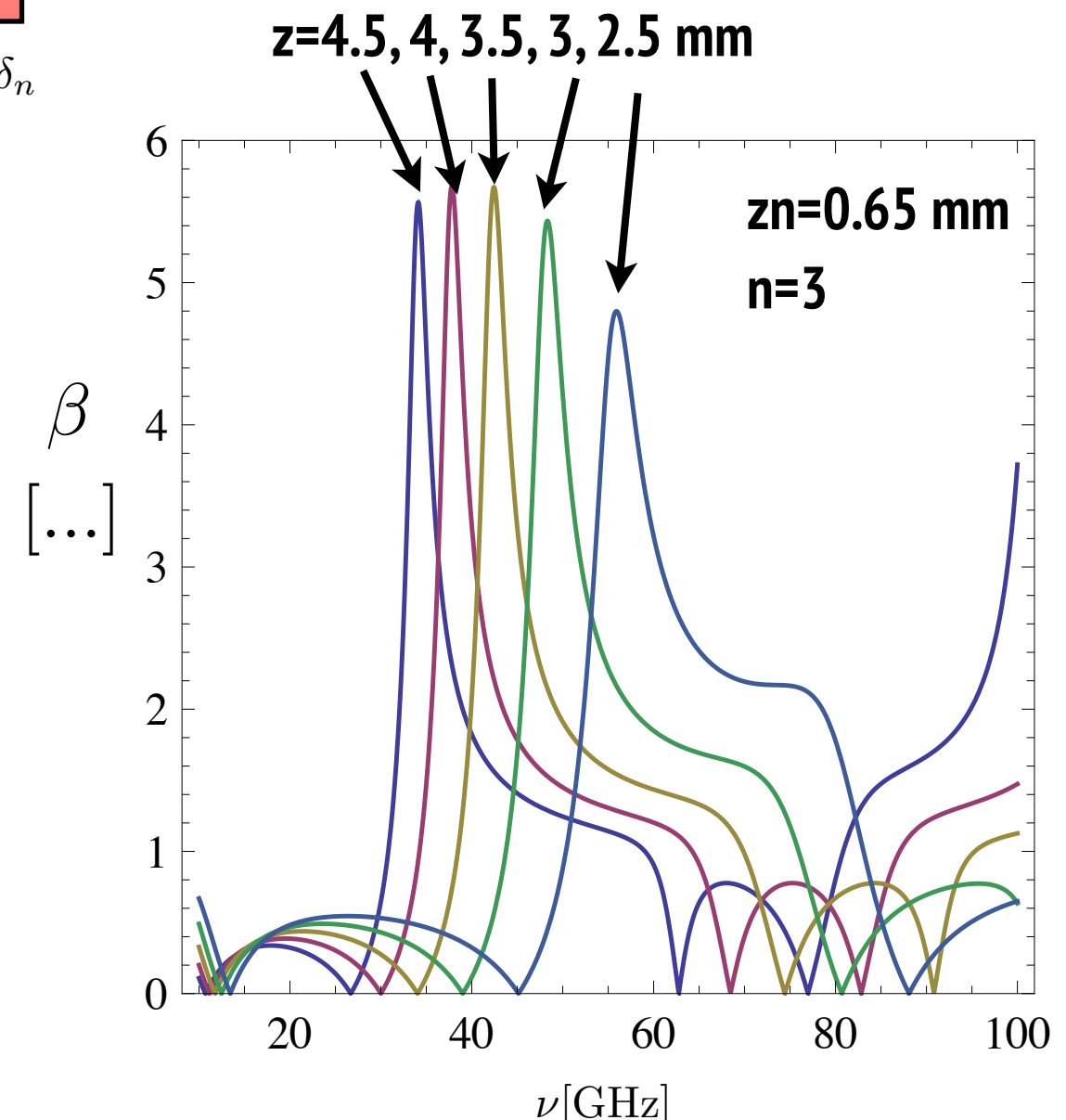
Nice boost, excellent bandwidth!

Close to $2\nu_0$, cavity effects

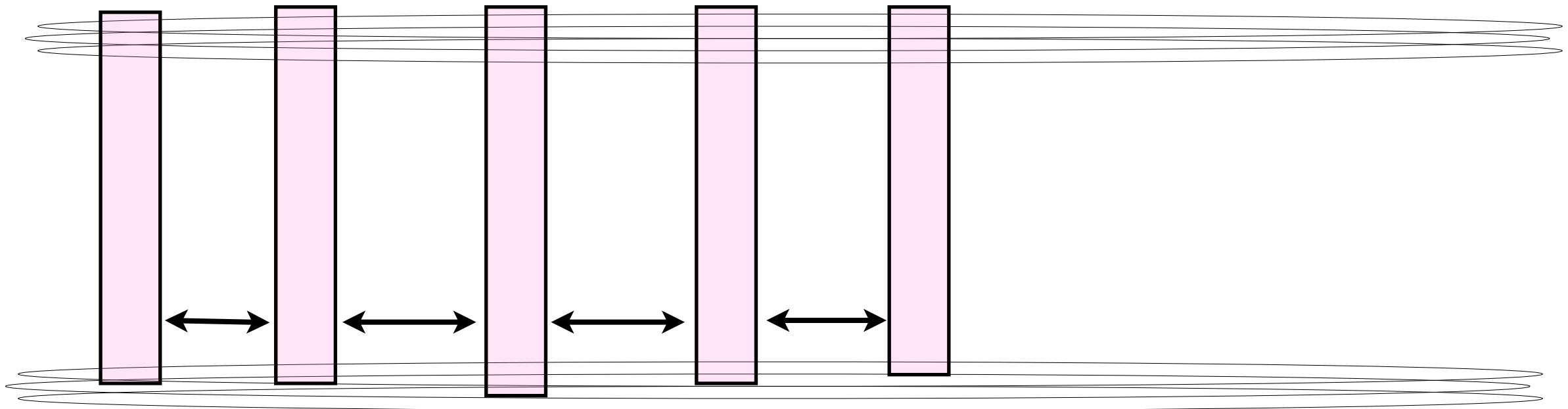
When dielectrics are non-transparent, mild resonances appear that allow higher boosts



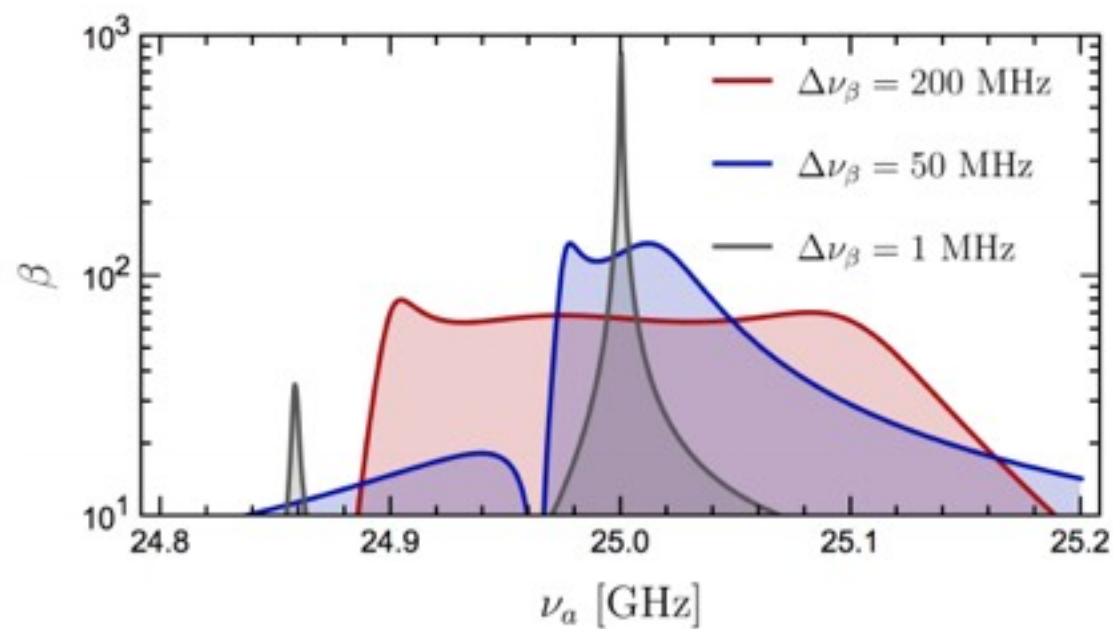
Maximum effect for $z = \frac{\lambda}{2}, z_d = \frac{\lambda'}{4} = \frac{\lambda}{4n}$



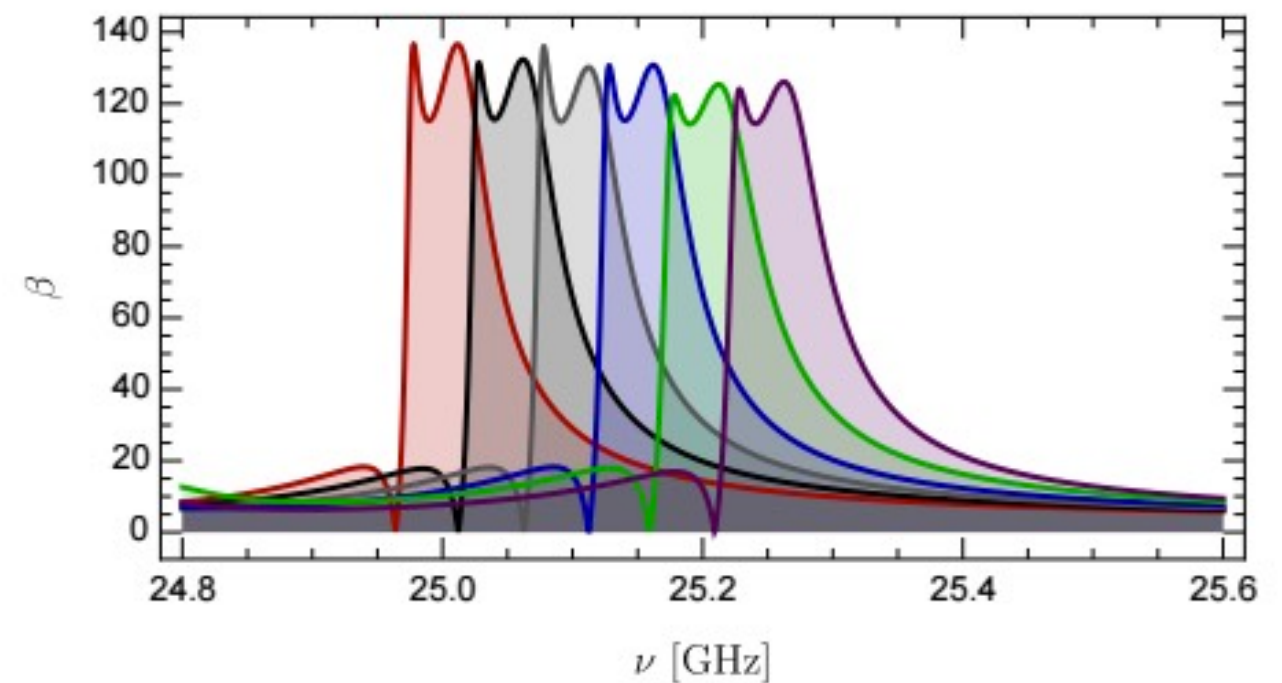
Distances between layers allow tuning (transfer matrix formalism)



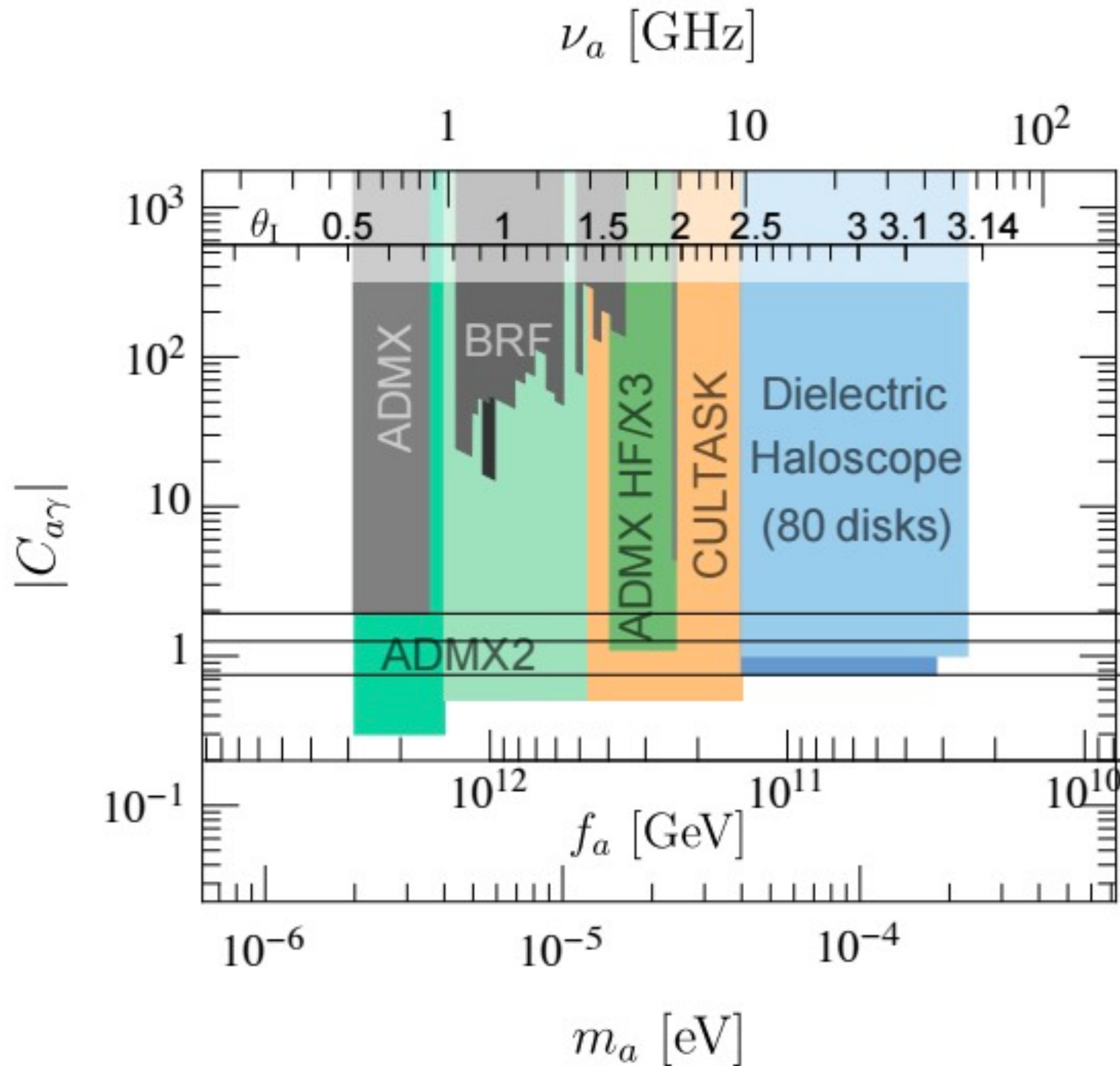
Broadband vs narrowband



central frequency



Potential reach



$A = 1 \text{ m}^2$

$B\text{-field} = 10 \text{ T}$

80disk $n=5$

- why can it work?

Large volume + (relatively) simple tuning
long measurements broadband
short measurements narrowband

- What can go wrong?

- 10T dipole-like Magnet, 1 m² aperture !**
- 1 m² dielectrics (tiling smaller pieces)**
- Tolerances (several micron distances)**
- diffraction**

Conclusions

- Strong CP problem and dark matter motivate Axions
- Most predictive model (N=1) mass~ 0.1 meV ($f_a \sim 10^{11}$ GeV)
- Many experimental efforts, solid player missing in that range
- MW emission from interfaces is weak , make layered haloscope
- Munich Axion Dark MAtter eXperiment