

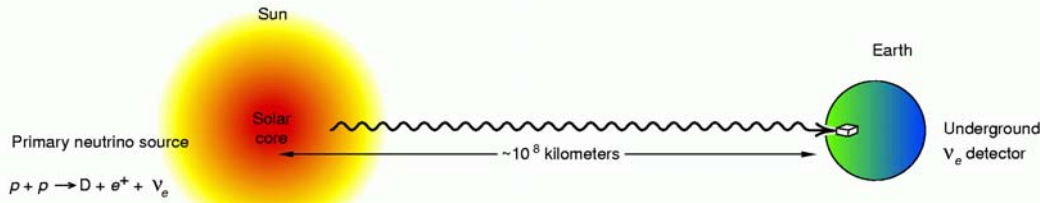
Neutrino Masses in Astrophysics and Cosmology

International School on
Astroparticle and Neutrino Physics
Villa Cipressi, Varenna, Italy
10 - 15 June 2002

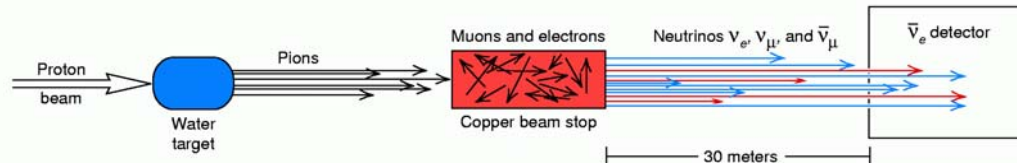
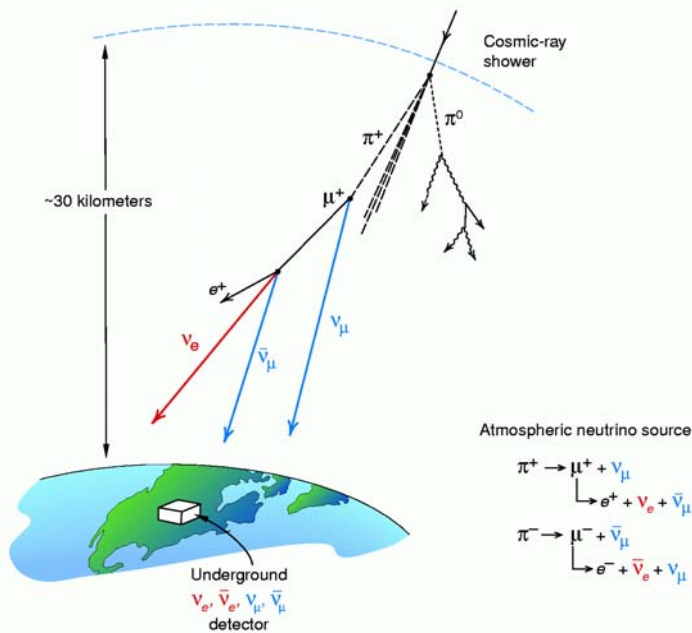
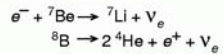


Georg Raffelt (Max-Planck-Institut für Physik, München)

Evidence for Neutrino Oscillations



Other sources of neutrinos:



Solar Neutrinos

Atmospheric Neutrinos

LSND Experiment

Status of Evidence for Neutrino Oscillations

System	Atmospheric	Solar	LSND
Channel	$\nu_\mu \rightarrow \nu_\tau$	$\nu_e \rightarrow \nu_{\mu\tau}$	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$
$\delta m^2 / \text{eV}^2$	$(1.5 - 4) \times 10^{-3}$	LMA $(0.2 - 2) \times 10^{-4}$	0.2-2
$\sin^2 2\theta$	0.9-1	0.2-0.6	0.001-0.03
Status	Established	Established	Unconfirmed
Mutually inconsistent with 3 mass eigenstates			
Test	Long Baseline	KamLAND 2002 ?	MiniBooNE 2003/4?
Simplest interpretation	Three mass eigenstates with $m_1 \ll m_2 \ll m_3 \sim 50 \text{ meV}$ (hierarchical) $m_1 \sim m_2 \sim m_3 \gg 50 \text{ meV}$ (degenerate)		Wrong

What is the absolute neutrino mass scale m_ν ?

Neutrino Masses in Astrophysics and Cosmology



**Time-Of-Flight Sensitivity
with Supernova Neutrinos**



**Neutrino Dark Matter and
Cosmic Structure Formation**



**Neutrino Chemical Potentials,
Big-Bang Nucleosynthesis,
and Flavor Oscillations**



**Highest-Energy Cosmic Rays
and the Cosmic Neutrino Sea**



**Massive Neutrinos and the
Cosmic Baryon Asymmetry**

Time-Of-Flight Sensitivity with Supernova Neutrinos



Sanduleak -69 202

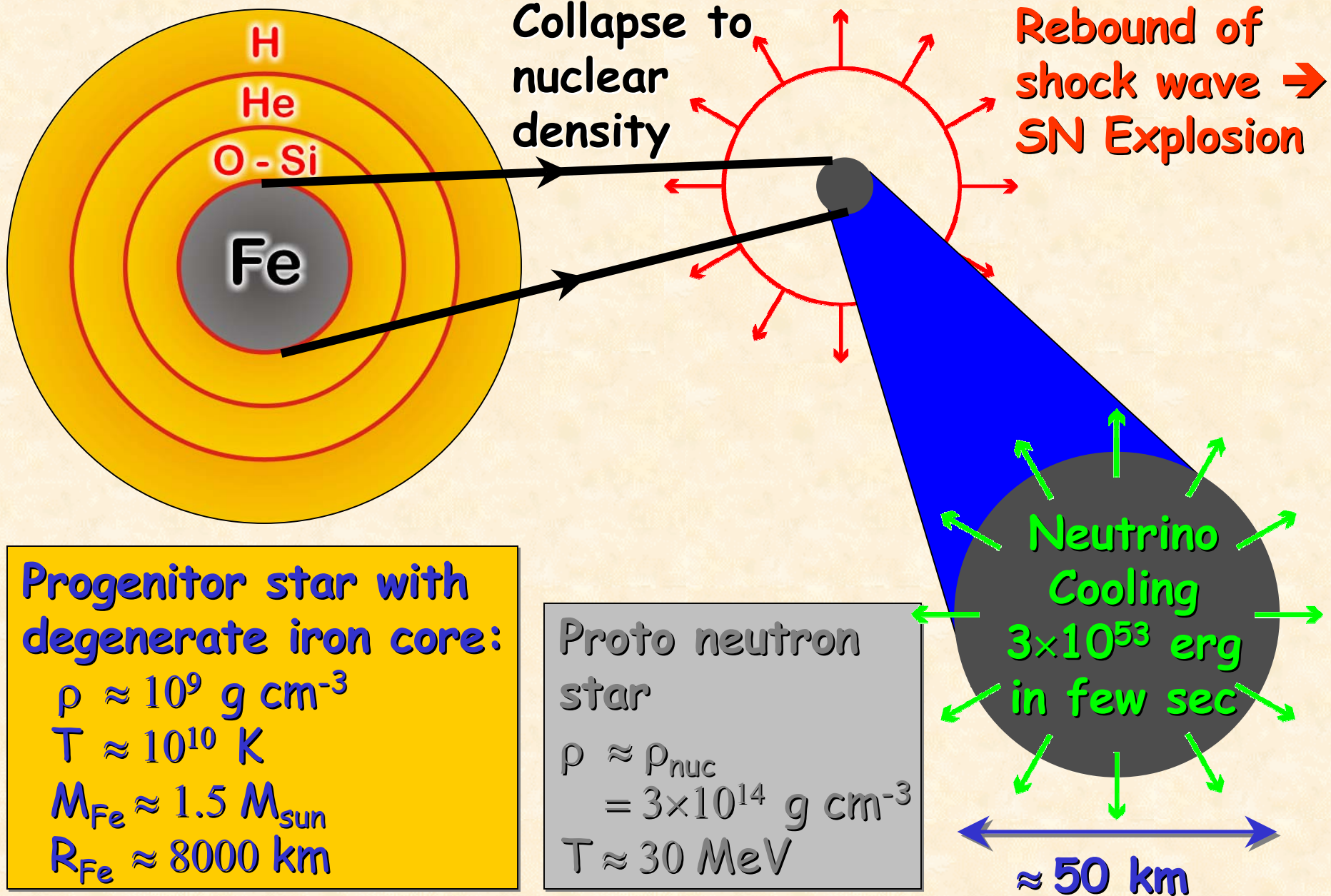


Supernova 1987A

23 February 1987



Stellar Collapse and Supernova Explosion



Core Collapse Supernova Energetics

Liberated gravitational binding energy of neutron star:

$$E_b \approx 3 \times 10^{53} \text{ erg} \approx 17\% M_{\text{SUN}} c^2$$

This shows up as

99%	Neutrinos
1%	Kinetic energy of explosion (1% of this into cosmic rays)
0.01%	Photons (outshine host galaxy)

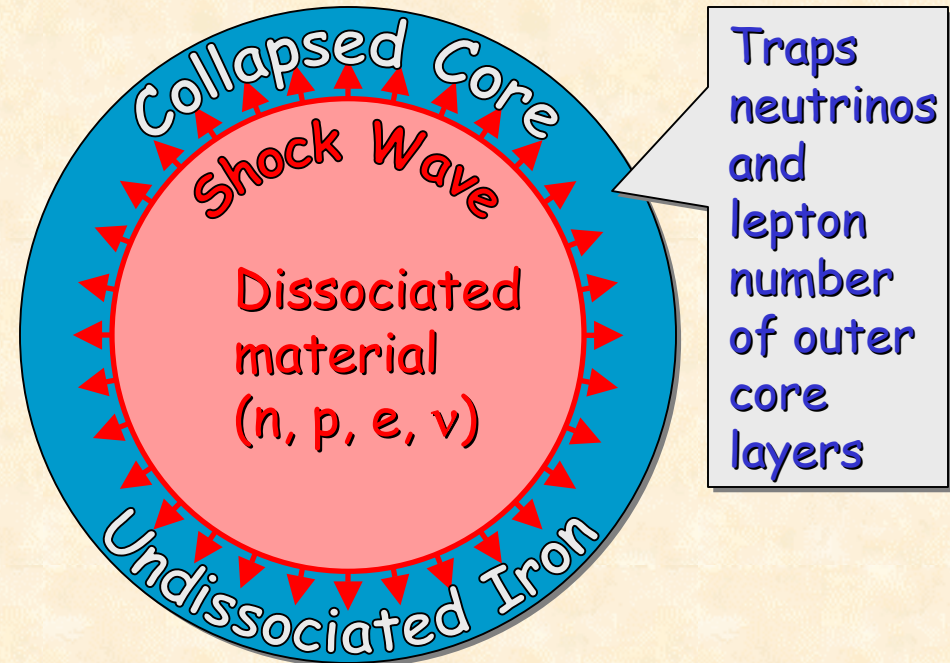
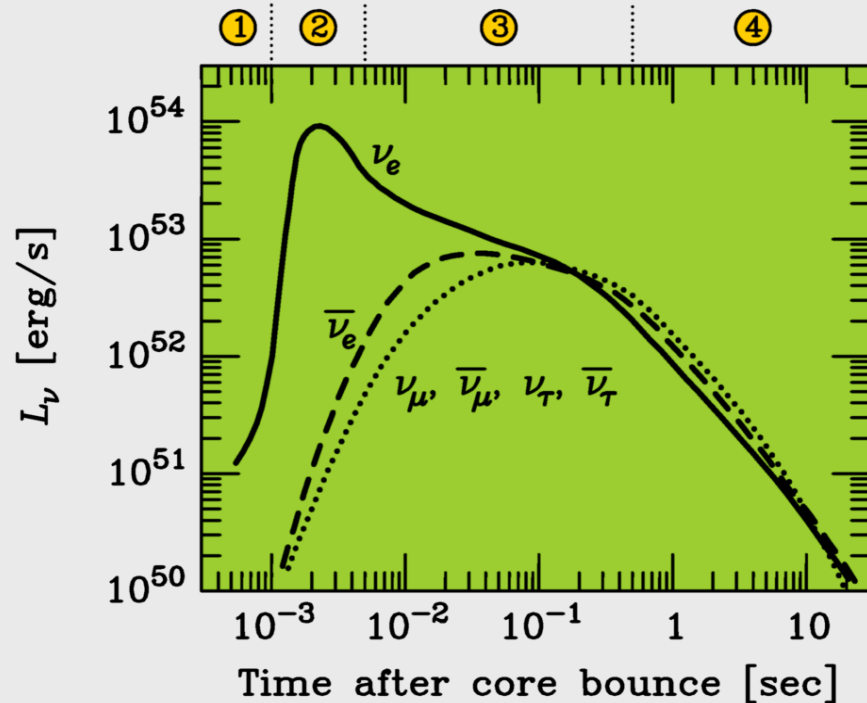
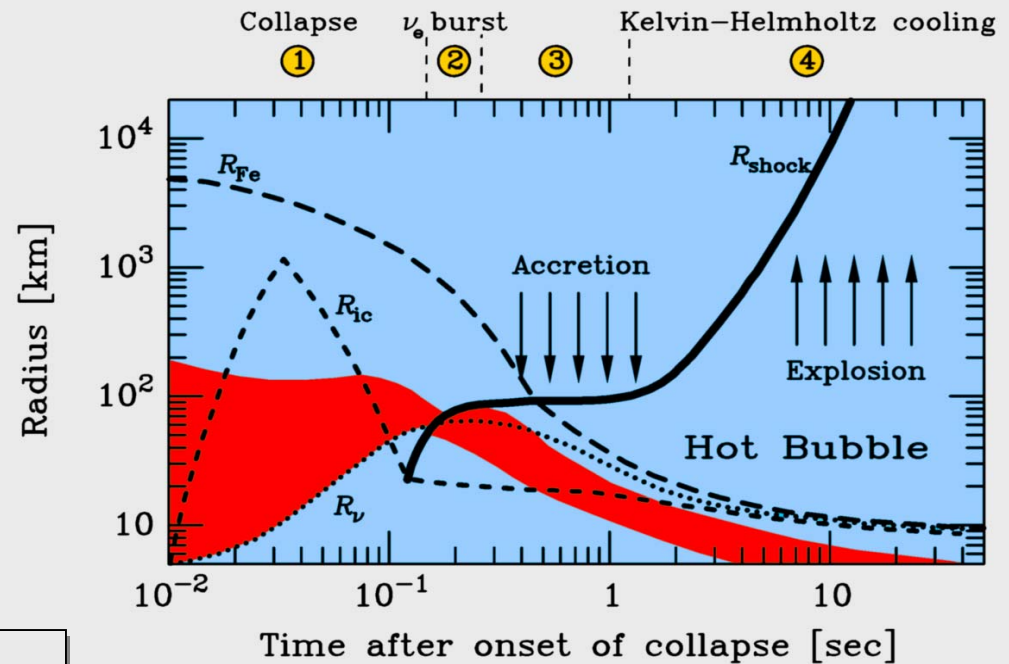
Neutrino luminosity

$$L_\nu \approx 3 \times 10^{53} \text{ erg} / 3 \text{ sec} \approx 3 \times 10^{19} L_{\text{SUN}}$$

While it lasts, outshines the photon luminosity of the entire visible universe!

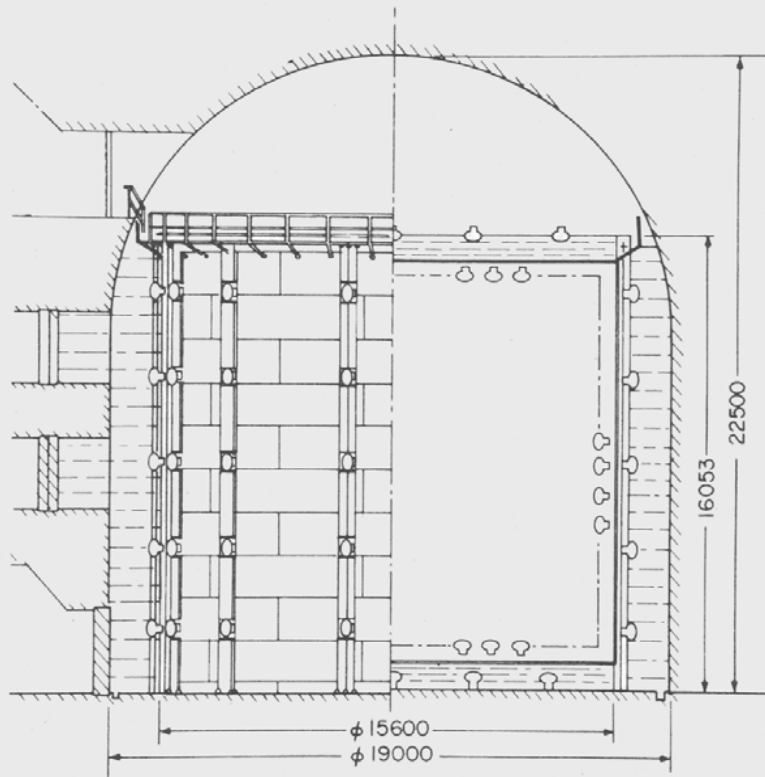
Supernova Neutrino Signal

1. Collapse (infall phase)
2. Shock break out
3. Matter accretion
4. Kelvin-Helmholtz cooling

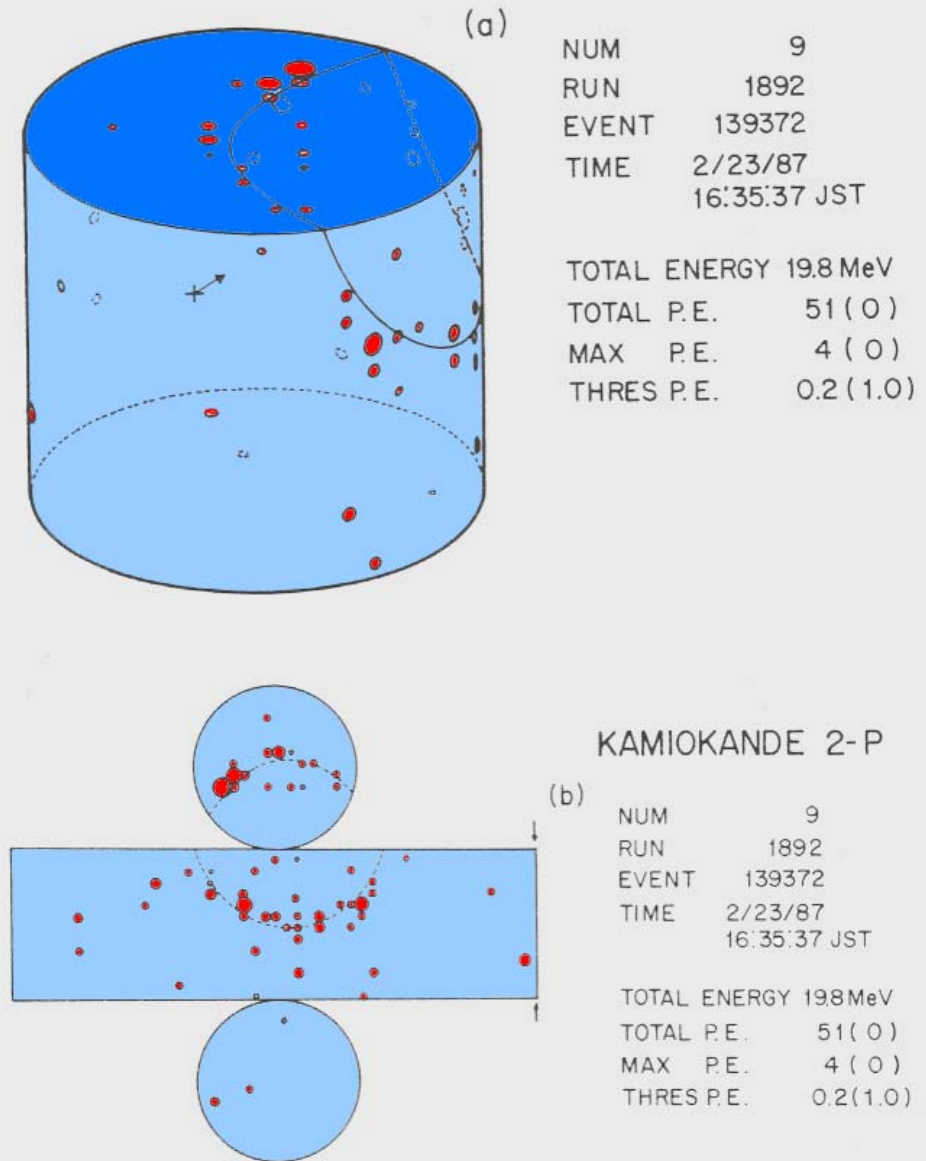


SN 1987A Event No.9 in Kamiokande

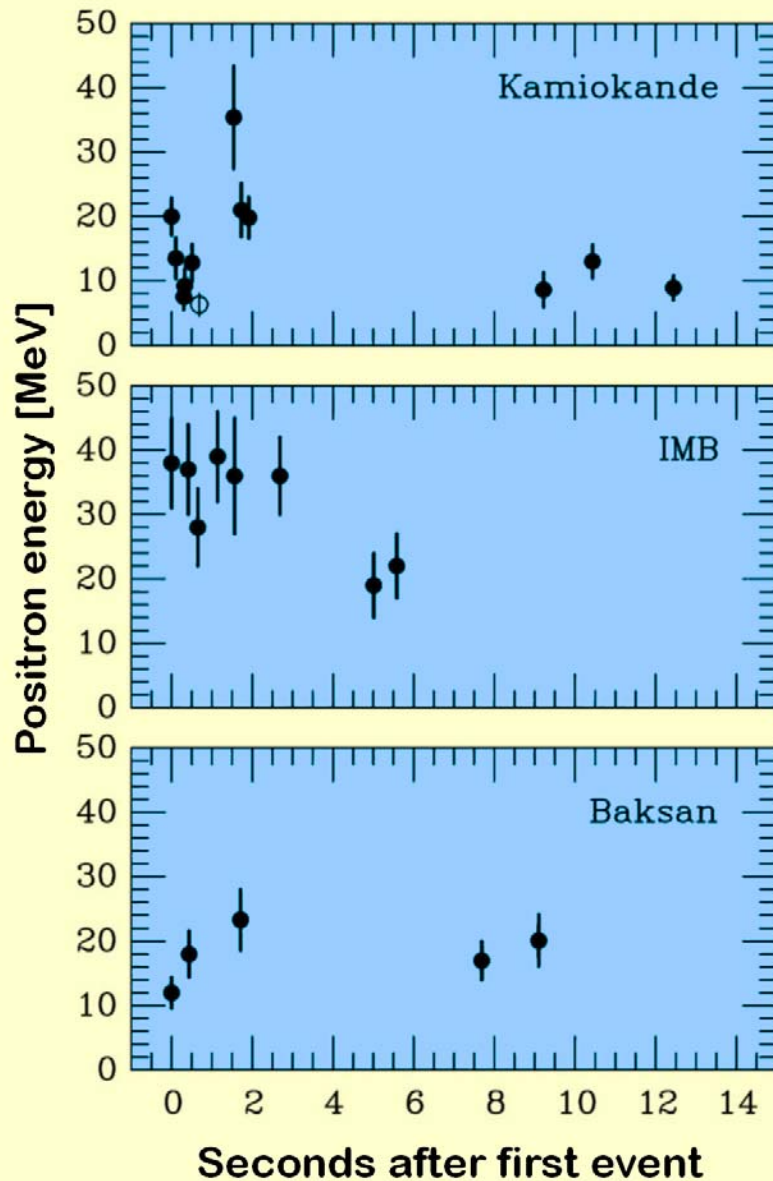
Kamiokande Detector



Hirata et al., PRD 38 (1988) 448



Neutrino Signal of Supernova 1987A



Kamiokande (Japan)
Water Cherenkov detector
Clock uncertainty ± 1 min

Irvine-Michigan-Brookhaven
(USA)
Water Cherenkov detector
Clock uncertainty ± 50 ms

Baksan Scintillator Telescope
(Soviet Union)
Clock uncertainty $+2/-54$ s

Within clock uncertainties,
signals are contemporaneous

Neutrino Limits by Intrinsic Signal Dispersion

Time-of-flight delay
of massive neutrinos

$$\Delta t = 2.57 \text{ s} \left(\frac{D}{50 \text{ kpc}} \right) \left(\frac{10 \text{ MeV}}{E_\nu} \right)^2 \left(\frac{m_\nu}{10 \text{ eV}} \right)^2$$

SN 1987A

$E \approx 20 \text{ MeV}$, $\Delta t \approx 10 \text{ s}$, $D \approx 50 \text{ kpc}$

$m_{\nu_e} < 20 \text{ eV}$

Detailed maximum-likelihood analysis similar result
In 1987 competitive with tritium end-point limits

Future
Galactic SN
(Super-K)

$D \approx 50 \text{ kpc}$, Rise-time 0.01 s

Sensitivity approximately

$m_{\nu_e} \sim 3 \text{ eV}$

Neutrino Astronomy

Events
from a
Supernova
at 10 kpc

many σ

800

8000

370

940

240

400

20

70

330

80

Water Cherenkov

Scintillator

Radio
chemical

SN 1987A

Kamiokande II+III

IMB-3

MACRO

LVD (Gran Sasso)

LSD (Mont Blanc)

Baksan Scintillator Telescope (BST)

Kamland

Borexino

Gallex

GNO

Sage

Homestake

Antares

Baikal

Amanda

SNO

Super K

1970

1975

1980

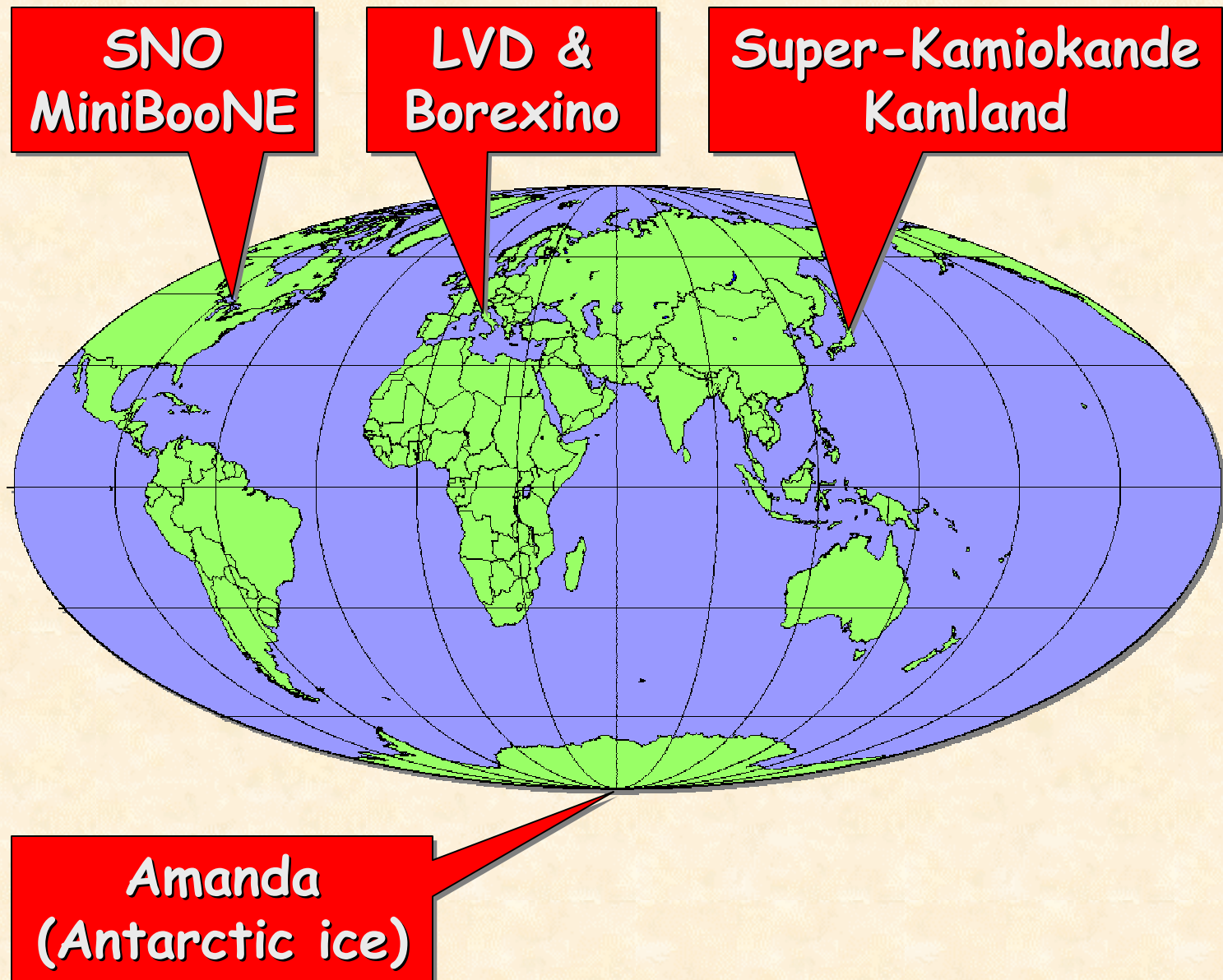
1985

1990

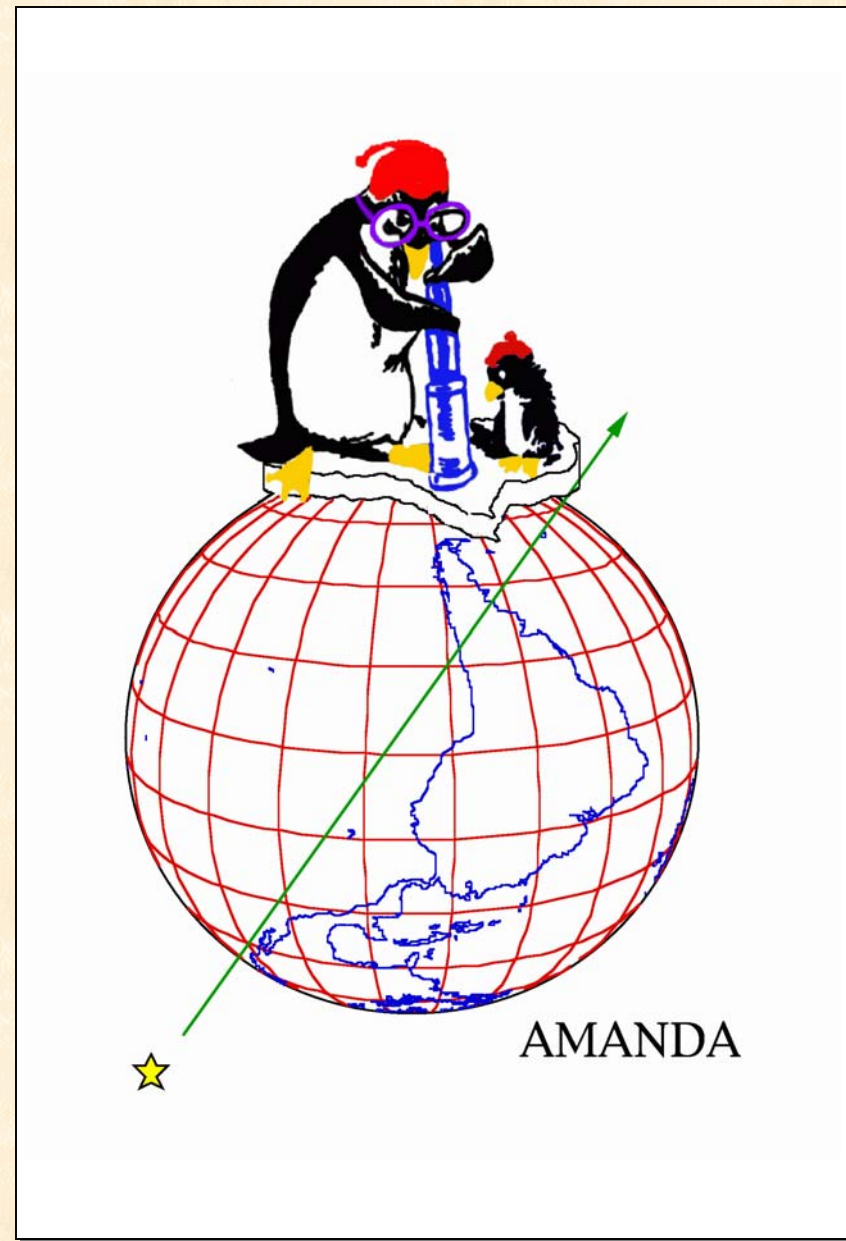
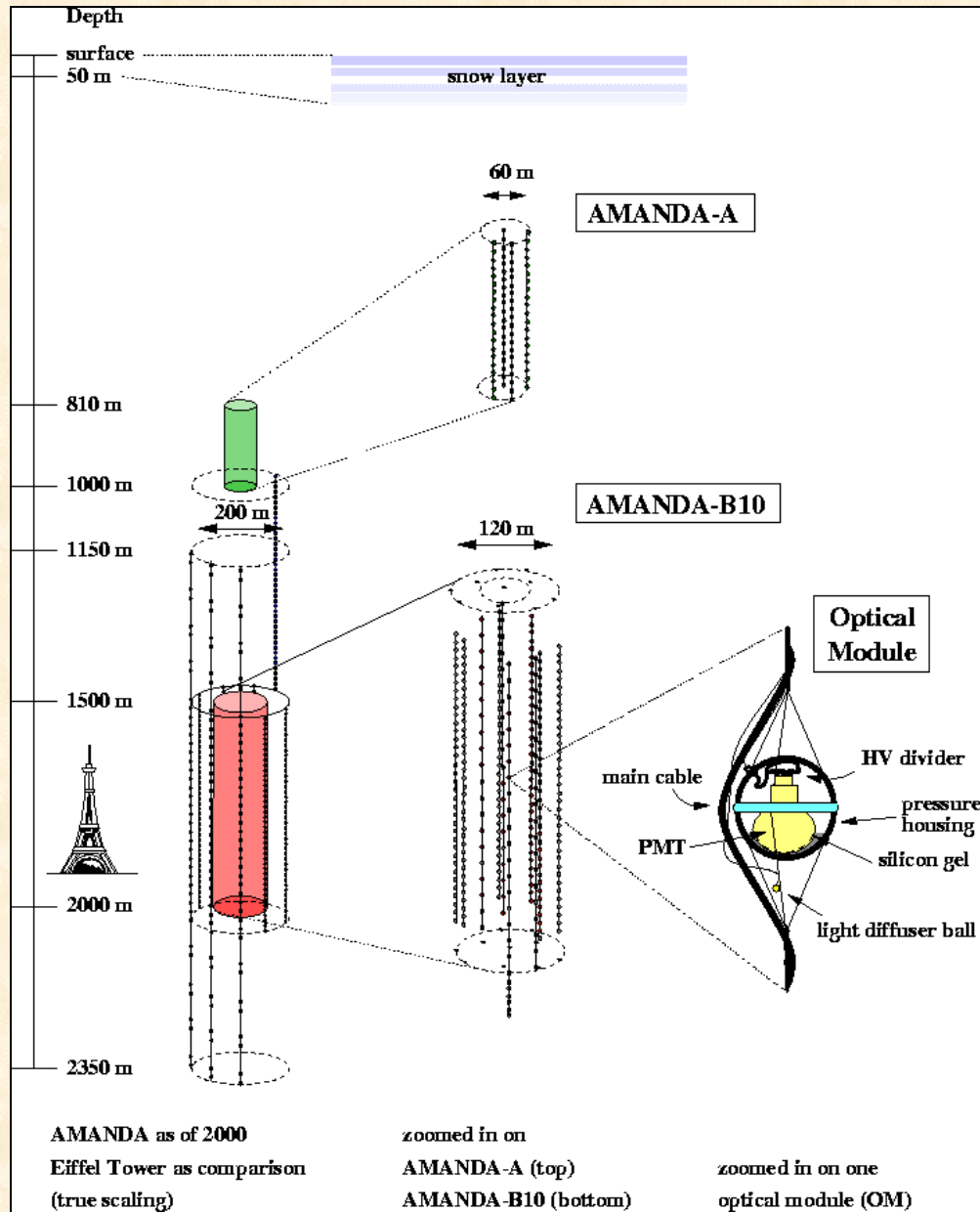
1995

2000

Large Detectors for SN Neutrinos

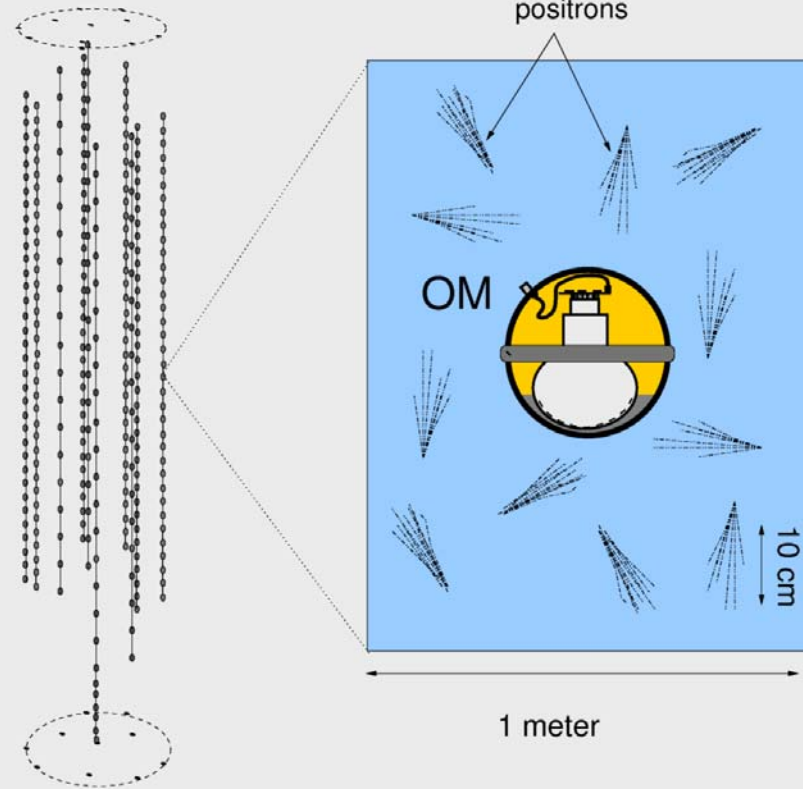


AMANDA - Neutrino Telescope at the Southpole

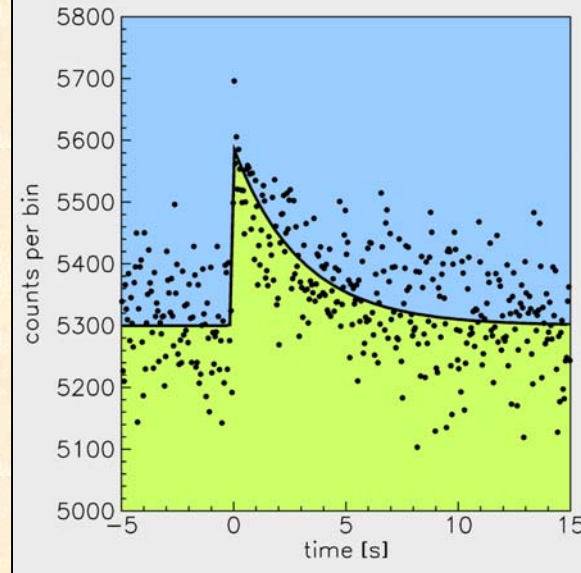


Amanda as a Supernova Detector

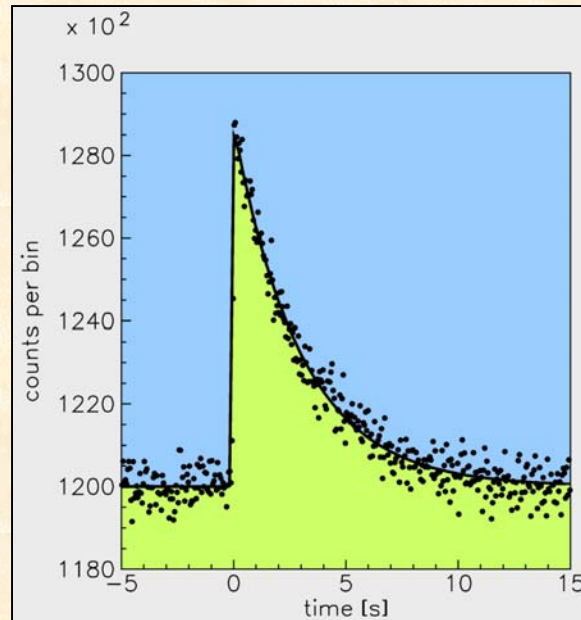
AMANDA-B10



Each optical module (OM) picks up Cherenkov light from its neighborhood. SN appears as correlated "noise" between OMs



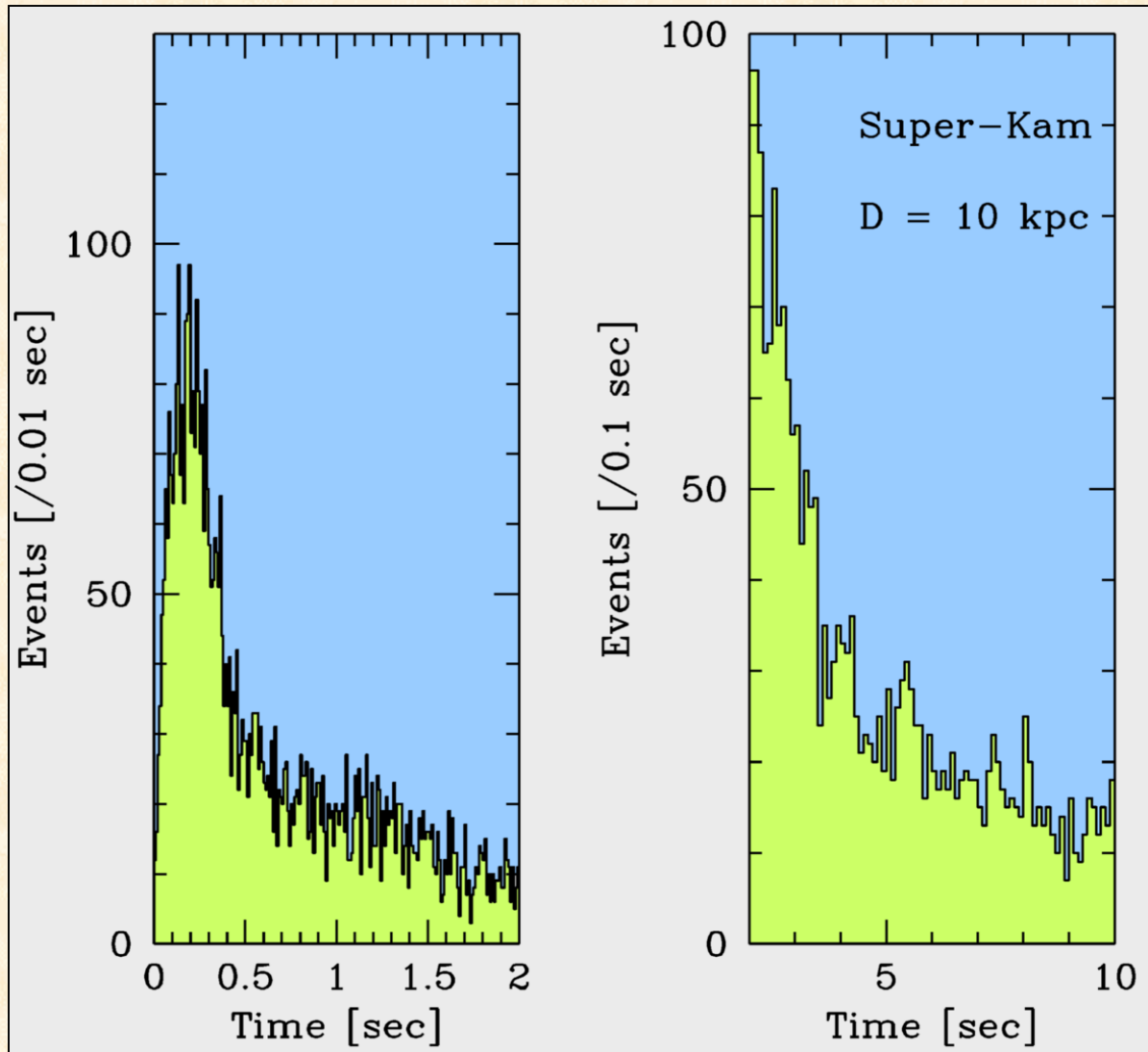
SN @ 8.5 kpc
Signal in
Amanda



SN @ 8.5 kpc
Signal in
Ice Cube

Amanda Coll.
(2001)

Galactic Supernova Signal in Super-Kamiokande

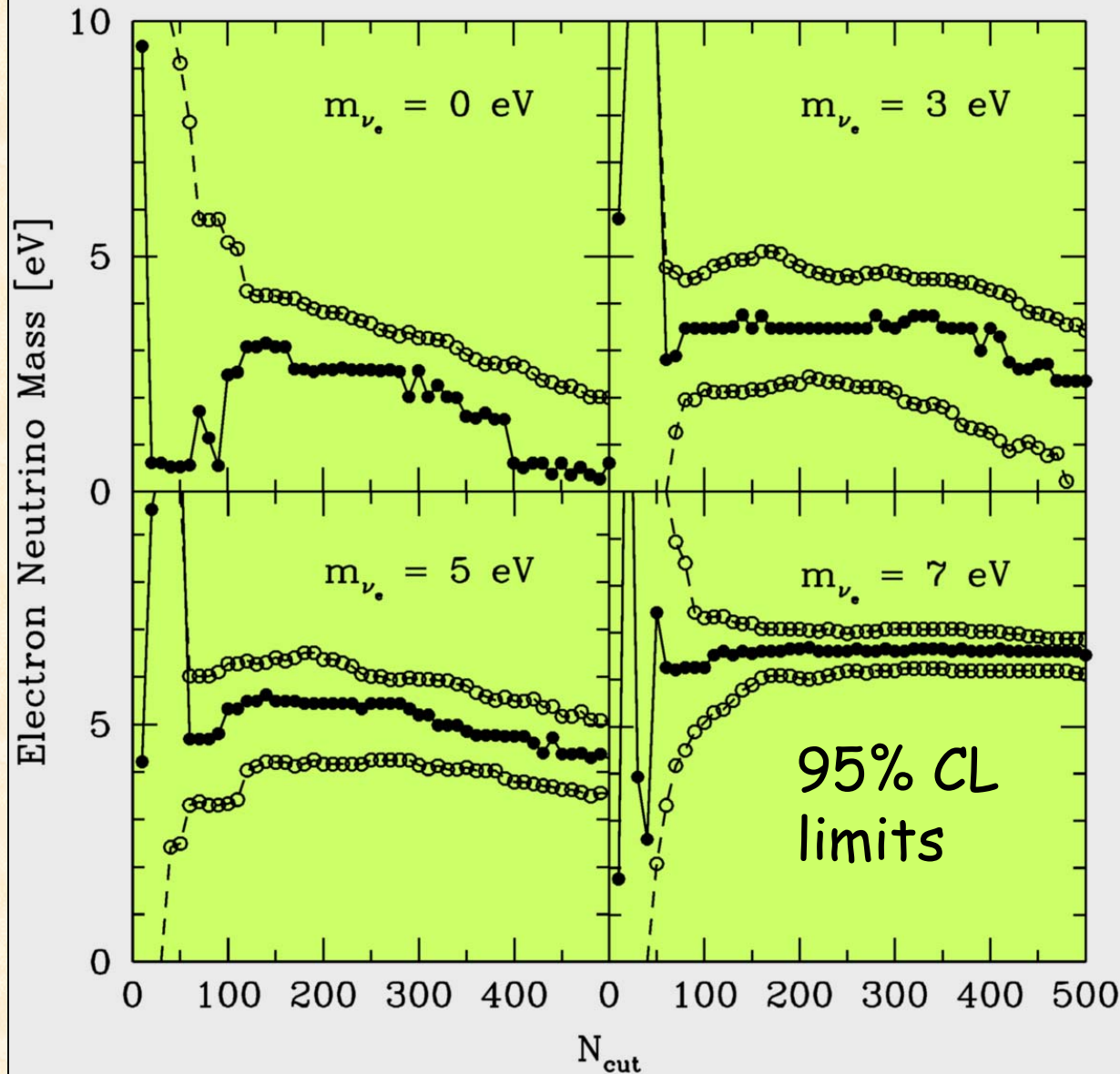


Monte-Carlo simulation
for Super-K signal of
SN at 10 kpc,
based on a
numerical model
with Livermore code

Total of about 8300
events for $t < 18$ s

Totani, Sato, Dalhed & Wilson, ApJ 496 (1998) 216

Neutrino Mass from a Future Galactic SN



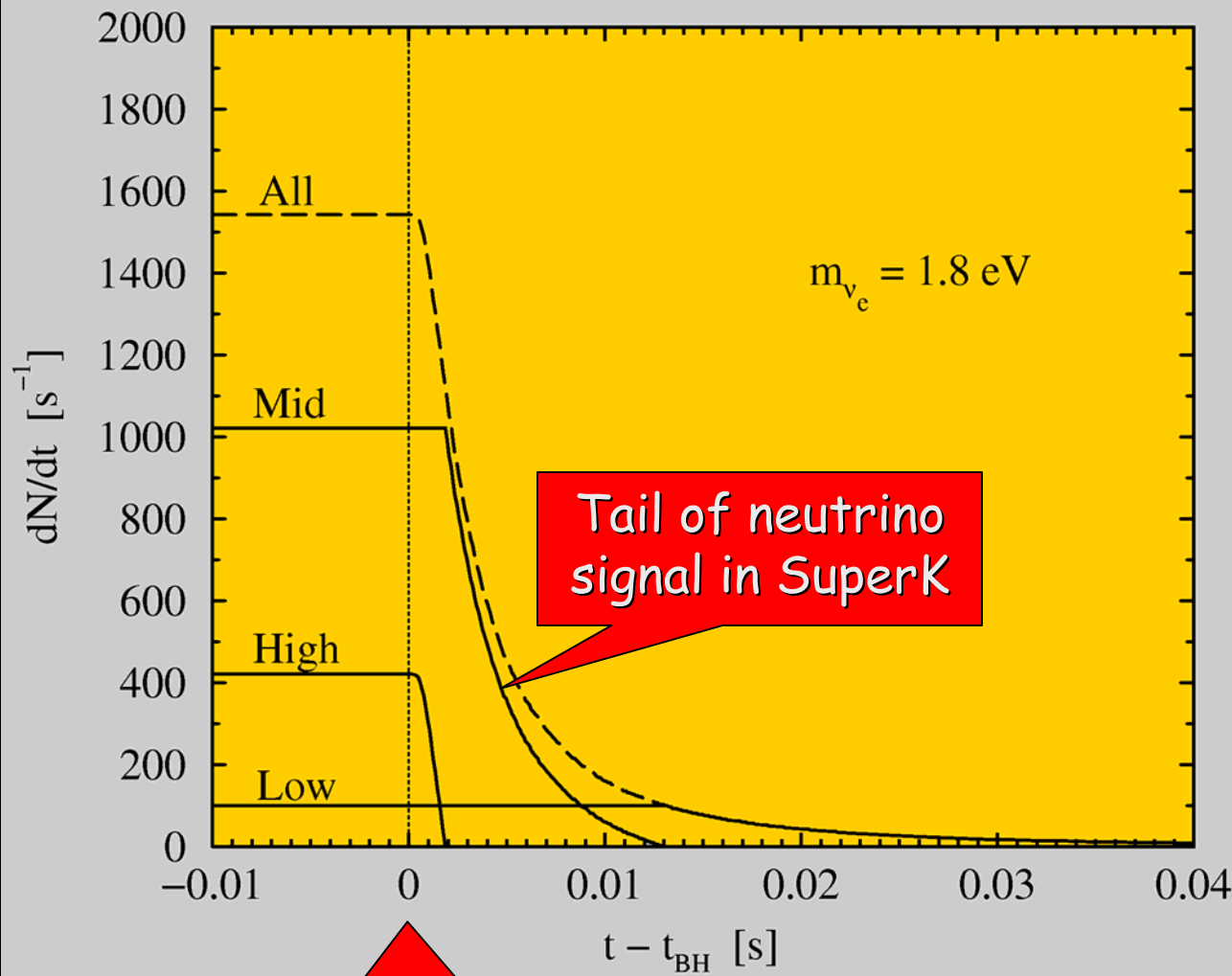
Use correlation between energy & time of arrival as a statistical measure of first N_{cut} events in Monte Carlo simulation of SN signal in SuperK

For SN at 10 kpc sensitive to approx. 3 eV

T.Totani,
PRL 80 (1998) 2040

SN Neutrino Mass from Early Black Hole Formation

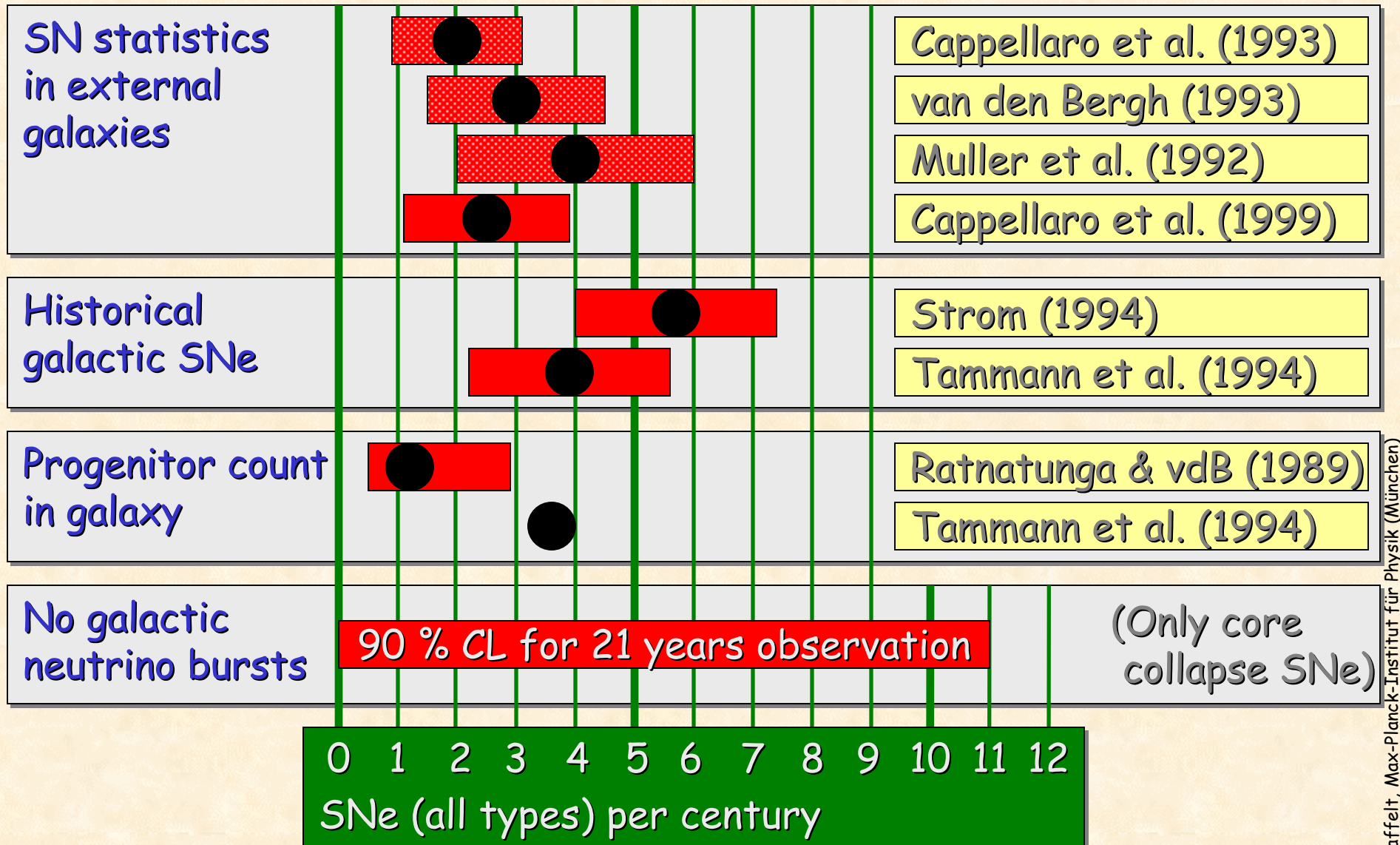
Beacom, Boyd &
Mezzacappa,
hep-ph/0006015



Black Hole Formation

Super-Kamiokande
sensitivity
 $m_{\nu} \sim 1.8 \text{ eV}$

Estimates of the Galactic SN Rate



The Future: A Megatonne Detector?

Megatonne detector motivated by

- Long baseline neutrino oscillations
 - Proton decay
 - Atmospheric neutrinos
 - Solar neutrinos
 - Supernova neutrinos
- ($\sim 10^5$ events for SN at 10 kpc)

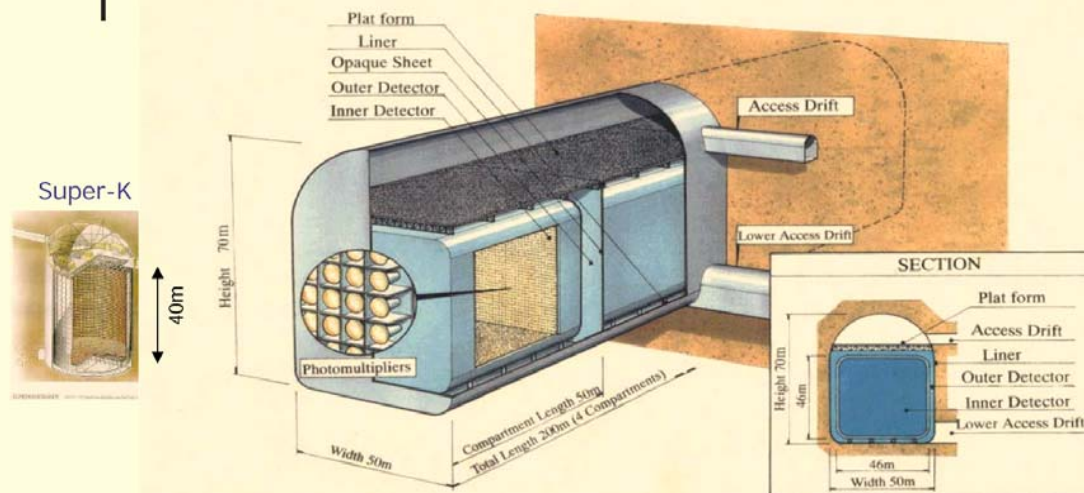
1. Overview of the experiment

(expect to start in 2007)



Possible Design of Hyper-Kamiokande

May.-2001 JHF-Kamioka ν workshop

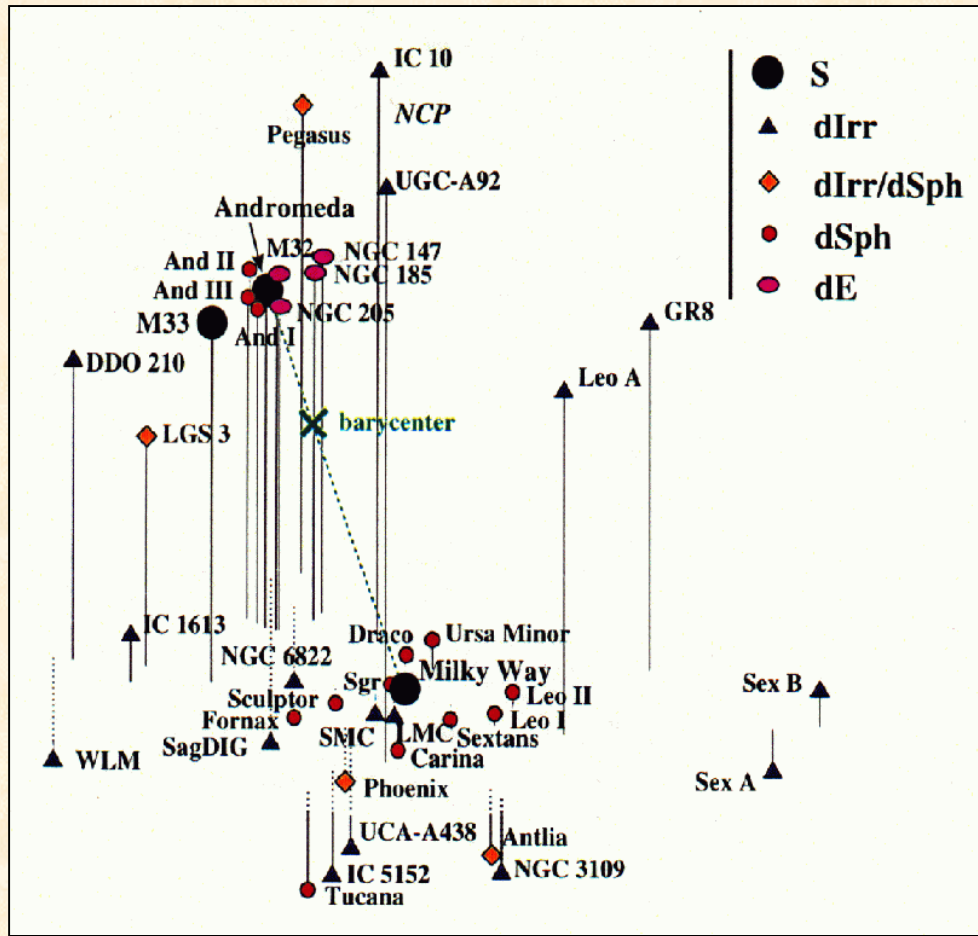


1 Mton fiducial volume: Total Length 800m (16 Compartments)

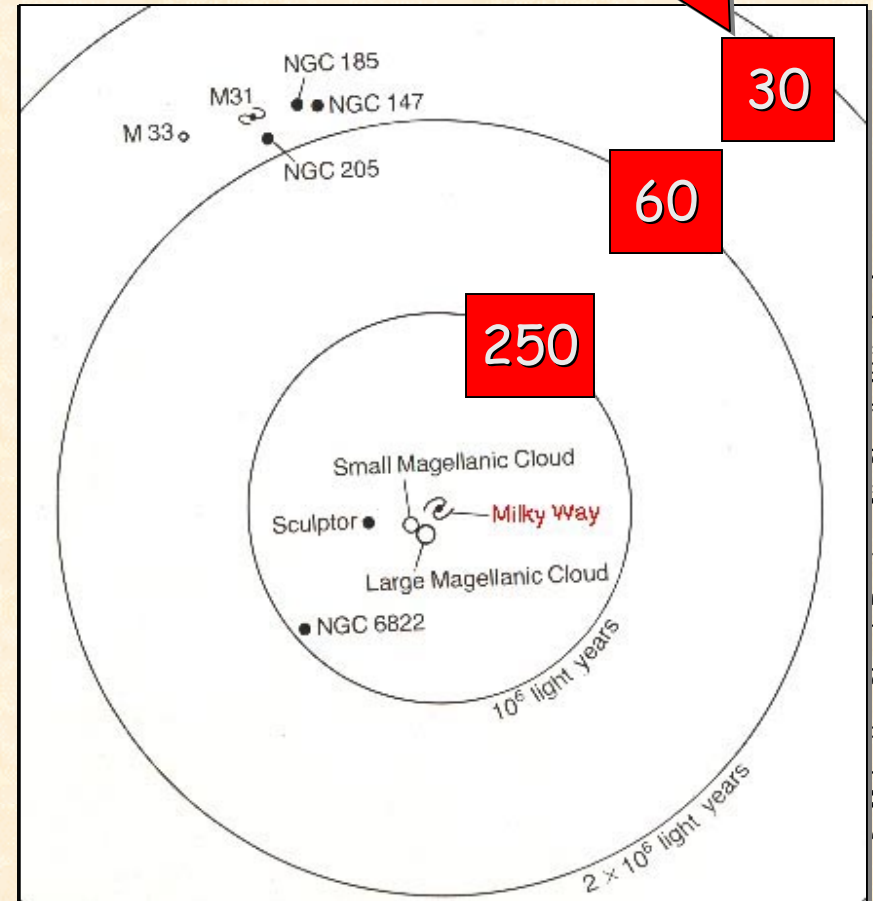
(STRAIGHT TYPE)

- ## Similar discussions in
- USA (UNO project)
 - Europe (Frejus Tunnel)

Local Group of Galaxies



Events in a detector with
30 x SuperK fiducial volume
(8000 events in SK at 10 kpc)



Brightest Members of Local Group

	Type	Lum	D kpc	Neutrino events	SN (all types) per century
Milky Way	S(B)bc	1	8.5	330,000	1-6
LMC	Ir	0.11	50	9,600	0.1 / 0.23 / 0.49
SMC	Ir	0.030	60	6,600	0.065 / 0.12
NGC 6822	Ir	0.011	500	96	0.04
IC 10 (UGC 192)	Ir	0.015	660	55	0.082-0.11
NGC 205	Sph	0.016	760	42	
M32 (NGC 221)	E2	0.017	760	42	
Andromeda (M31)	Sb	1.3	760	42	0.9 / 1.21 / 1.25
Triangulum (M33)	Sc	0.16	790	38	0.28 / 0.35 / 0.68

- Luminosity: Visual in units of the Milky Way
- Neutrino events in $30 \times$ SK fiducial vol. (8000 events in SK for SN at 10 kpc)
- Refs. for SN rates in Pavlidou & Fields, Ap. J. 558 (2001) 63.

Neutrino Limits by Intrinsic Signal Dispersion

Time-of-flight delay
of massive neutrinos

$$\Delta t = 2.57 \text{ s} \left(\frac{D}{50 \text{ kpc}} \right) \left(\frac{10 \text{ MeV}}{E_\nu} \right)^2 \left(\frac{m_\nu}{10 \text{ eV}} \right)^2$$

SN 1987A

$E \approx 20 \text{ MeV}$, $\Delta t \approx 10 \text{ s}$, $D \approx 50 \text{ kpc}$

$m_{\nu_e} < 20 \text{ eV}$

Detailed maximum-likelihood analysis similar result
In 1987 competitive with tritium end-point limits

Future
Galactic SN
(Super-K)

$D \approx 50 \text{ kpc}$, Rise-time 0.01 s

Sensitivity approximately

$m_{\nu_e} \sim 3 \text{ eV}$

Future SN
in Andromeda
(Megatonne)

$D \approx 750 \text{ kpc}$, $\Delta t \approx 10 \text{ s}$

Sensitivity approximately

$m_{\nu_e} \sim 1\text{-}2 \text{ eV}$

Neutrino Mass and SNe - Good or Bad News?

Tritium Limits $< 2\text{--}3\text{ eV}$
Future $< 0.3\text{ eV}$
Cosmological LSS similar limits
Oscillations: All mass differences $< 0.1\text{ eV}$

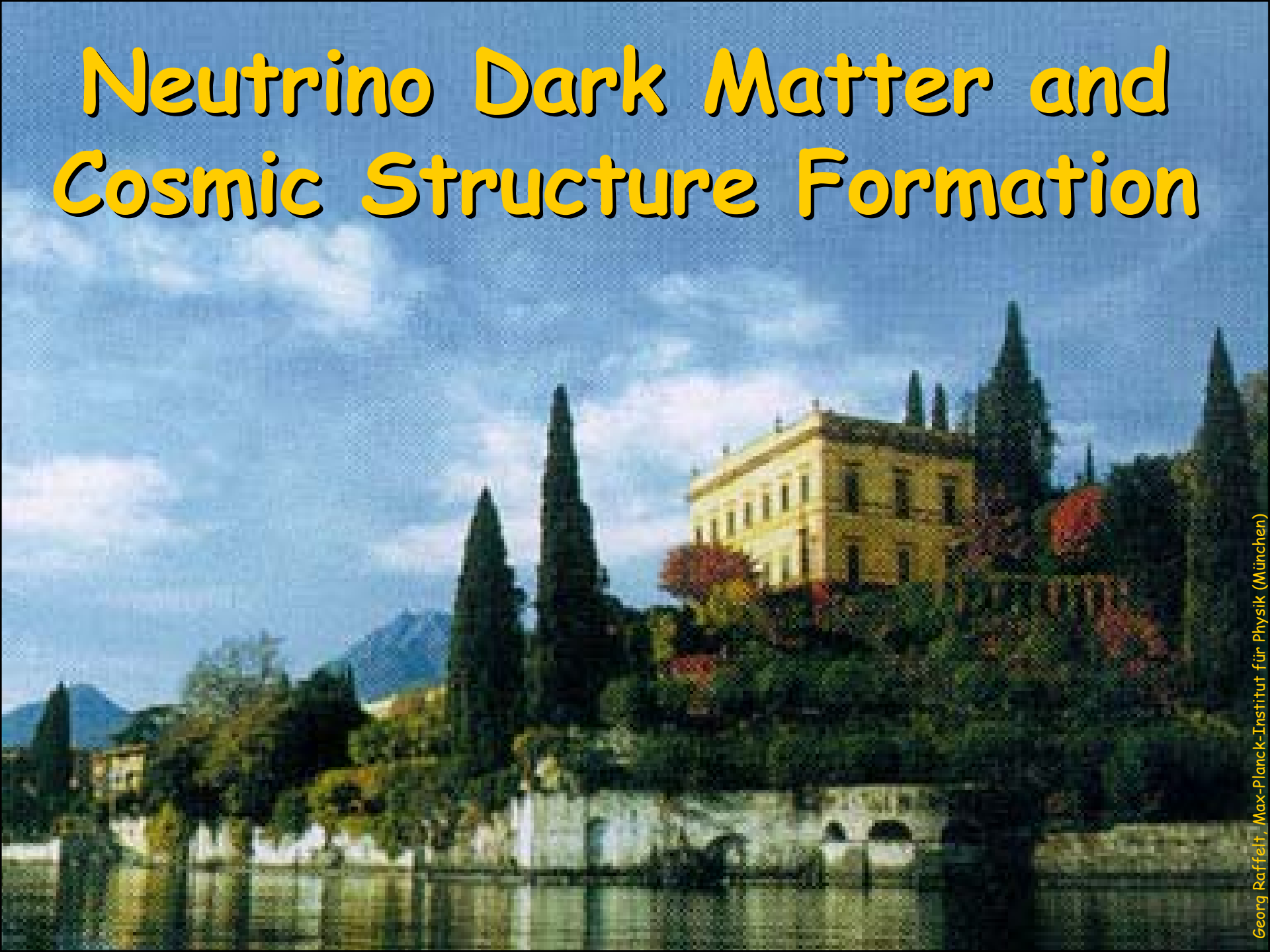
All active flavors
likely to have
sub-eV masses
or less

Supernova time-of-flight dispersion

$$\Delta t = 0.46\text{ms} \left(\frac{D}{10\text{kpc}} \right) \left(\frac{10\text{MeV}}{E_\nu} \right)^2 \left(\frac{m_\nu}{0.3\text{eV}} \right)^2$$

Measure pristine
neutrino light
curve, unspoilt
by dispersion
effects

Neutrino Dark Matter and Cosmic Structure Formation



Cosmological Limit on Neutrino Masses

Cosmic neutrino "sea":

Approximately **112 cm⁻³** neutrinos + antineutrinos per flavor

$$\Omega_\nu h^2 = \sum \frac{m_\nu}{94 \text{ eV}} < 0.4$$

$$m_\nu < 40 \text{ eV} \quad \text{For all stable flavors}$$

REST MASS OF MUONIC NEUTRINO AND COSMOLOGY

S. S. Gershtein and Ya. B. Zel'dovich

Submitted 4 June 1966

ZhETF Pis'ma 4, No. 5, 174-177, 1 September 1966

A classic paper:

Gershtein & Zeldovich

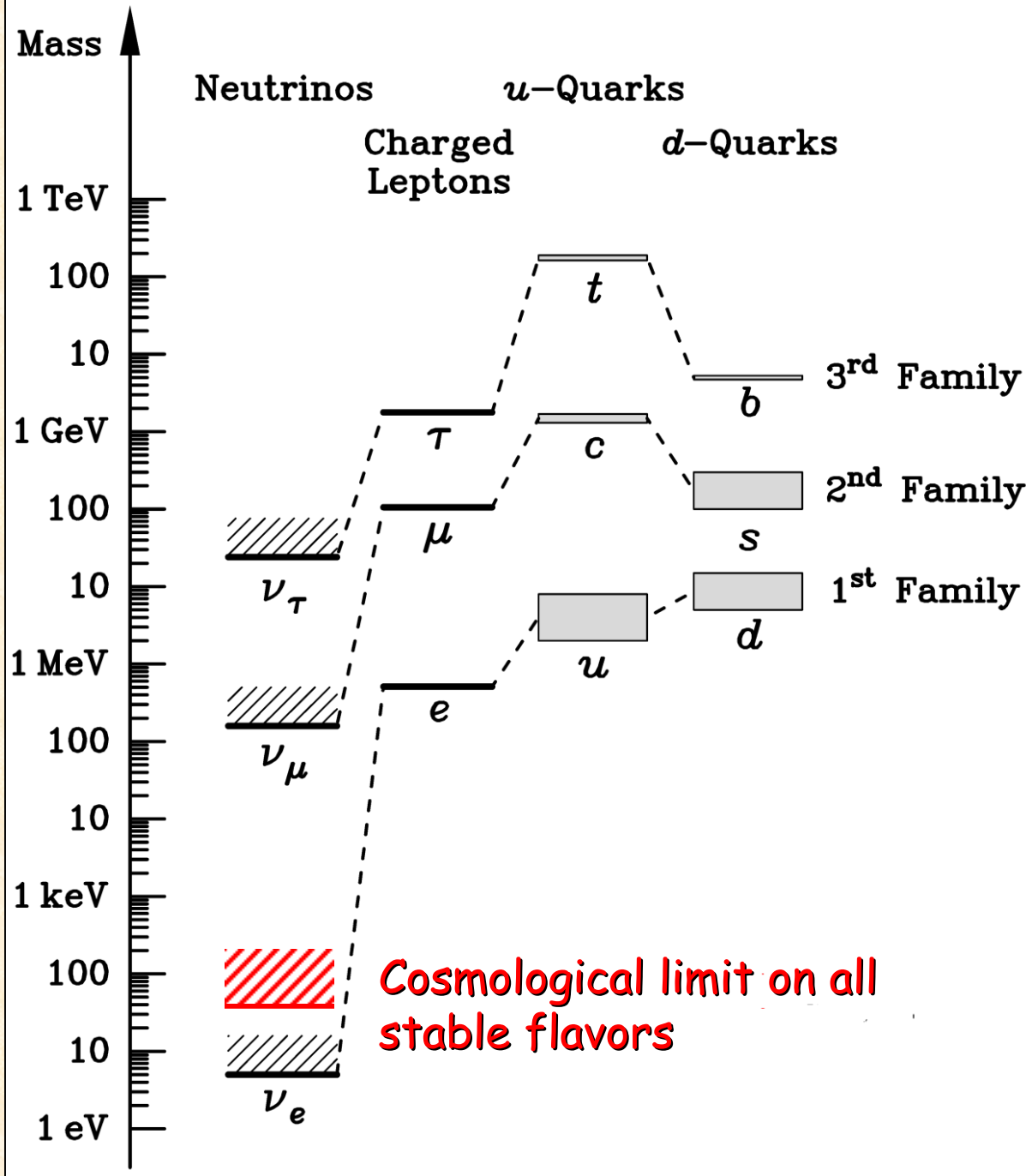
JETP Lett. 4 (1966) 120

Low-accuracy experimental estimates of the rest mass of the neutrino [1] yield $m(\nu_e) < 200 \text{ eV}/c^2$ for the electronic neutrino and $m(\nu_\mu) < 2.5 \times 10^6 \text{ eV}/c^2$ for the muonic neutrino.

Cosmological considerations connected with the hot model of the Universe [2] make it possible to strengthen greatly the second inequality. Just as in the paper by Ya. B. Zel'dovich and Ya. A. Smorodinskii [3], let us consider the gravitational effect of the neutrinos on the dynamics of the expanding Universe. The age of the known astronomical objects is not smaller than 5×10^9 years, and Hubble's constant H is not smaller than 75 km/sec-Mpc = $(13 \times 10^9 \text{ years})^{-1}$. It follows therefore that the density of all types of matter in the Universe is at the present time ¹⁾

$$\rho < 2 \times 10^{-28} \text{ g/cm}^3.$$

Fermion Mass Spectrum



GRAVITY OF NEUTRINOS OF NONZERO MASS IN ASTROPHYSICS

R. COWSIK* AND J. MCCLELLAND
 Department of Physics, University of California, Berkeley
 Received 1972 July 24

ABSTRACT

If neutrinos have a rest mass of a few eV/c^2 , then they would dominate the gravitational dynamics of the large clusters of galaxies and of the Universe. A simple model to understand the virial mass discrepancy in the Coma cluster on this basis is outlined.

Subject headings: cosmology — galaxies, clusters of — neutrinos

The possibility of a finite rest mass for the neutrinos has fascinated astrophysicists (Kuchowicz 1969). A recent discussion of such a possibility has been in the context of the solar-neutrino experiments (Bahcall, Cabibbo, and Yahil 1972). Here we wish to point out some interesting consequences of the gravitational interactions of such neutrinos. These considerations become particularly relevant in the framework of big-bang cosmologies which we assume to be valid in our discussion here.

In the early phase of such a Universe when the temperature was $\sim 1 \text{ MeV}$, several processes of neutrino production (Ruderman 1969) would have led to copious production of neutrinos and antineutrinos (Steigman 1972; Cowsik and McClelland 1972). Conditions of thermal equilibrium allow an easy estimate of their number densities (Landau and Lifshitz 1969):

$$n_{vi} = \frac{1}{\pi^2 \hbar^3} \int_0^\infty \frac{p^2 dp}{\exp [E/kT(z_{eq})] + 1} . \quad (1)$$

Here n_{vi} = number density of neutrinos of the i th kind (notice that in writing this expression we have assumed that both the helicity states are allowed for the neutrinos because of finite rest mass); $E = c(p^2 + m^2 c^2)^{1/2}$; k = Boltzmann's constant; $T(z_{eq}) = T_r(z_{eq}) = T_v(z_{eq}) = T_e(z_{eq}) \dots$ = the common temperature of radiation, neutrinos, electrons, etc., at the latest epoch characterized by redshift z_{eq} when they may be assumed to have been in thermal equilibrium; $kT(z_{eq}) \simeq 1 \text{ MeV}$.

Since the masses of the neutrinos are expected to be small, $kT(z_{eq}) \gg m_{vi} c^2$, in the extreme-relativistic limit equation (1) reduces to

$$n_{vi}(z_{eq}) \simeq 0.183 [T(z_{eq})/\hbar c]^3 . \quad (2)$$

As the Universe expands, only the neutrinos (in contrast to all other known particles) survive annihilation because of extremely low cross-sections (deGraf and Tolhock 1966), and their number density decreases with increasing volume of the Universe, simply as $\sim V(z_{eq})/V(z) = [(1+z)/(1+z_{eq})]^3$. Noting that $(1+z_{eq})/(1+z) = T_r(z_{eq})/T_r(z)$, the number density at the present epoch ($z = 0$) is given by

$$n_{vi}(0) = n_{vi}(z_{eq})/(1+z_{eq})^3 \simeq 0.183 [T_r(0)/\hbar c]^3 \simeq 300 \text{ cm}^{-3} , \quad (3)$$

* On leave from the Tata Institute of Fundamental Research, Bombay, India.

Weakly Interacting Particles as Dark Matter

Almost 30 years ago:

Beginnings of the idea of
weakly interacting particles
(neutrinos) as dark matter.

Massive neutrinos are no
longer thought to be a good
dark matter candidate.

However, the idea of weakly
interacting massive particles
constituting the cosmic dark
matter is now standard.

What is wrong with neutrino dark matter?

→ Galactic Phase Space ("Tremaine-Gunn-Limit")

Maximum mass density of a Fermi gas:

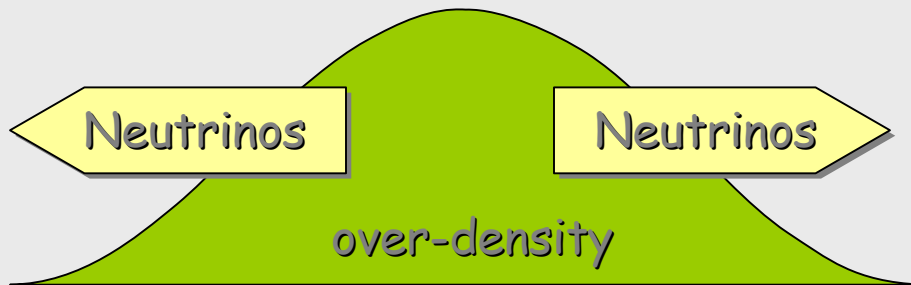
$$\rho_{\max} = m_\nu n_{\max} = m_\nu p_{\max}^3 / 3\pi^2 = m_\nu (m_\nu v_{\text{escape}})^3 / 3\pi^2$$

$$m_\nu > (20 - 40) \text{ eV}$$

Spiral galaxies, more restrictive
from dwarf galaxies: $m_\nu > 100\text{-}200 \text{ eV}$

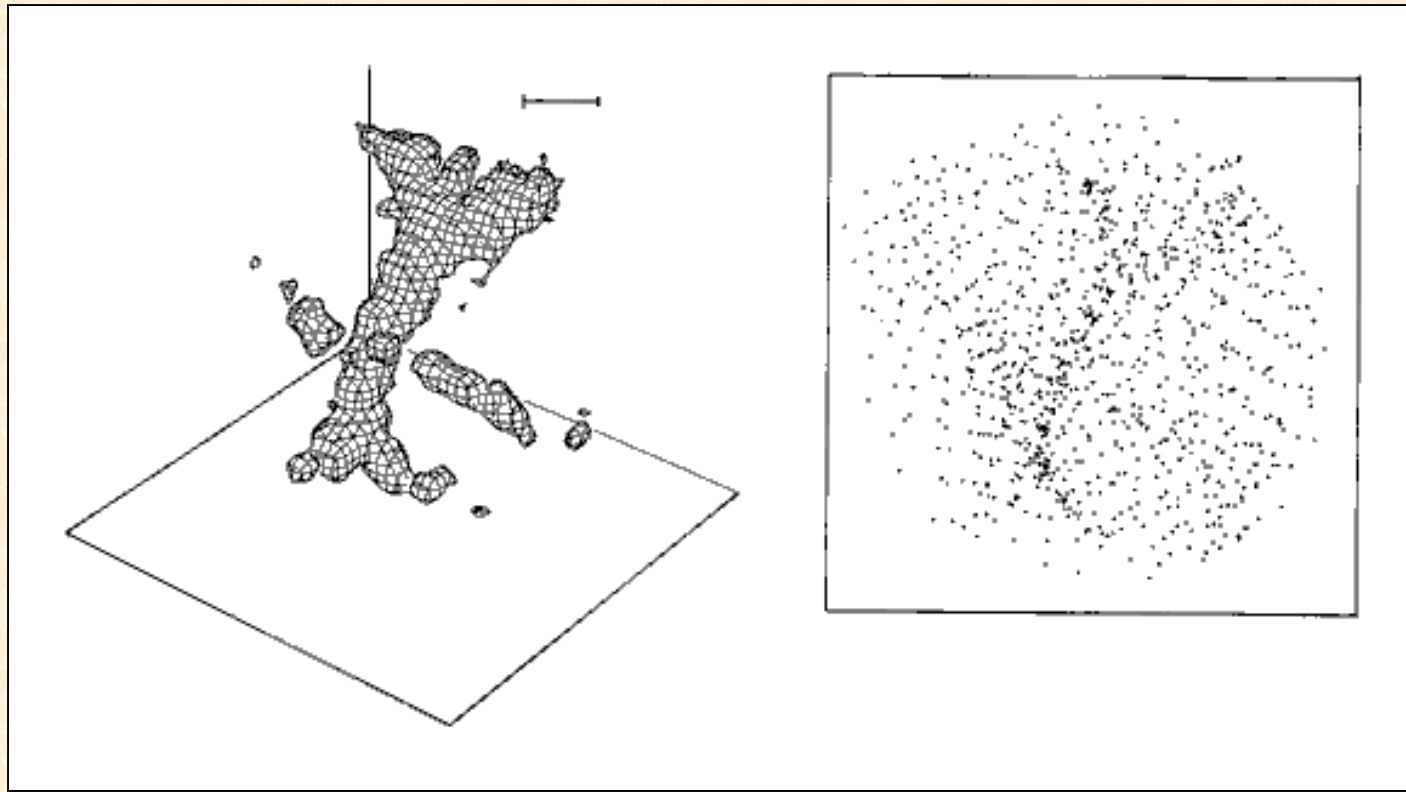
→ Neutrino Free Streaming (Collisionless Phase Mixing)

At $T < 1 \text{ MeV}$ neutrino scattering ineffective, stream freely until nonrelativistic, wash out density contrasts on small scales.



Neutrinos are
"Hot Dark Matter",
excluded by
structure formation

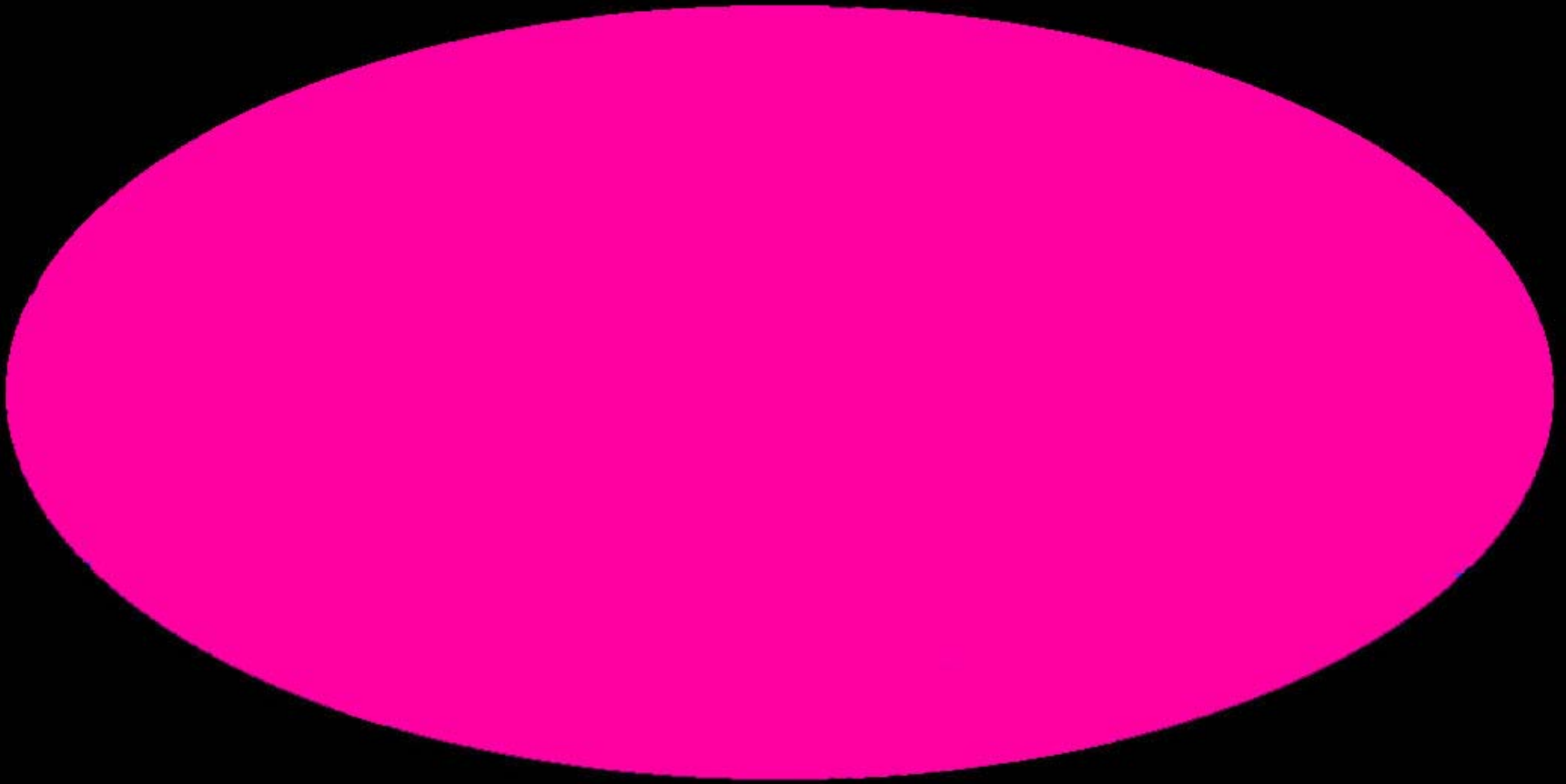
Hot dark matter ruled out in 1983



1000 particle simulation by Frenk, White & Davis, ApJ 271 (1983) 417

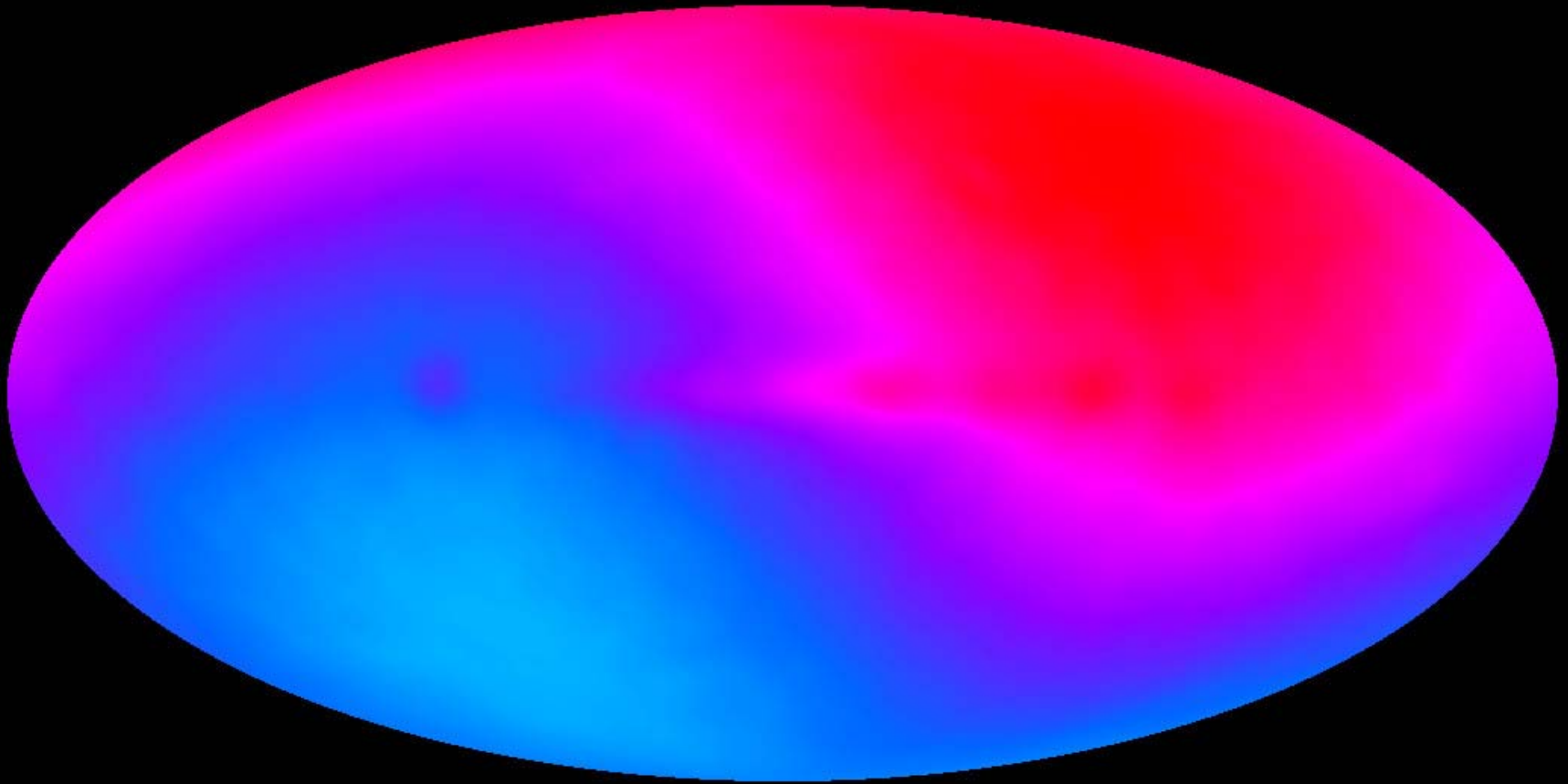
The coherence length of the neutrino distribution ... is too large to be consistent with the observed clustering scale of galaxies ...
The conventional neutrino-dominated picture appears to be ruled out.
[White, Frenk & Davis, ApJ 274 (1983) L1]

COBE Sky Map of the CMBR Temperature



$T = 2.728 \text{ K}$ (uniform on the sky)

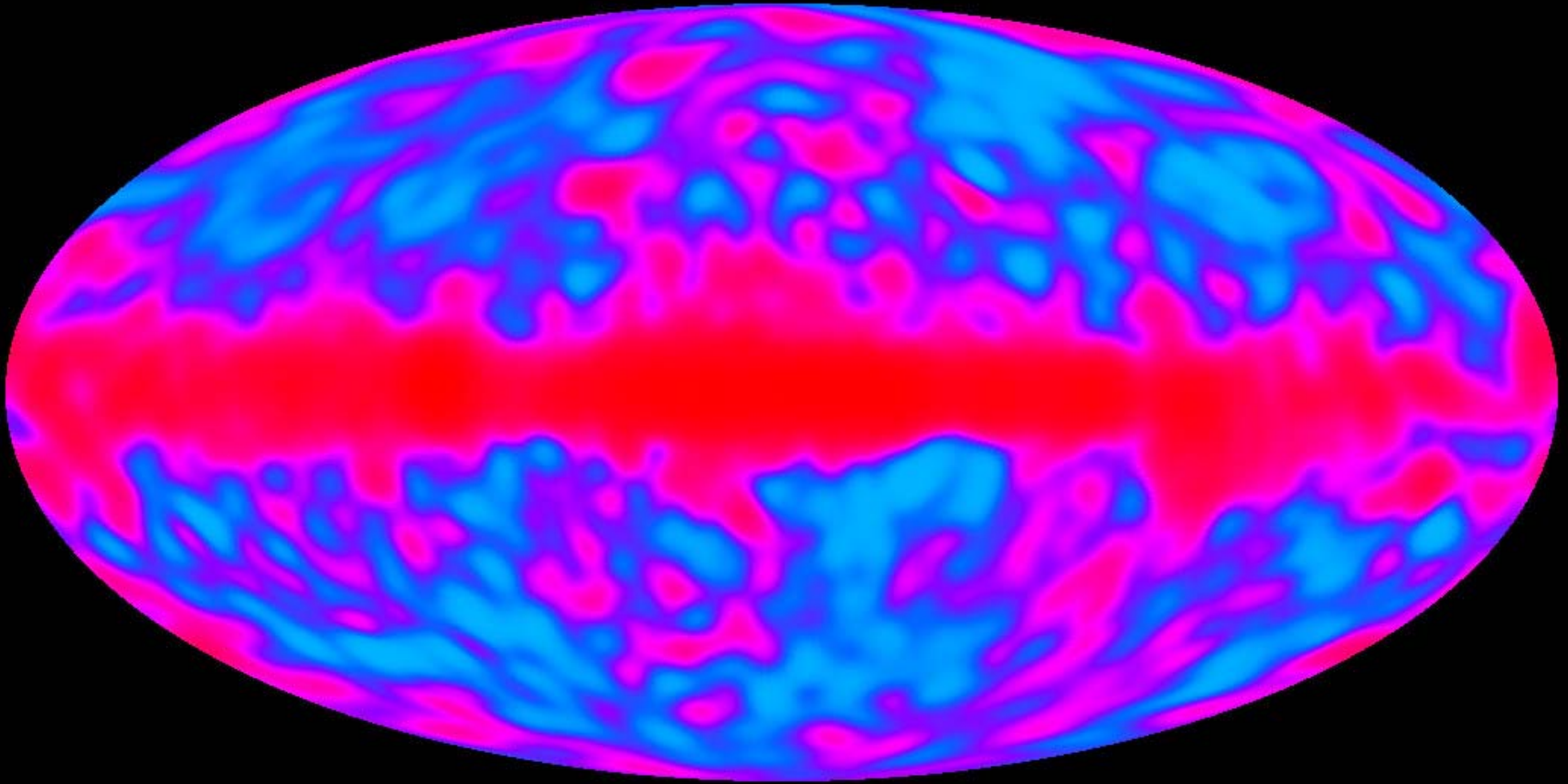
COBE Sky Map of the CMBR Temperature



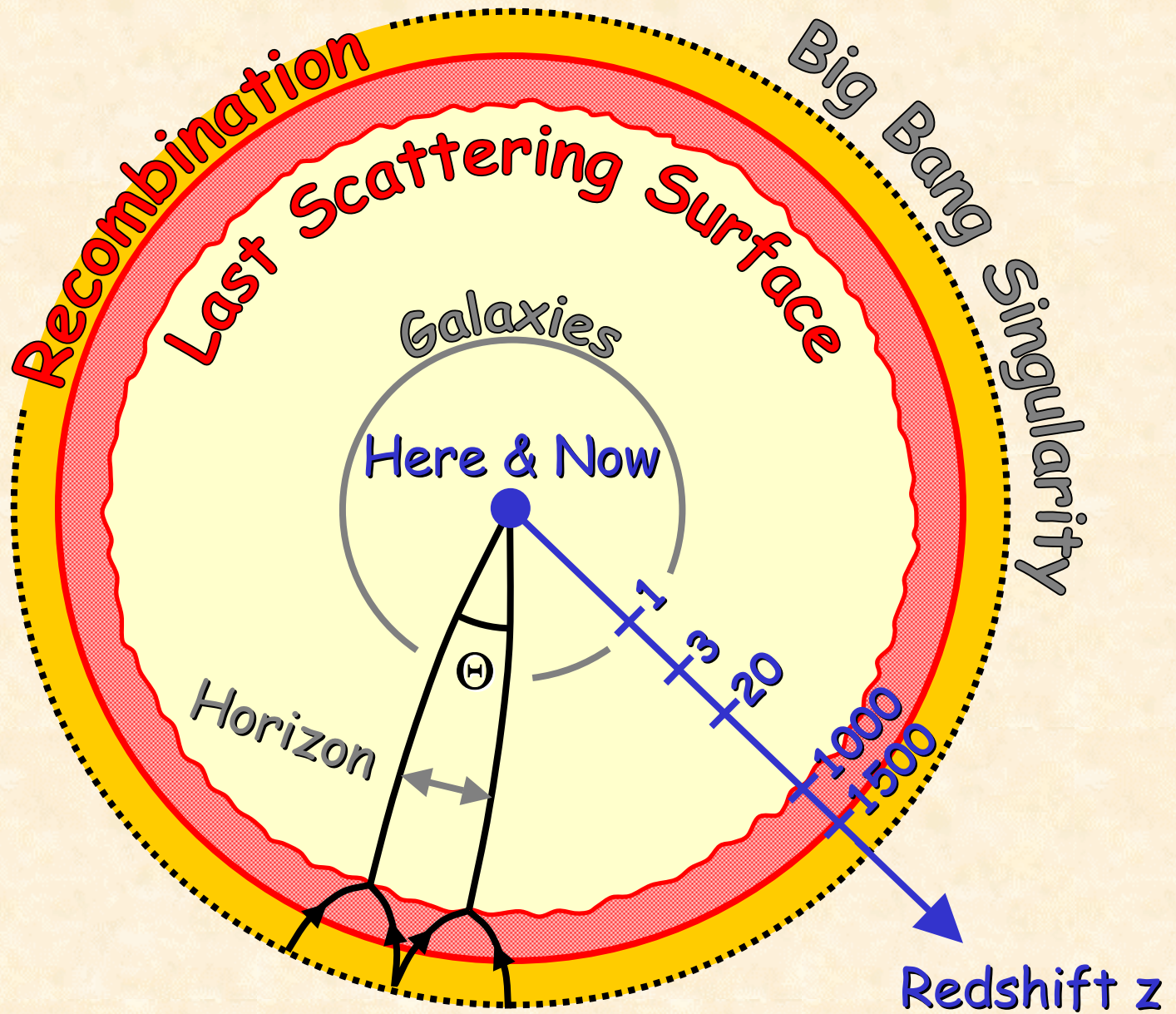
Dynamical range $\Delta T = 3.353 \text{ mK}$

Dipole temperature distribution from Doppler effect due to our motion relative to the cosmic frame

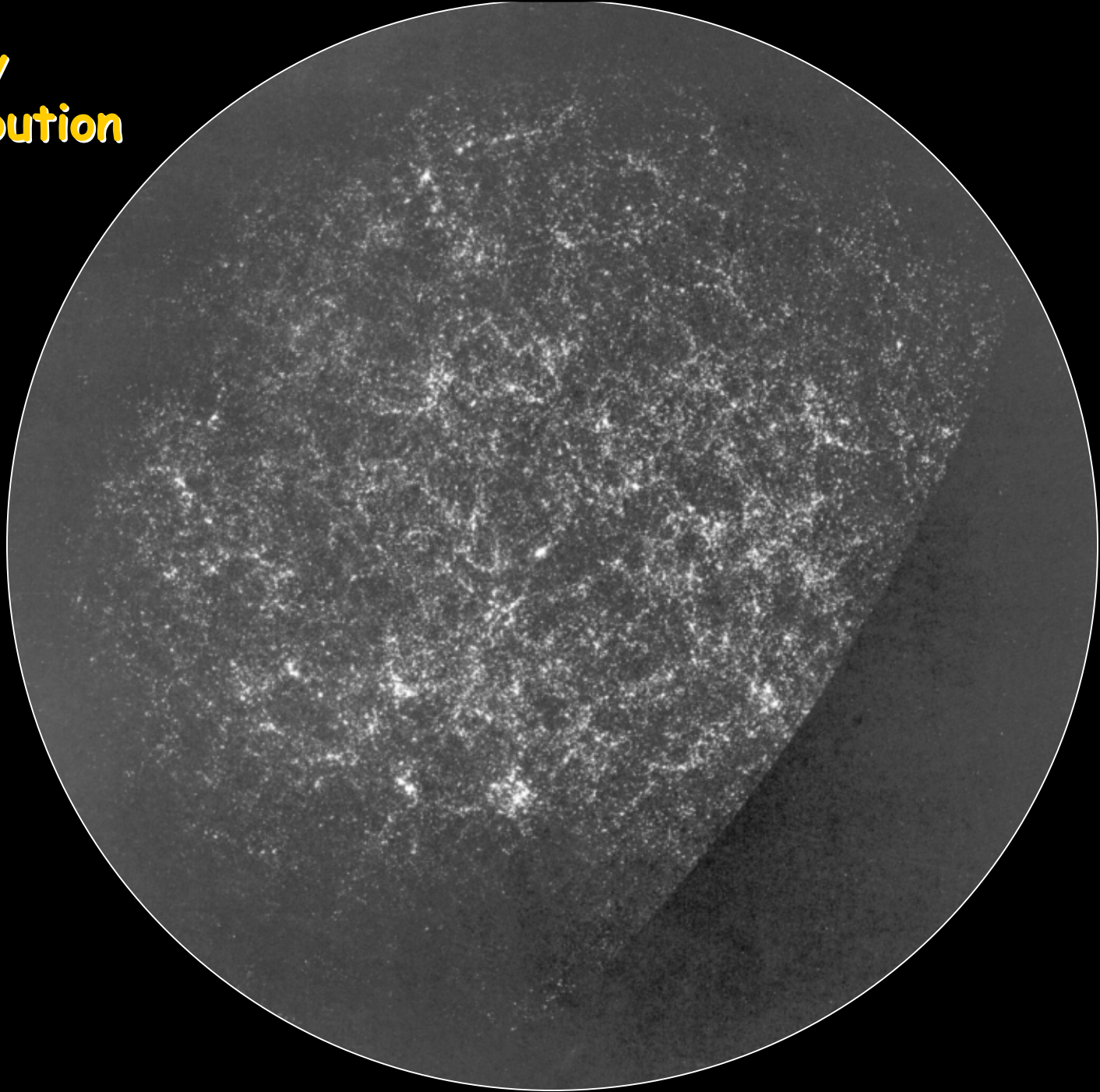
COBE Sky Map of the CMBR Temperature



Dynamical range $\Delta T = 18 \mu\text{K}$
Primordial temperature fluctuations!



**Galaxy
distribution
in the
sky**



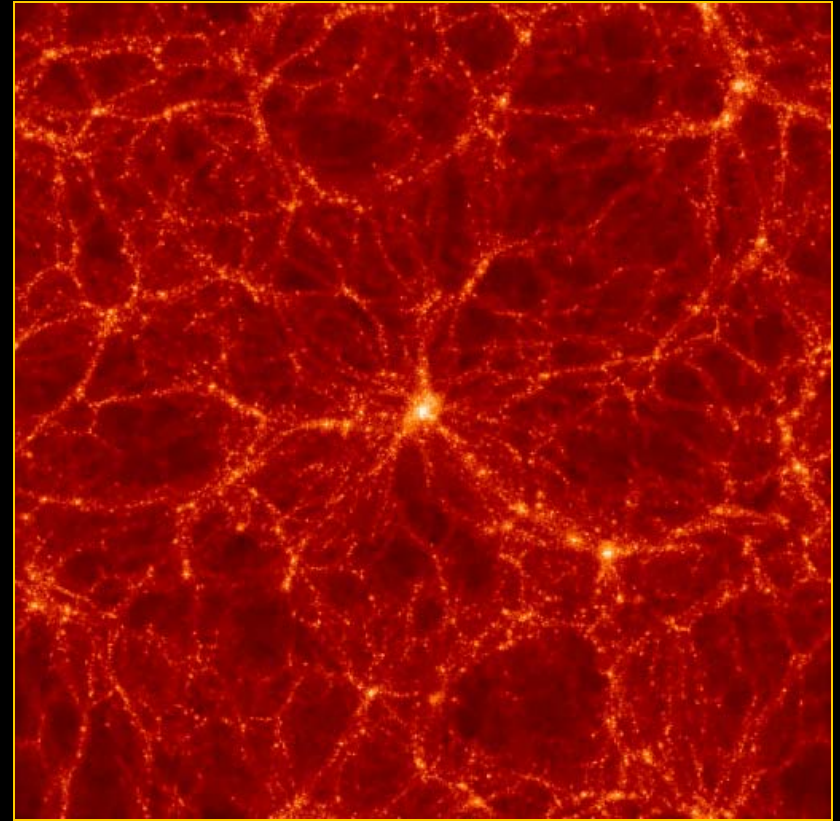
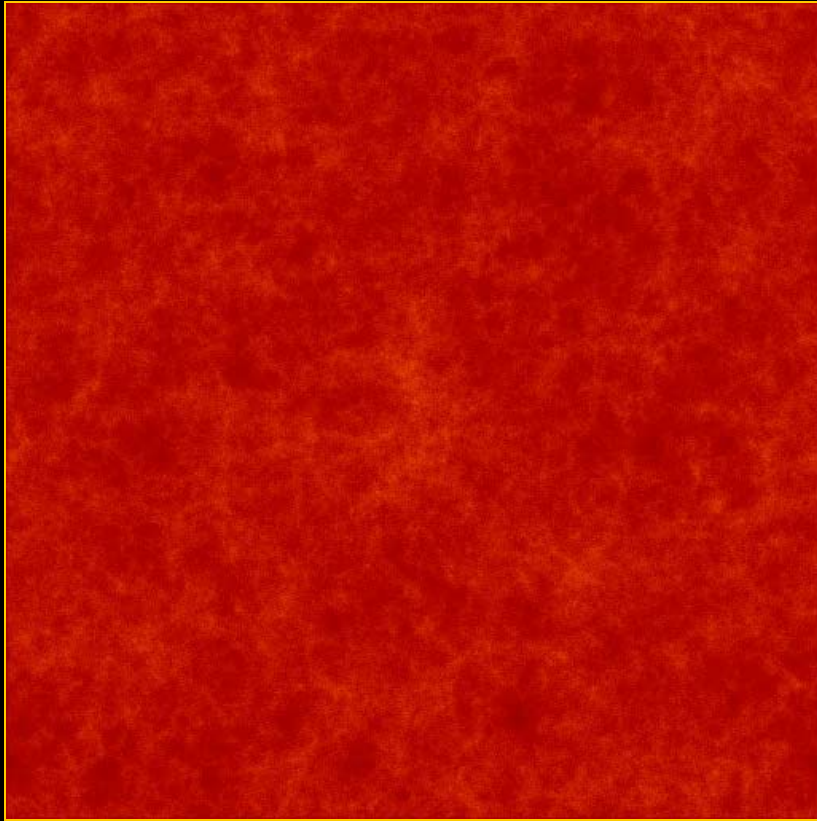
Formation of Structure

[Numerical Simulation Max-Planck-Institut für Astrophysik, Garching]

Smooth



Structured



Power Spectrum of Density Fluctuations

Power Spectrum

Field of Density Fluctuations:

$$\delta(\mathbf{x}) = \frac{\delta\rho(\mathbf{x})}{\bar{\rho}}$$

Fourier Transform:

$$\delta(\mathbf{k}) = \int d^3\mathbf{x} e^{-i\mathbf{k}\cdot\mathbf{x}} \delta(\mathbf{x})$$

Power Spectrum:

$$\langle \delta(\mathbf{k}) \delta(\mathbf{k}') \rangle = (2\pi)^3 \hat{\delta}(\mathbf{k} - \mathbf{k}') P(\mathbf{k})$$

with $\hat{\delta}$ the delta function

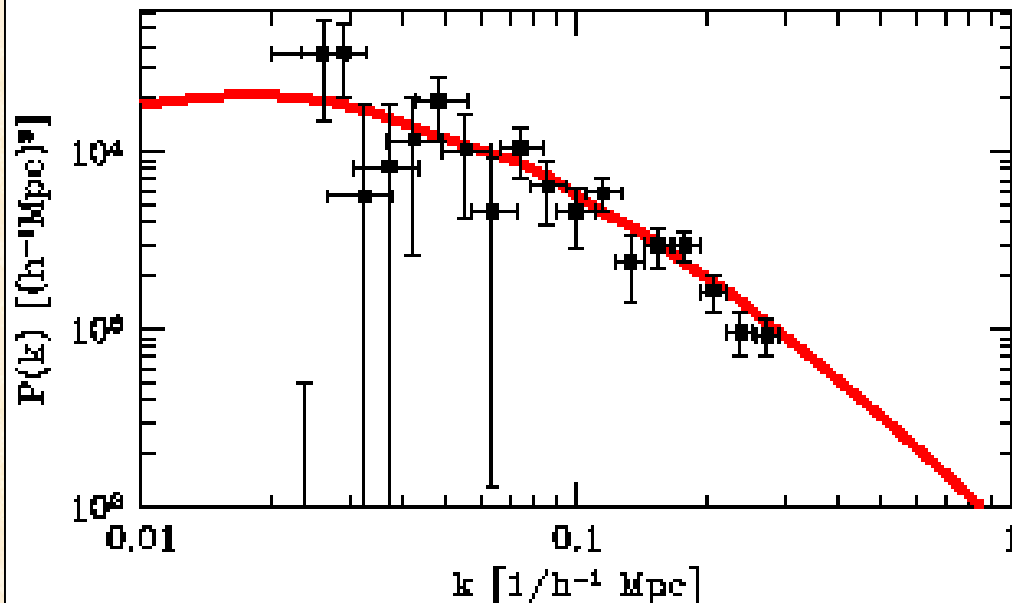
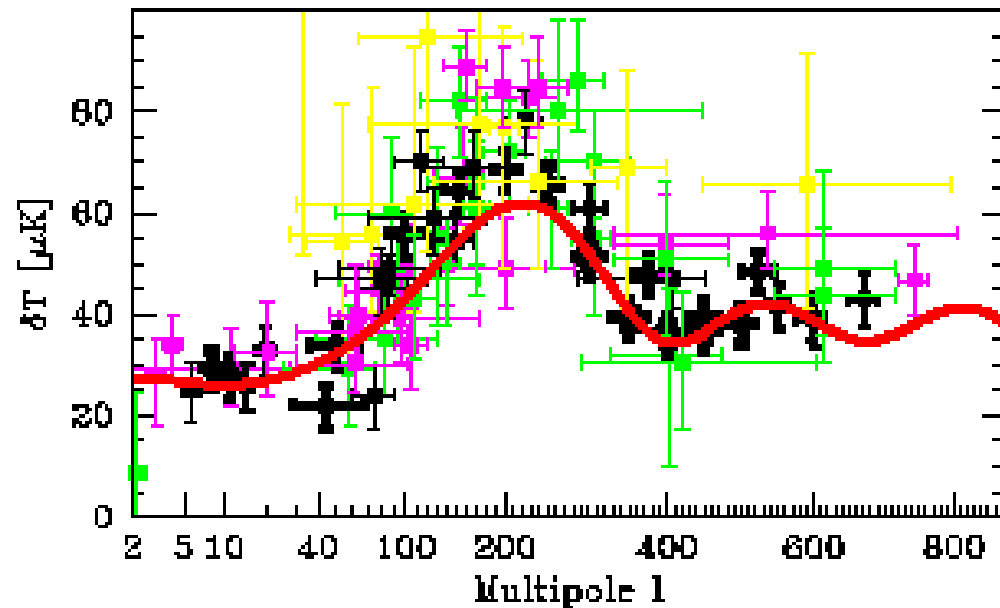
Two-Point Correlation Function:

$$\begin{aligned} \xi(\mathbf{x}) &= \langle \delta(\mathbf{x}_2) \delta(\mathbf{x}_1) \rangle \\ &= \int \frac{d^3\mathbf{k}}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{x}} P(\mathbf{k}) \end{aligned}$$

where $\mathbf{x} = \mathbf{x}_2 - \mathbf{x}_1$.

Gaussian random field
fully characterized
by power spectrum

Fitting the Cosmological Model - Neutrinos



$$\tau = 0.000$$

Ionization parameter

$$\Omega_k = 0.000$$

Curvature

$$\Omega_\Lambda = 0.61$$

Cosmological constant

$$\omega_d = 0.13$$

Dark Matter ($\Omega_M h^2$)

$$\omega_b = 0.020$$

Baryons ($\Omega_B h^2$)

$$f_\nu = 0.000$$

Neutrino fraction of DM

$$n_s = 0.90$$

Scalar power-law index

$$n_t = -1.000$$

Tensor power-law index

$$A_s = 0.44$$

Scalar Amplitude

$$A_t = 0.000$$

Tensor Amplitude

$$b = 1.2$$

Biasing parameter

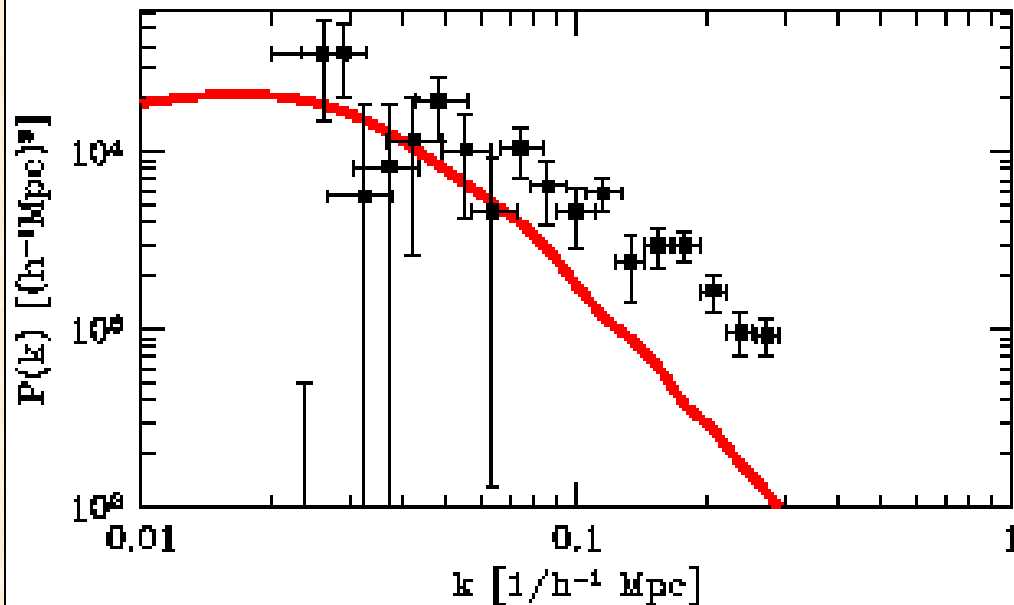
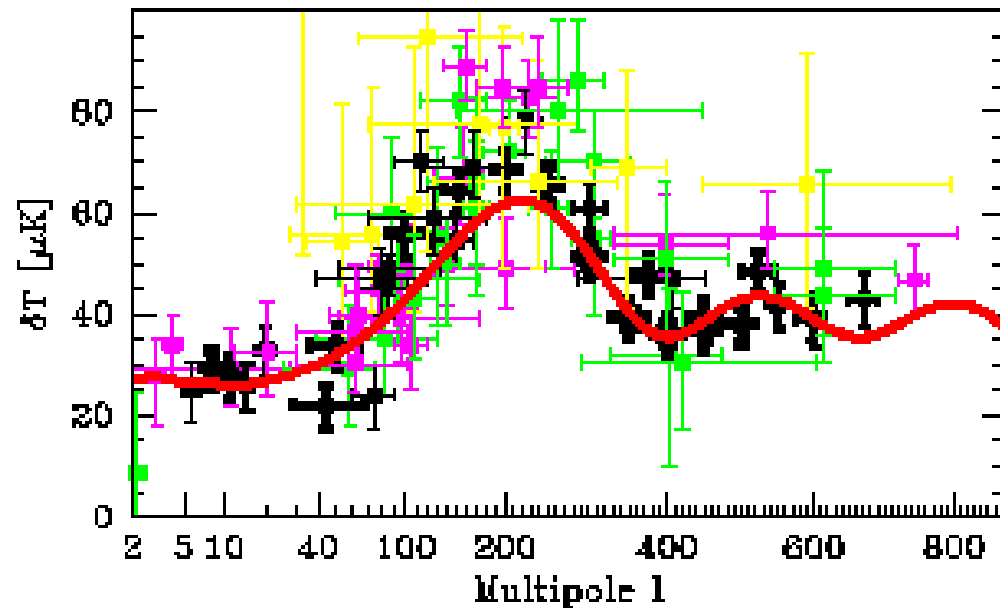
$$h = 0.62$$

Hubble constant

$$\chi^2 = 0.000$$

Max Tegmark,
www.hep.upenn.edu/~max/concordance.html

Fitting the Cosmological Model - Neutrinos



$$\tau = 0.000$$

Ionization parameter

$$\Omega_c = 0.000$$

Curvature

$$\Omega_\Lambda = 0.61$$

Cosmological constant

$$\omega_d = 0.13$$

Dark Matter ($\Omega_M h^2$)

$$\omega_b = 0.020$$

Baryons ($\Omega_B h^2$)

$$f_\nu = 0.36$$

Neutrino fraction of DM

$$n_s = 0.90$$

Scalar power-law index

$$n_t = -1.000$$

Tensor power-law index

$$A_s = 0.44$$

Scalar Amplitude

$$A_t = 0.000$$

Tensor Amplitude

$$b = 1.2$$

Biasing parameter

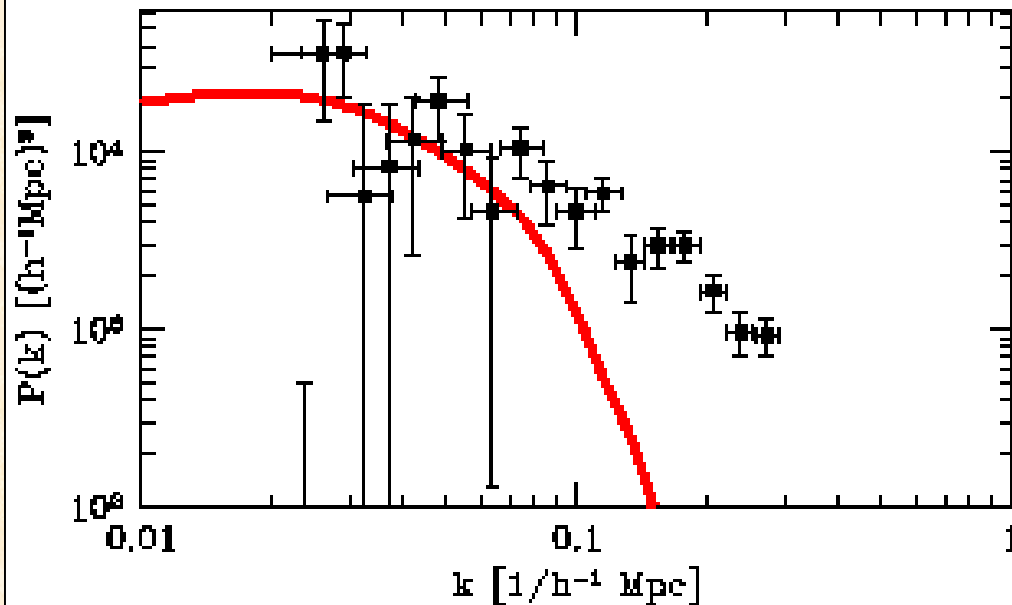
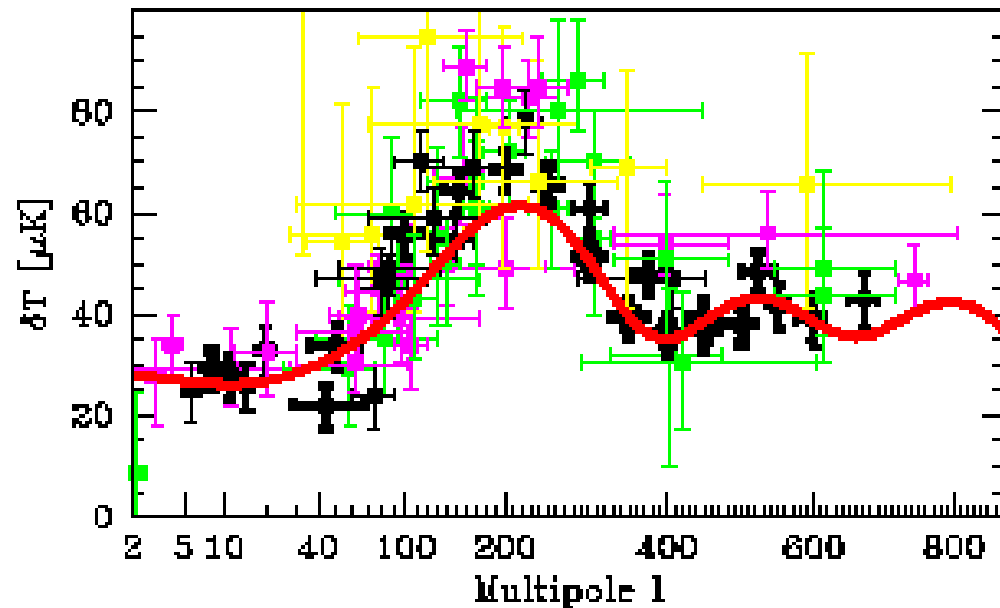
$$h = 0.62$$

Hubble constant

$$\chi^2 = 0.000$$

Max Tegmark,
www.hep.upenn.edu/~max/concordance.html

Fitting the Cosmological Model - Neutrinos



$$\tau = 0.000$$

Ionization parameter

$$\Omega_k = 0.000$$

Curvature

$$\Omega_\Lambda = 0.61$$

Cosmological constant

$$\omega_d = 0.13$$

Dark Matter ($\Omega_M h^2$)

$$\omega_b = 0.020$$

Baryons ($\Omega_B h^2$)

$$f_\nu = 0.99$$

Neutrino fraction of DM

$$n_s = 0.90$$

Scalar power-law index

$$n_t = -1.000$$

Tensor power-law index

$$A_s = 0.44$$

Scalar Amplitude

$$A_t = 0.000$$

Tensor Amplitude

$$b = 1.2$$

Biasing parameter

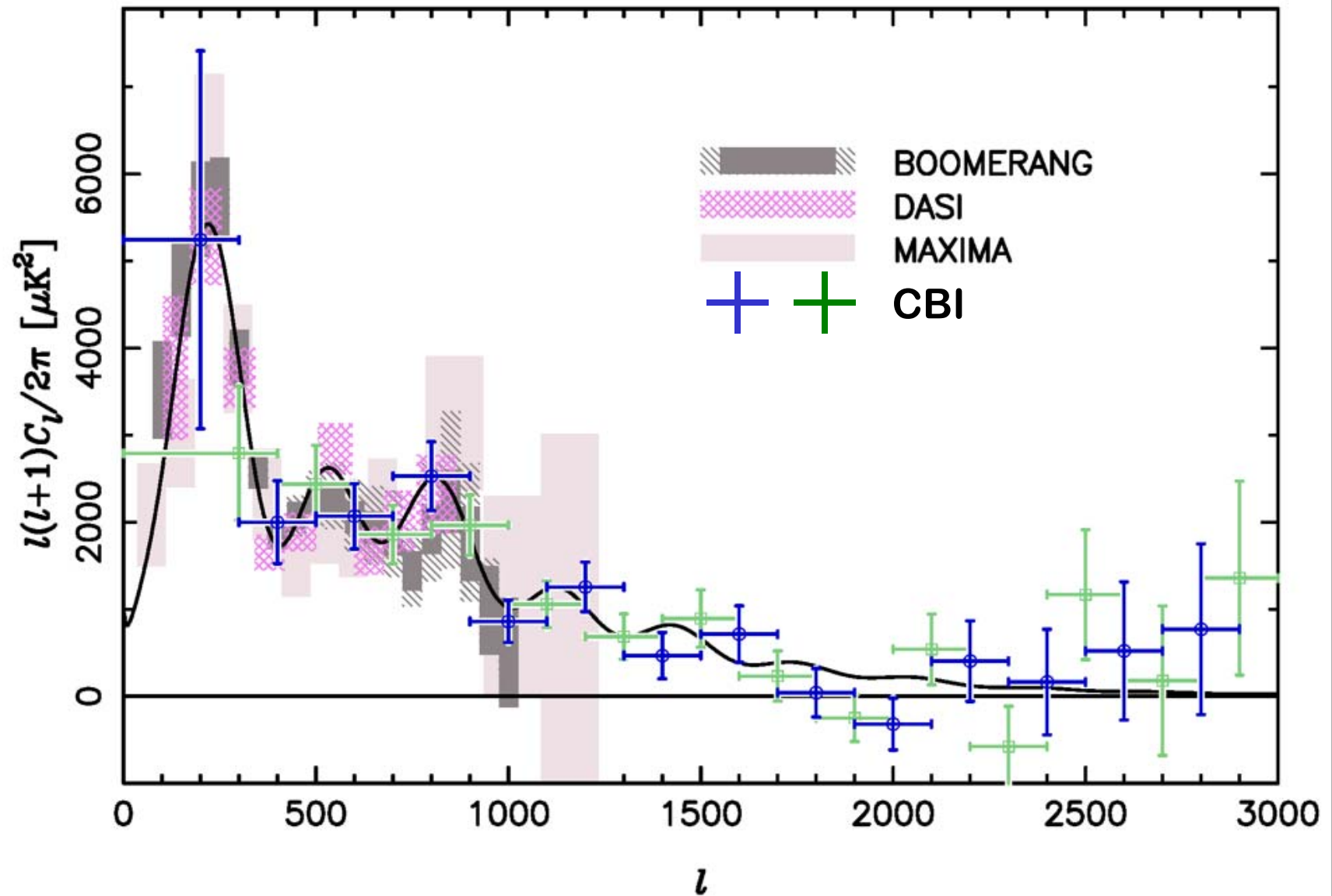
$$h = 0.62$$

Hubble constant

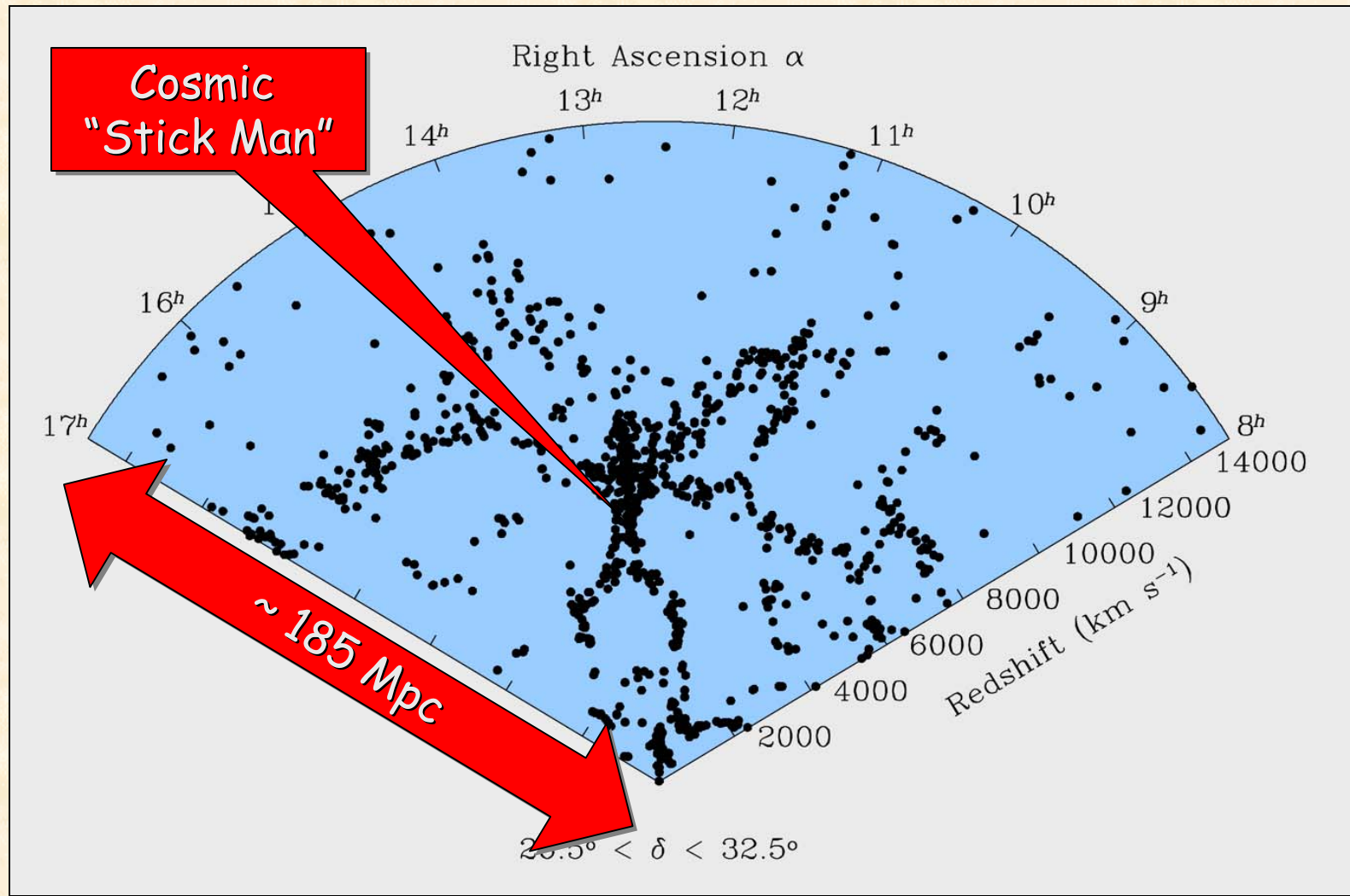
$$\chi^2 = 0.000$$

Max Tegmark,
www.hep.upenn.edu/~max/concordance.html

Multiple Peaks in CMBR Angular Power Spectrum

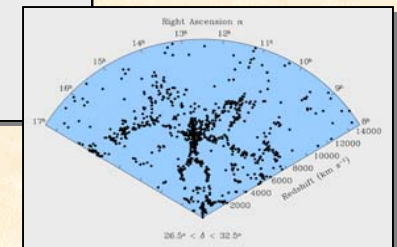
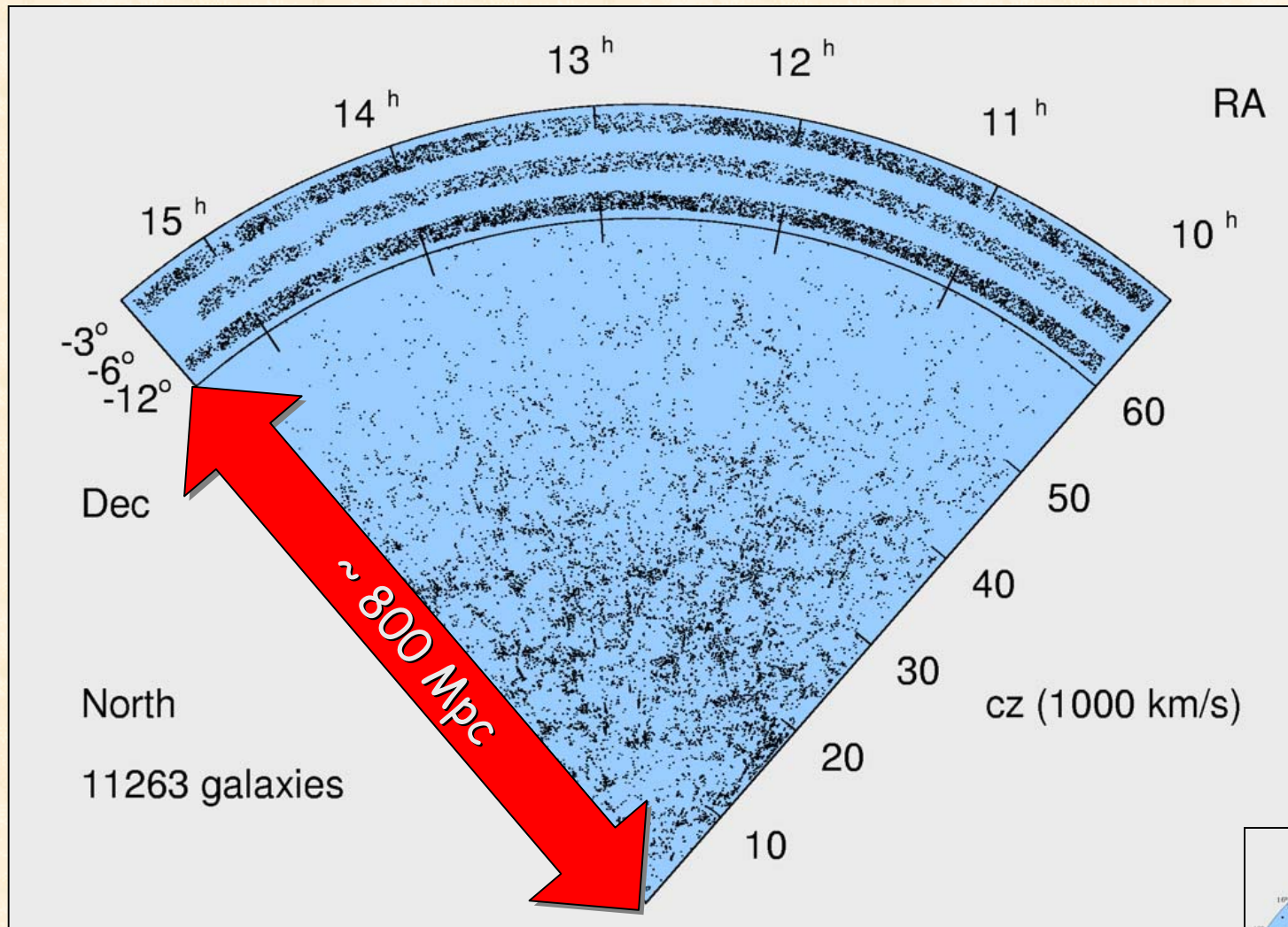


A Slice of the Universe

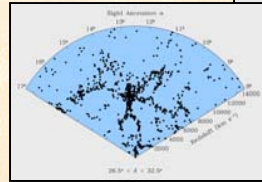
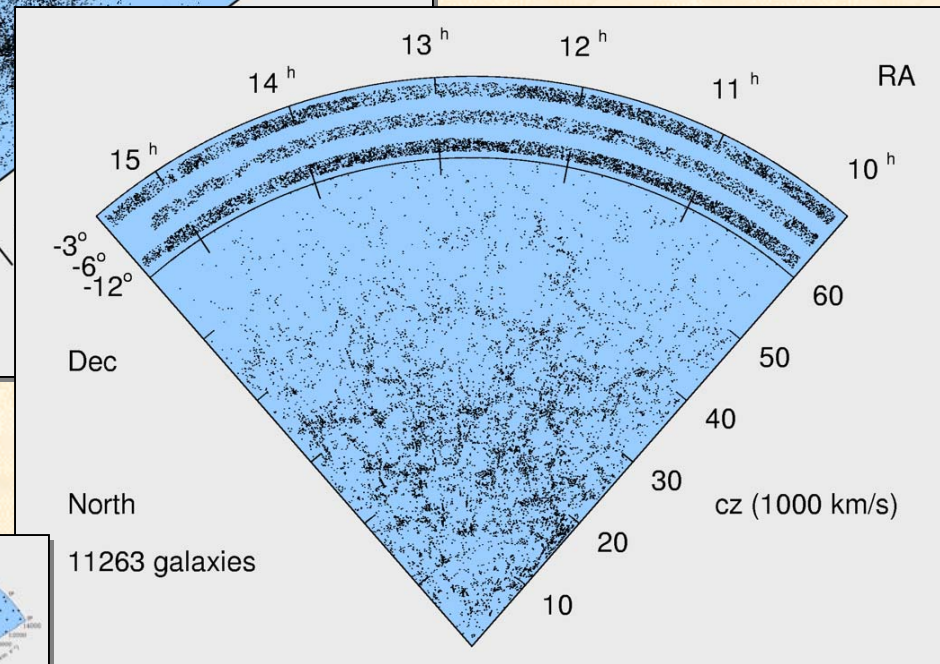
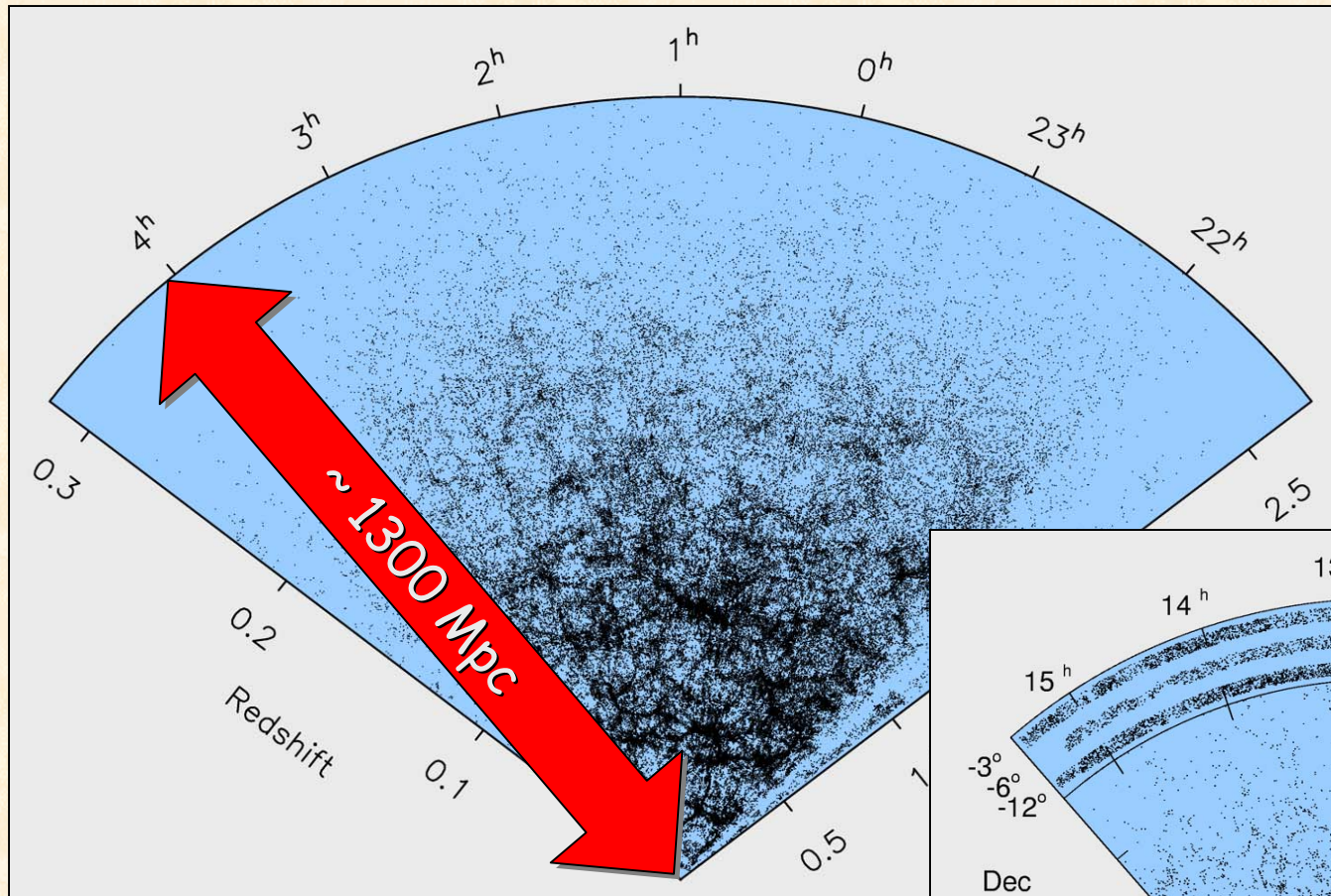


Galaxy distribution from the CfA redshift survey
[ApJ 302 (1986) L1]

Las Campanas Redshift Survey



2dF Galaxy Redshift Survey (15 May 2002)

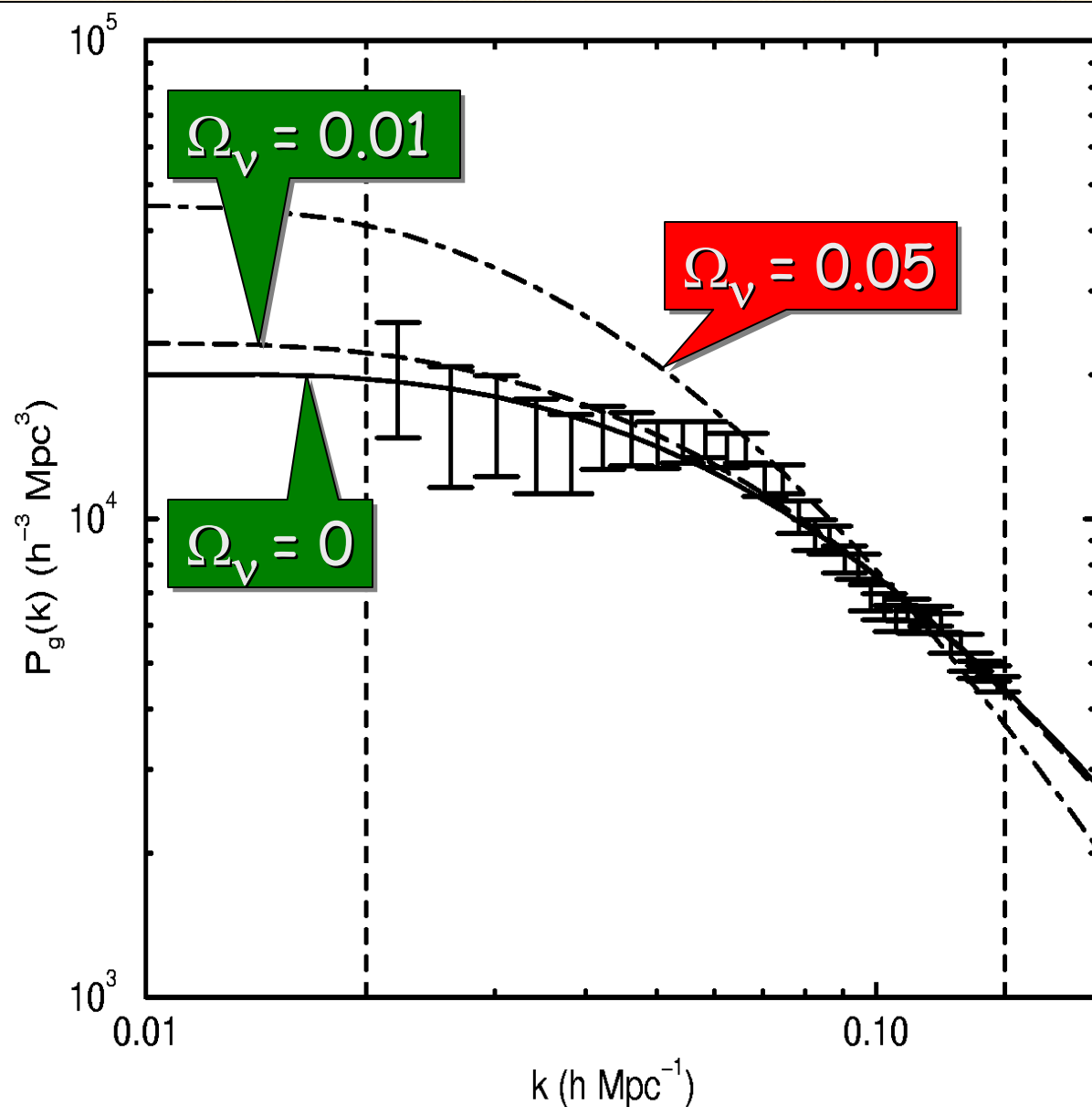


2dF Galaxy Redshift Survey



Animations from 2dFGRS Homepage <http://www.mso.anu.edu.au/2dFGRS/>

Neutrino Mass Limit from 2dF Galaxy Survey



Elgaroy et al.,
astro-ph/0204152

$\Sigma m_\nu < 2.2 \text{ eV}$
at 95% CL

For a similar limit
based on the
2dF data
see Hannestad,
astro-ph/0205223

Neutrino Mass Limits from Cosmic Structure

- Global fits to CMBR and LSS data,
- 3 flavors of equal-mass neutrinos

$$m_\nu < 2.8 \text{ eV at 95\% CL}$$

Novosyadlyj, Durrer & Apunevych,
astro-ph/0011055

$$m_\nu < 1.4 \text{ eV at 95\% CL}$$

Wang, Tegmark & Zaldarriaga,
astro-ph/0105091

$$m_\nu < 0.8 \text{ eV at 95\% CL}$$

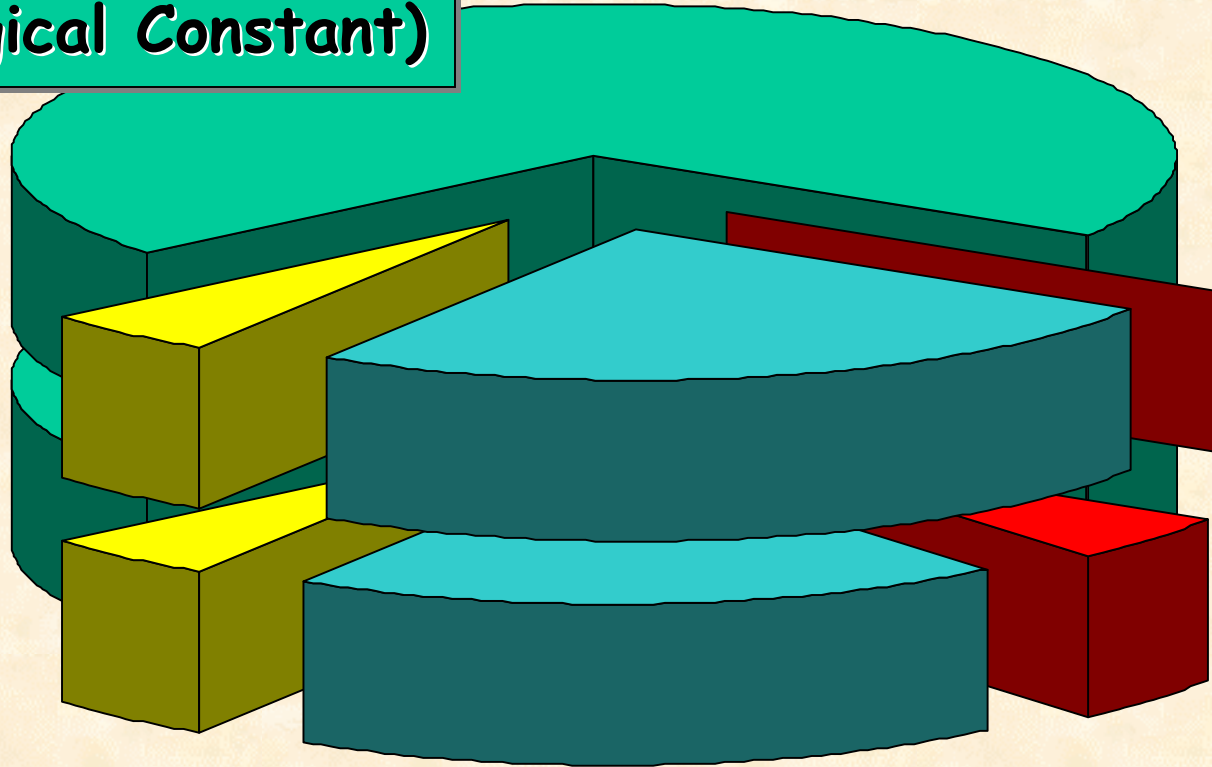
Elgaroy et al., astro-ph/0204152
Hannestad, astro-ph/0205223

For similar older limits see

- Croft, Hu & Davé, PRL 83 (1999) 1092
- Fukugita, Liu & Sugiyama, PRL 84 (2000) 1082
- Gawiser, astro-ph/0005475

Matter Inventory of the Universe

**Dark Energy
(Cosmological Constant)**



**Baryonic Matter
(of this about
10% luminous)**

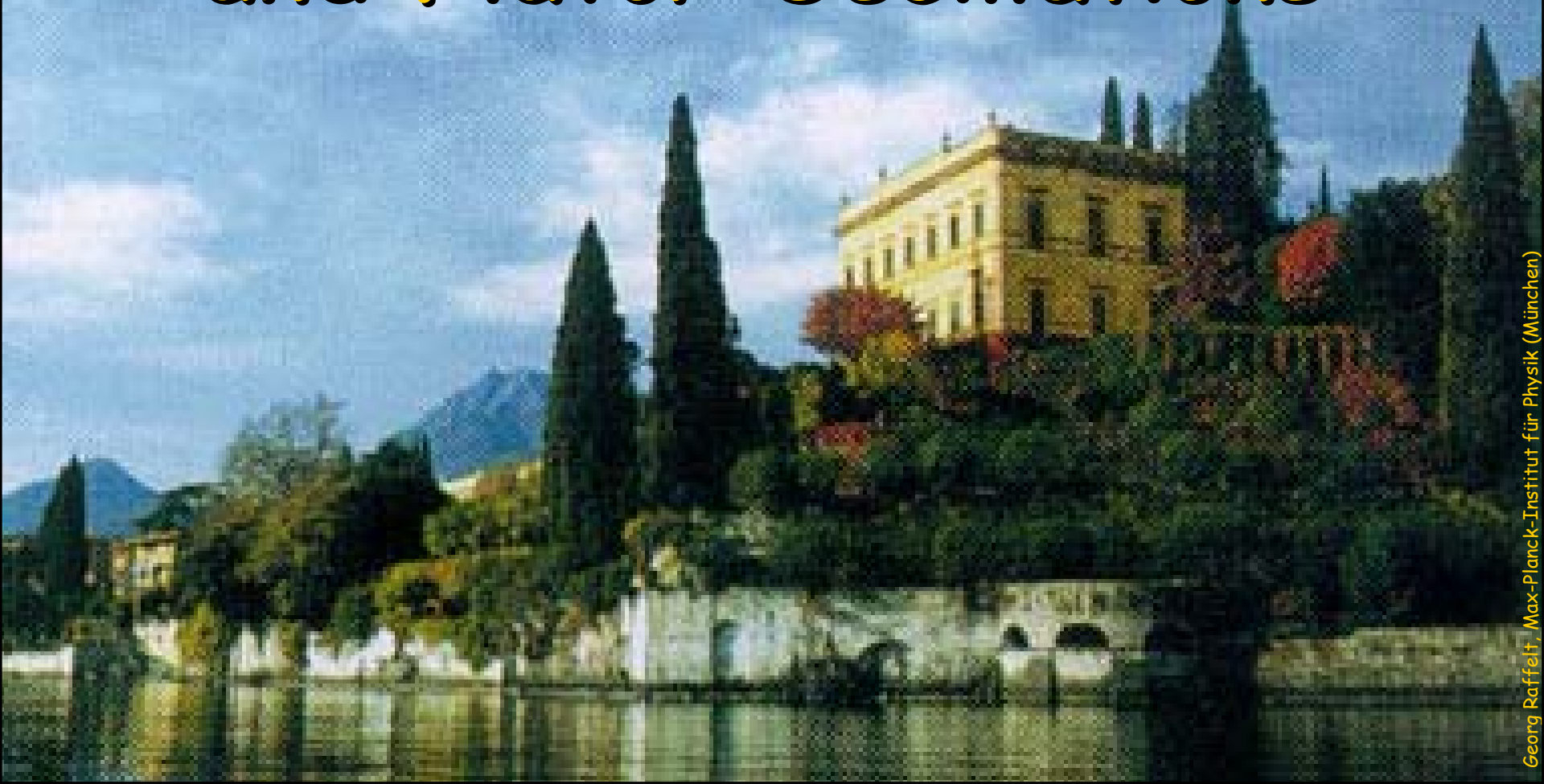
**Dark
Matter**

**Neutrinos
min. 0.1%
max. 6%**

Neutrino Mass Limits and Future Sensitivity

Tritium endpoint	Mainz/Troitsk	2.5 eV
	KATRIN	0.3 eV
Supernova Neutrinos Time-of-flight	SN 1987A	20 eV
	Super-Kamiokande	3 eV
	with black hole	2 eV
	with gravity waves	1 eV
Cosmic structure	2DF Redshift Survey	$m_\nu < 0.8 \text{ eV}$
	Sloan Digital Sky Survey	$m_\nu < 0.3 \text{ eV}$

Neutrino Chemical Potentials, Big-Bang Nucleosynthesis, and Flavor Oscillations



How Many Relic Neutrinos?

Standard thermal population in one flavor: $n_{\nu\bar{\nu}} = \frac{3}{11}n_\gamma \approx 112 \text{ cm}^{-3}$

Additional active neutrinos beyond standard population of ν_e, ν_μ, ν_τ

Additional families

Excluded by Z^0 width: $N_\nu = 3$

Chemical potentials for ν_e, ν_μ, ν_τ

o.k.

Sterile (right-handed) states, populated by $\nu_L \rightarrow \nu_R$

Dirac mass

Too small in eV range

Electromagnetic dipole moments

Excluded by energy loss of globular cluster stars

Right-handed currents

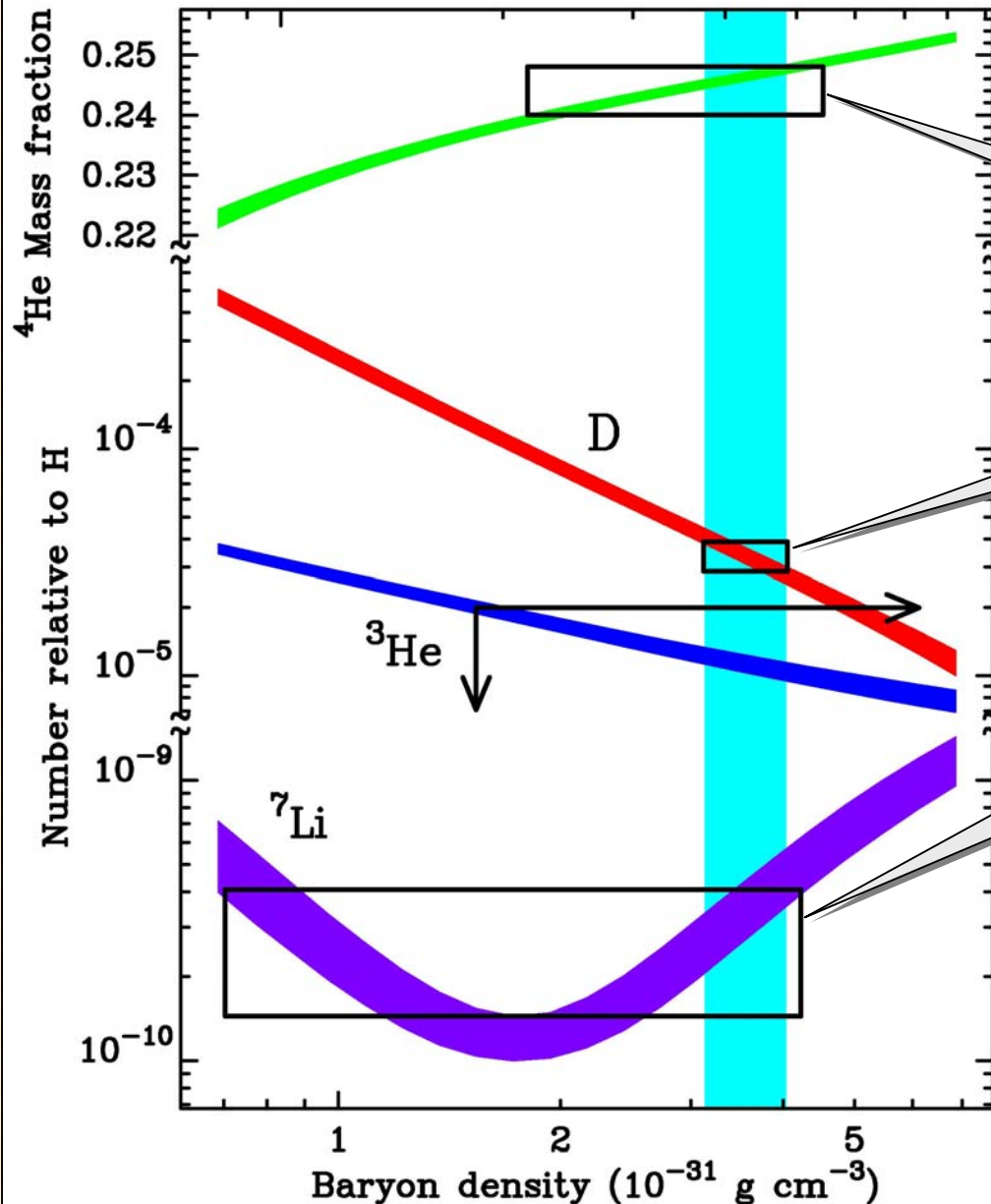
Excluded by energy loss of SN 87A

Oscillations/collisions

o.k., hot/warm DM possible

BBN Concordance

[astro-ph/9903300]



Helium 4

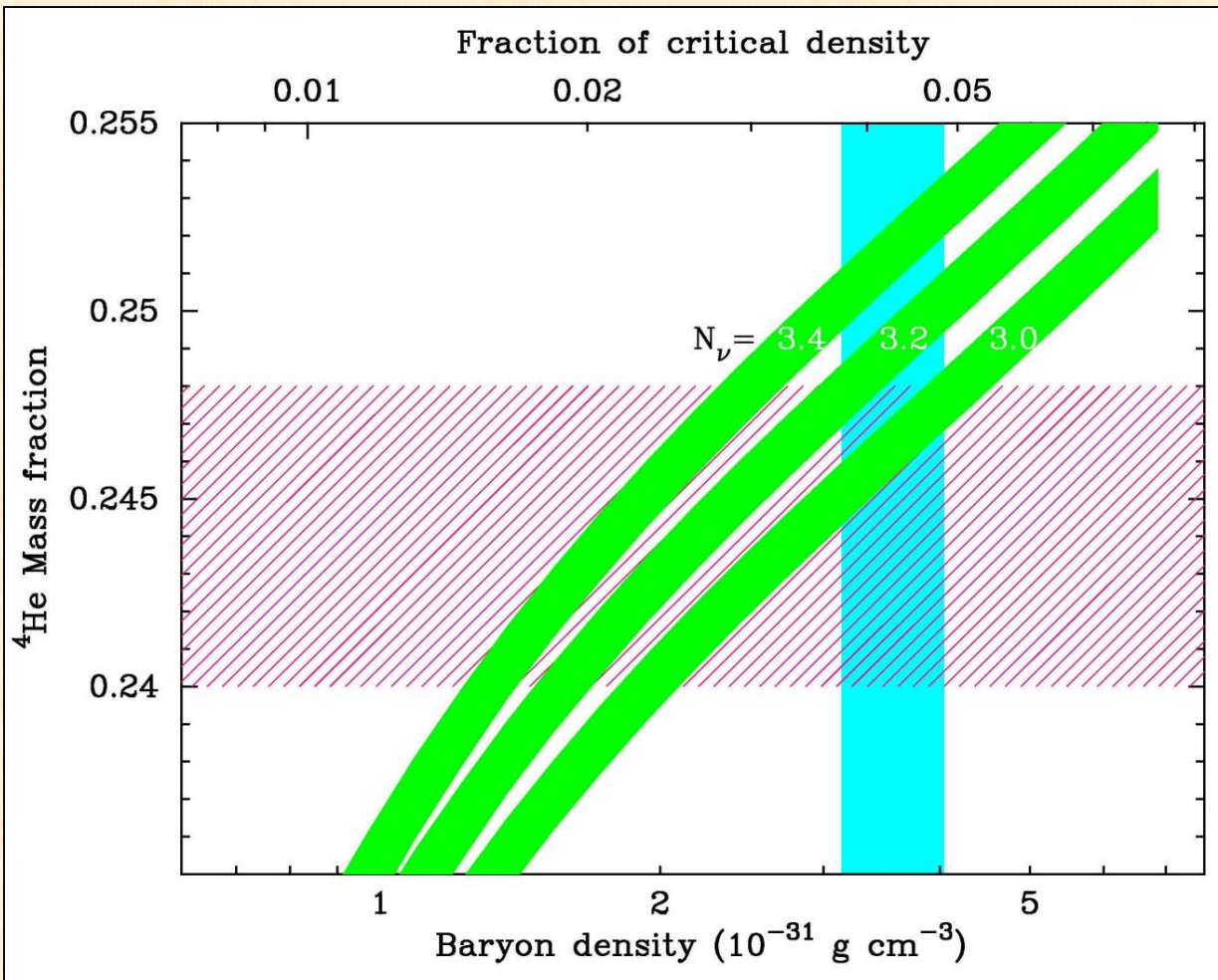
Deuterium

Lithium

Cosmic baryon
density implied by
deuterium abundance:

$$\Omega_B h^2 = 0.019 \pm 0.0024$$

BBN Limits on Neutrino Flavors



At BBN one flavor contributes about 16% to cosmic mass-energy density. Extra flavors modify expansion parameter accordingly.

Conservative limit:
 $|\Delta N_{\text{eff}}| < 1$

[astro-ph/9903300]

BBN and Neutrino Chemical Potentials

Expansion-rate effect:

Energy density in one neutrino flavor with degeneracy parameter $\xi = \eta/T$

$$\rho_{\nu\bar{\nu}} = \frac{7\pi^2}{120} T_\nu^4 \left[1 + \underbrace{\frac{30}{7} \left(\frac{\xi}{\pi}\right)^2 + \frac{15}{7} \left(\frac{\xi}{\pi}\right)^4}_{\Delta N_{\text{eff}}} \right]$$

Beta equilibrium effect for electron flavor: $n + \nu_e \leftrightarrow p + e^-$

Helium abundance essentially fixed by n/p ratio at beta freeze-out

$$\frac{n}{p} = e^{-(m_n - m_p)/T - \eta/T}$$

Numerically effect on helium analogous to $\Delta N_{\text{eff}} \sim -18 \xi_{\nu_e}$.

$$|\Delta N_{\text{eff}}| < 1 \rightarrow |\xi_{\nu_e}| < 0.057$$

Beta effect of ν_e can compensate expansion-rate effect of ν_μ or ν_τ
No significant BBN limit on neutrino number density

Chemical Potentials and Flavor Oscillations

Flavor mixing (neutrino oscillations)

Flavor lepton numbers not conserved

Only one common chemical potential ξ

Stringent ξ_{ν_e} limit applies to all flavors

For all flavors $|\xi_{\nu_e, \mu, \tau}| < 0.07$

Extra neutrino density $\Delta N_{\text{eff}} < 0.0064$

Cosmic neutrino density very close to standard value.
Neutrino-mass experiments indeed measure or
constrain cosmic neutrino density

Chain of arguments only valid if relaxation to flavor equilibrium obtains
before n/p freeze-out, i.e. before $T \sim 1 \text{ MeV}$

Flavor Equilibrium Before N/P Freeze-Out ?

yes

- Atmospheric neutrinos: Maximally mixed $\nu_\mu \rightarrow \nu_\tau$ oscillations
- Solar neutrinos by LMA solution

maybe

Solar LOW solution
(Quasi-equilibrium depends on unknown third mixing angle θ_3)

no

Solar SMA solution

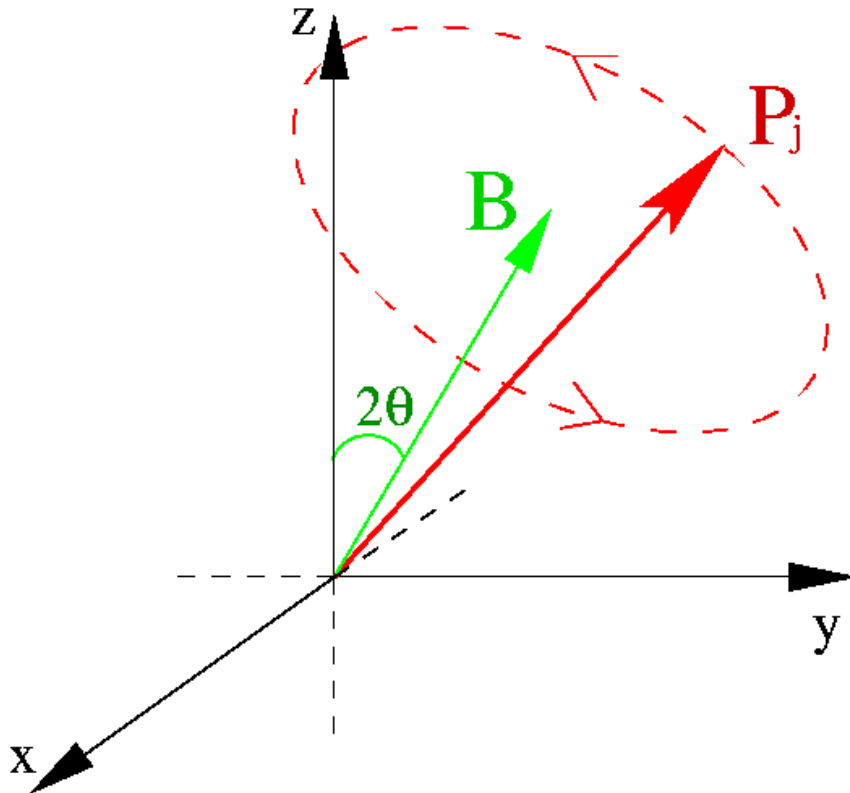
Our knowledge of the cosmic neutrino density depends on the solution of the solar neutrino problem: KamLAND most relevant current experiment

- Lunardini & Smirnov, hep-ph/0012056
- Dolgov, Hansen, Pastor, Petcov, Raffelt & Semikoz, hep-ph/0201287
- Abazajian, Beacom & Bell, astro-ph/0203442
- Wong, hep-ph/0203180

Two-Flavor Neutrino Oscillations

Evolution of neutrino ensemble described in terms of density matrices:

$$f_p \rightarrow \rho_p = \begin{pmatrix} f_p^e & f_p^{\mu e} \\ f_p^{e\mu} & f_p^{\mu} \end{pmatrix} = \frac{1}{2} (f_p + \vec{\sigma} \cdot \vec{P}_p)$$



Flavor oscillations in vacuum:

$$\partial_t \vec{P}_p = \frac{\delta m^2}{2p} \vec{B} \times \vec{P}_p$$

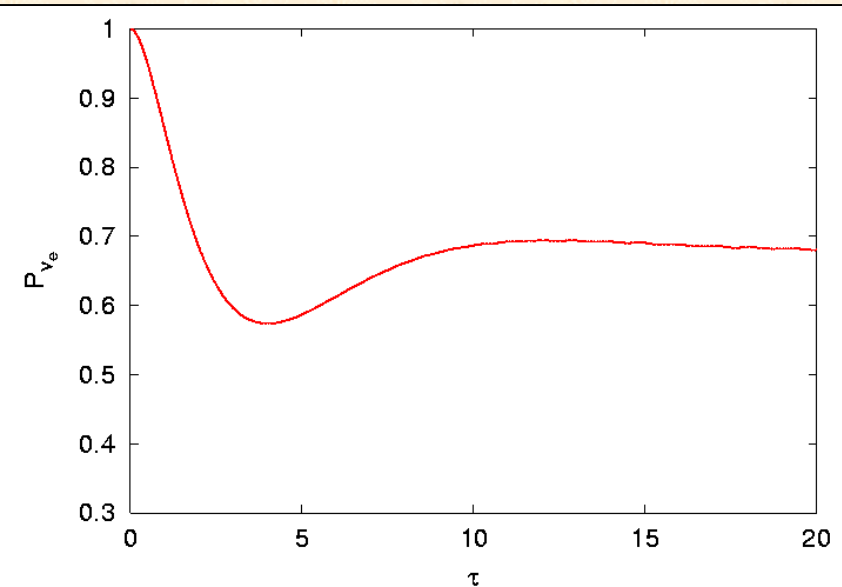
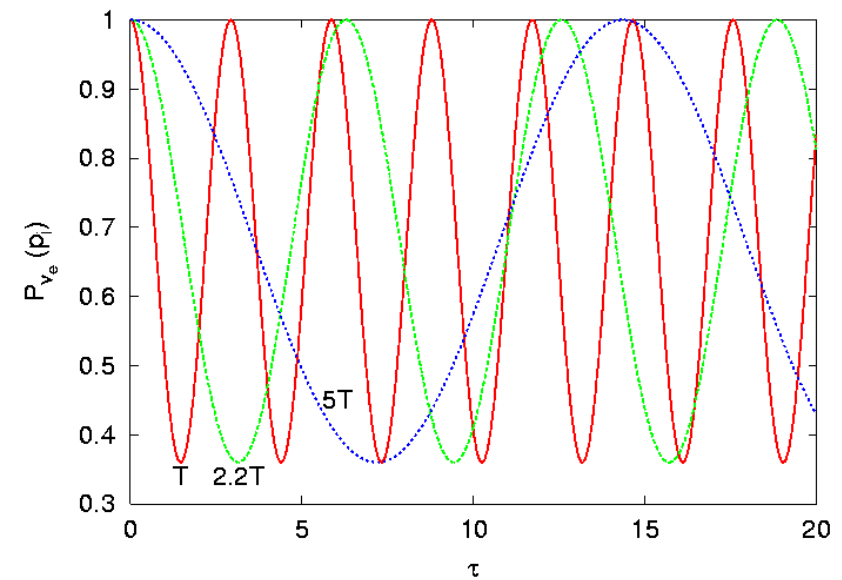
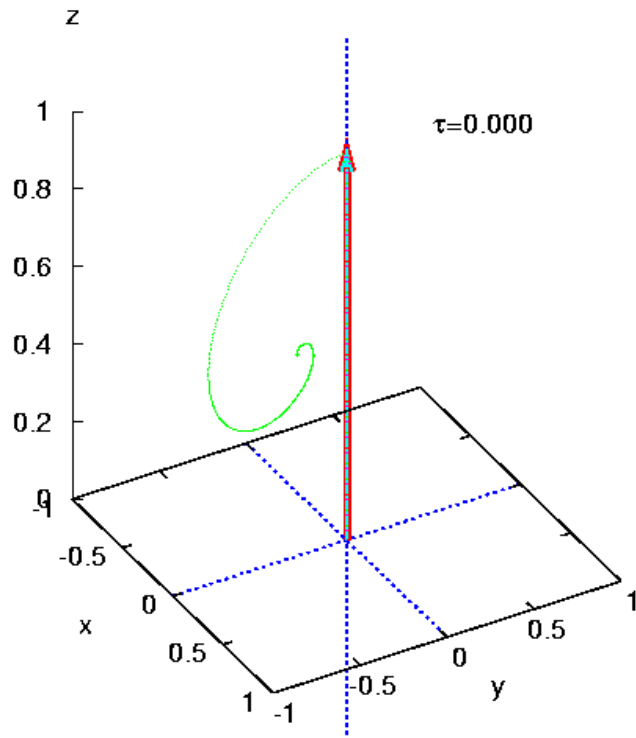
with

$$\vec{B} = \begin{pmatrix} \sin 2\theta \\ 0 \\ \cos 2\theta \end{pmatrix}$$

and

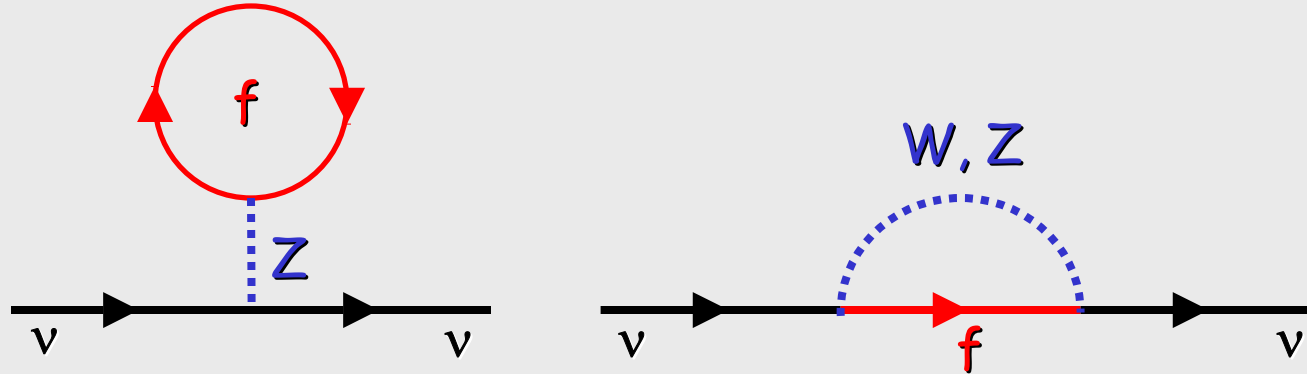
$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

Vacuum Oscillations with Distribution of Energies



Two-Flavor Oscillations in Media

Neutrinos propagating in a medium suffer refraction (Wolfenstein 1978)



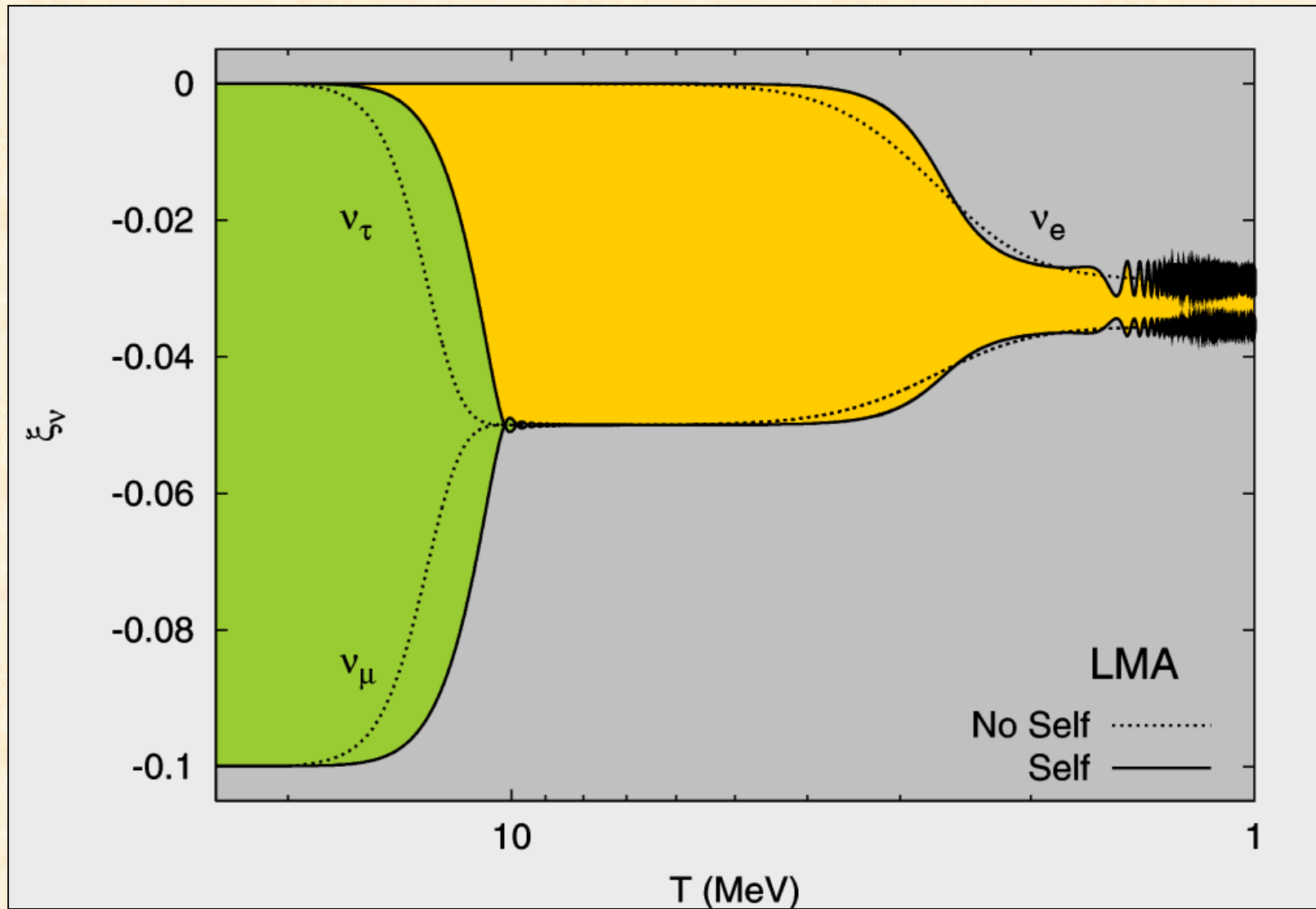
Effect is usually different for different flavors

Equation of motion in early universe (ignoring neutrino background)

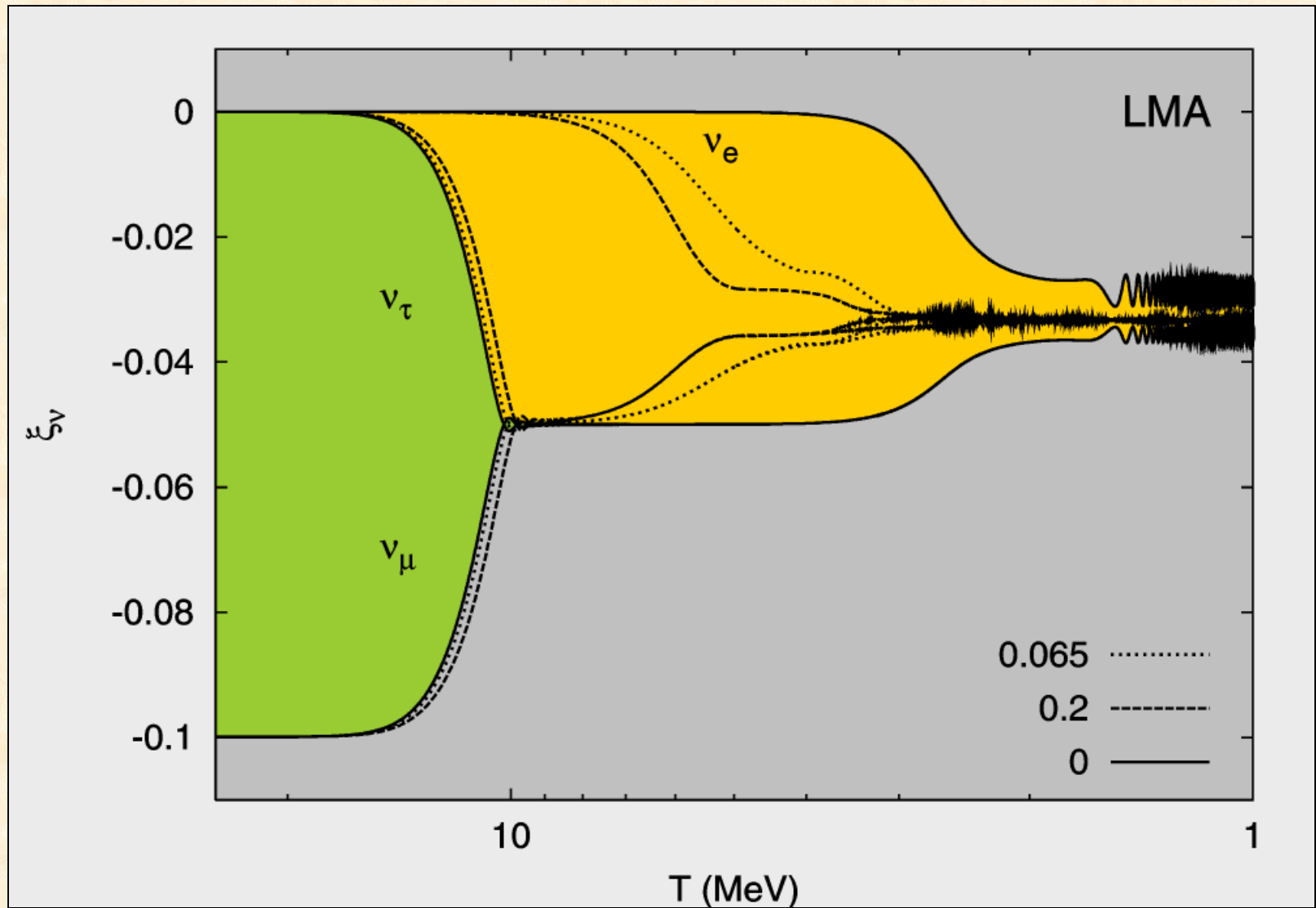
$$\partial_t \vec{P}_p = \left(\frac{\delta m^2}{2p} \vec{B} - \frac{8\sqrt{2}G_F p}{3m_W^2} \rho_e \vec{Z} \right) \times \vec{P}_p \quad \text{with } \rho_e \text{ the } e^+e^- \text{ energy density}$$

Oscillations begin when background medium is sufficiently diluted to avoid large medium effect compared to vacuum mixing

Flavor Equilibration: LMA Solution

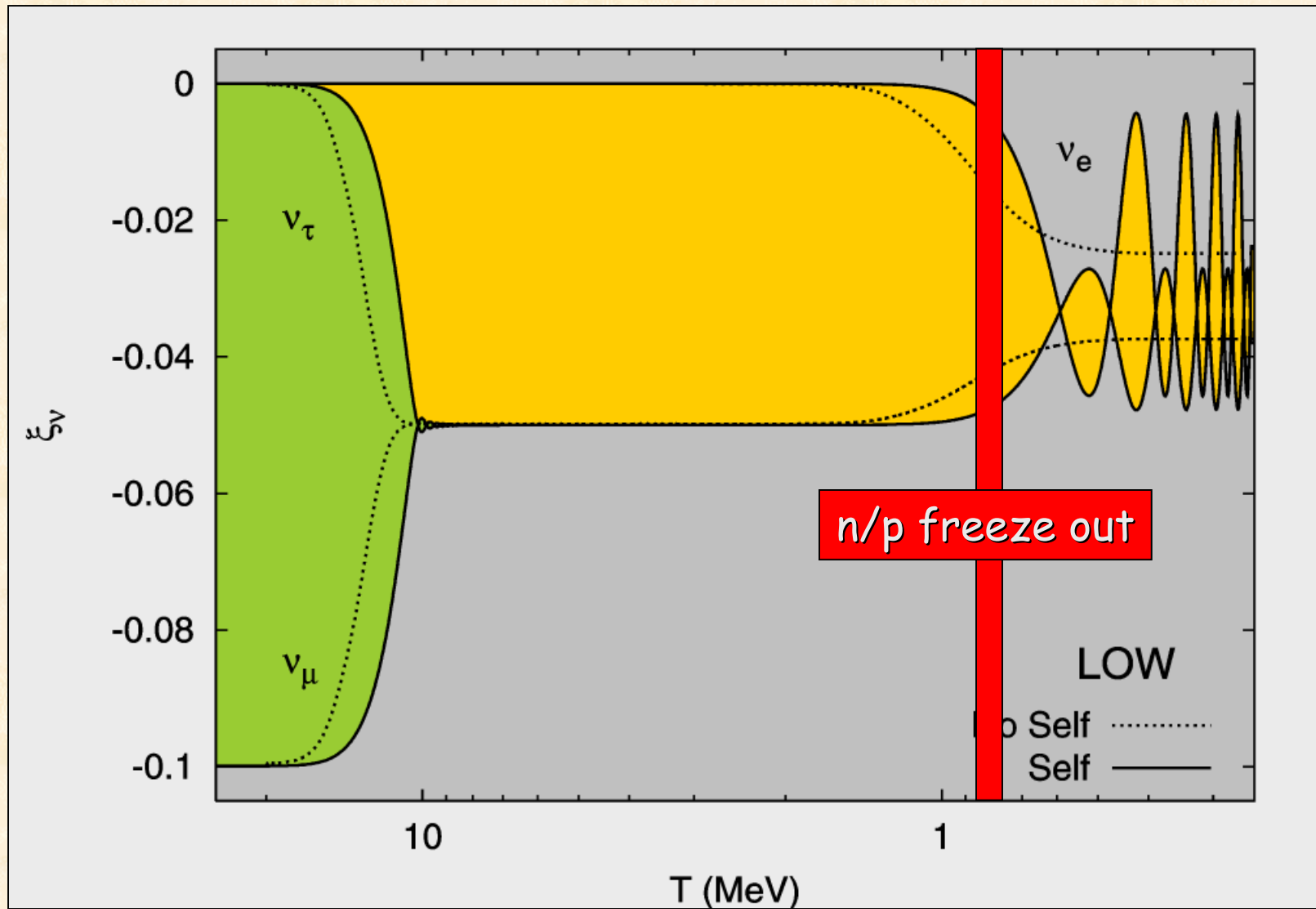


Flavor Equilibration: LMA With Non-Zero Θ_{13}

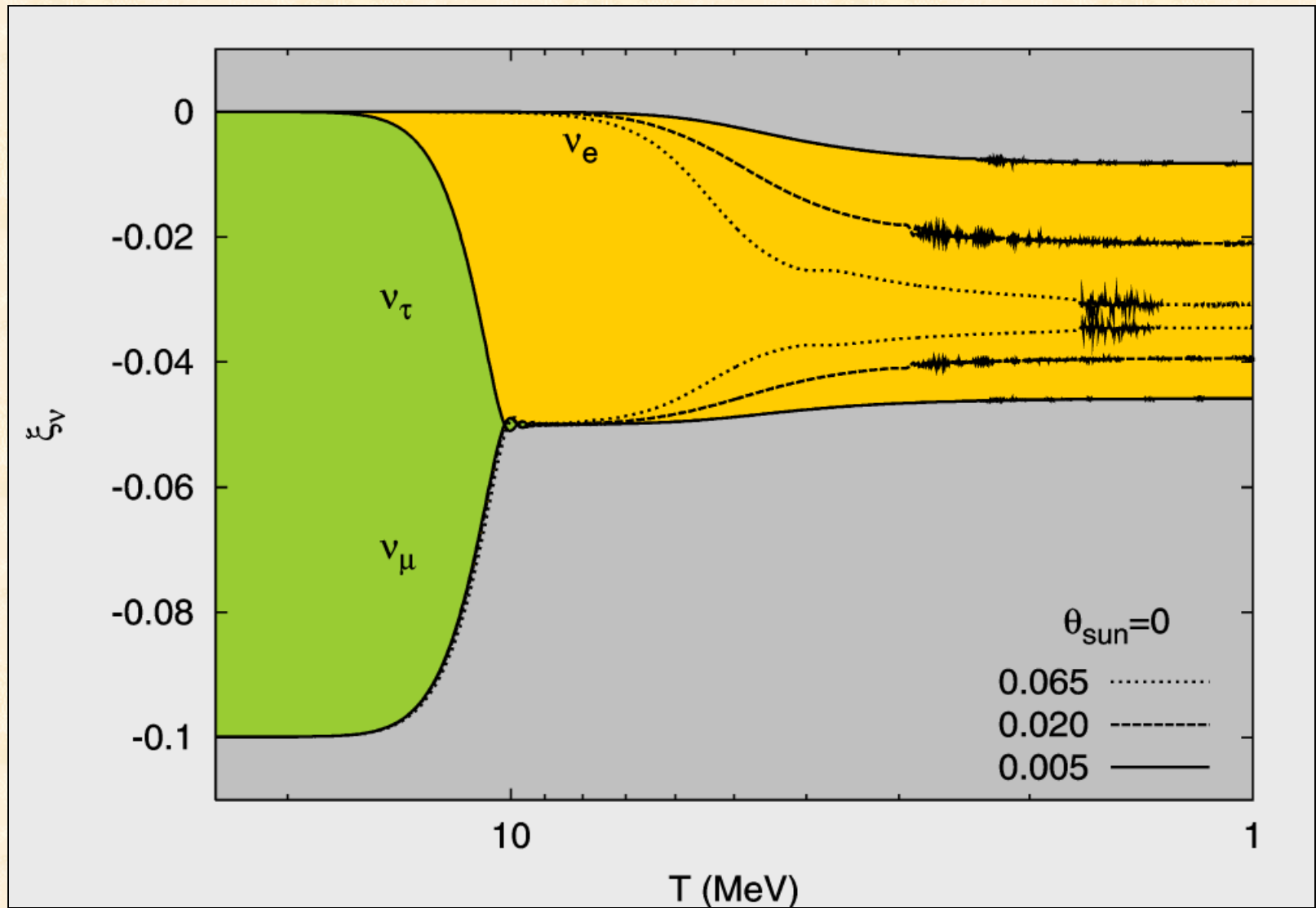


Dolgov, Hansen, Pastor, Petcov, Raffelt & Semikoz, hep-ph/0201287

Flavor Equilibration: LOW Solution



Flavor Equilibration: SMA With Non-Zero Θ_{13}



Dolgov, Hansen, Pastor, Petcov, Raffelt & Semikoz, hep-ph/0201287

Synchronized Oscillations by Self-Interactions

Equation of motion in early universe with neutrino background

$$\partial_t \mathbf{P}_p = + \left(\frac{\delta m^2}{2p} \mathbf{B} - \frac{8\sqrt{2}G_F p}{3m_W^2} \rho_e \mathbf{z} \right) \times \mathbf{P}_p + \sqrt{2}G_F (\mathbf{P} - \bar{\mathbf{P}}) \times \mathbf{P}_p \quad \text{neutrinos}$$

$$\partial_t \bar{\mathbf{P}}_p = - \left(\frac{\delta m^2}{2p} \mathbf{B} - \frac{8\sqrt{2}G_F p}{3m_W^2} \rho_e \mathbf{z} \right) \times \bar{\mathbf{P}}_p + \sqrt{2}G_F (\mathbf{P} - \bar{\mathbf{P}}) \times \bar{\mathbf{P}}_p \quad \text{anti-neutrinos}$$

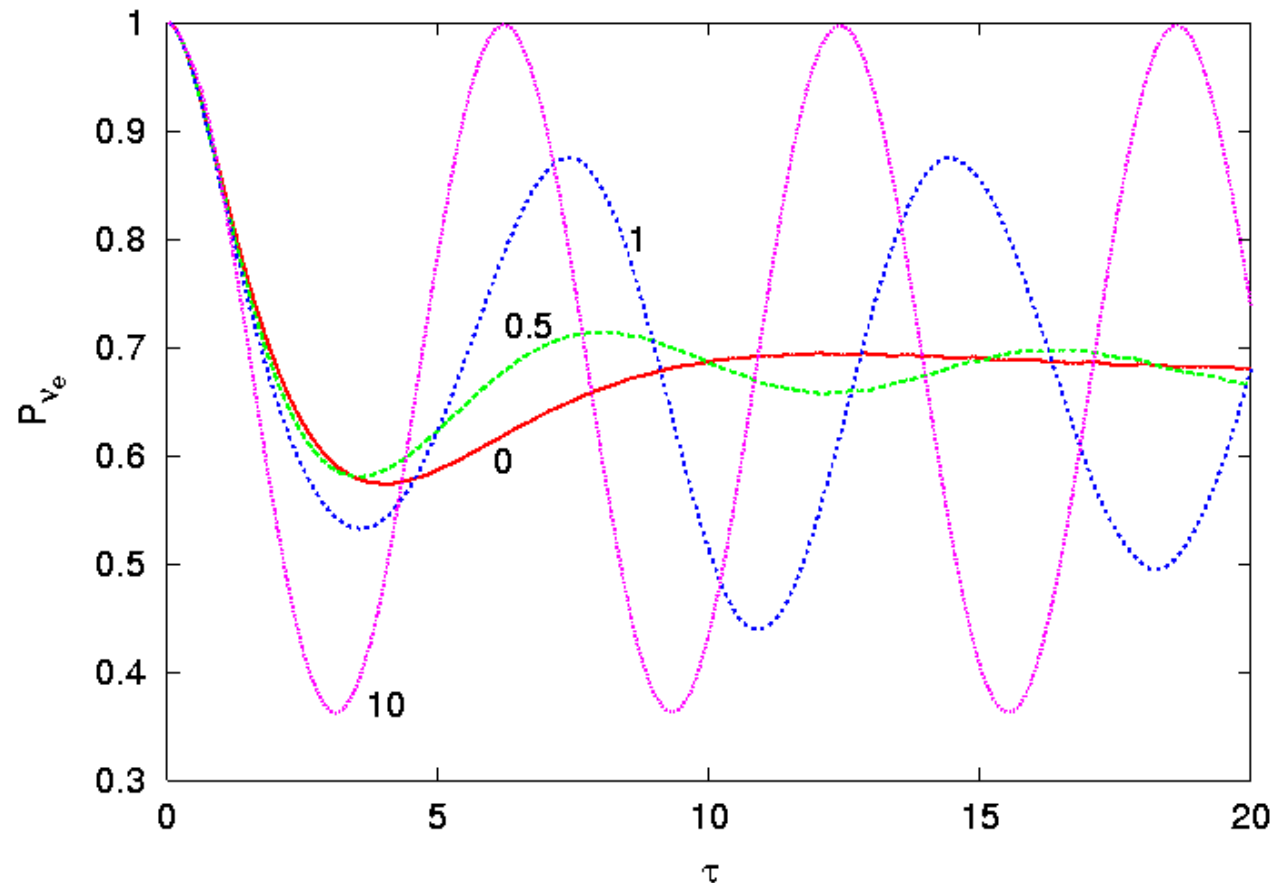
with the integrated neutrino polarization vectors

$$\mathbf{P} = \int \frac{d^3 p}{(2\pi)^3} \mathbf{P}_p \quad \text{and} \quad \bar{\mathbf{P}} = \int \frac{d^3 p}{(2\pi)^3} \bar{\mathbf{P}}_p$$

"Magnetic field" caused by neutrinos themselves much larger than vacuum or medium terms.

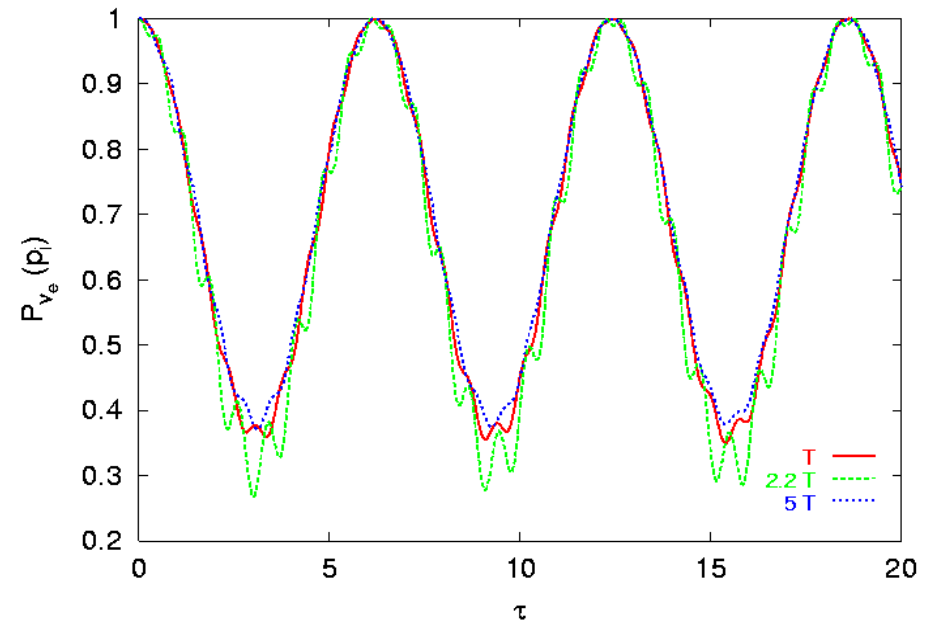
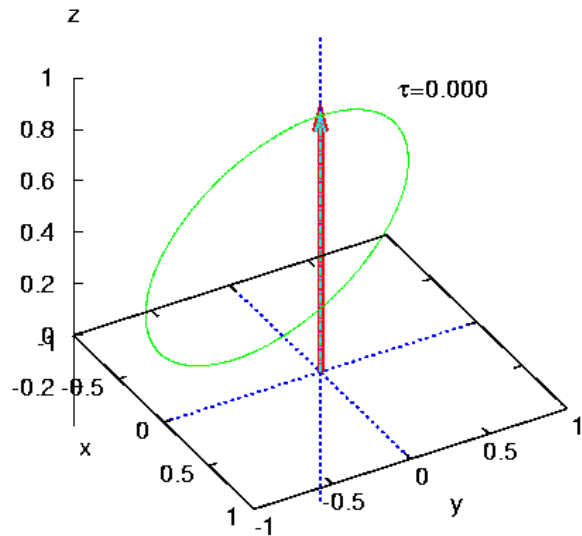
Couples "magnetic moments" to one large dipole which precesses with a single frequency.

Synchronized Oscillations by Self-Interactions



Pastor, Raffelt & Semikoz, PRD 65 (2002) 053011

Synchronized Oscillations by Self-Interactions



Pastor, Raffelt & Semikoz, PRD 65 (2002) 053011

Zero-Frequency Synchronized Oscillations

Synchronized oscillation frequency, assuming all polarization vectors initially parallel or antiparallel to z-axis

$$\omega_{\text{synch}} = \frac{\int \frac{d^3p}{(2\pi)^3} \frac{\delta m^2}{2p} (p_p + \bar{p}_p)}{\left| \int \frac{d^3p}{(2\pi)^3} (p_p - \bar{p}_p) \right|}$$

Only neutrinos (no anti-neutrinos):

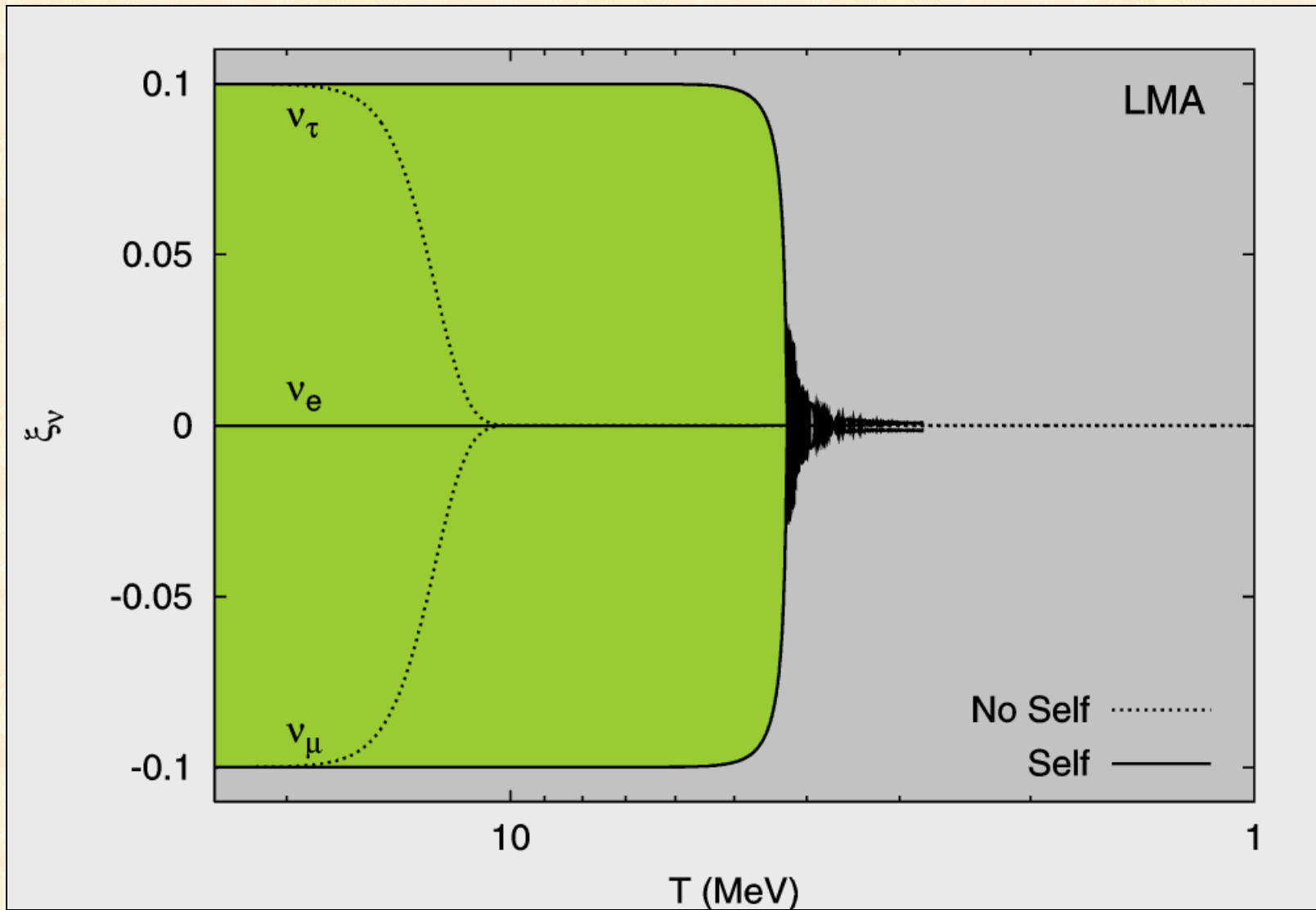
$$\omega_{\text{synch}} = \frac{\delta m^2}{2} \left\langle \frac{1}{p} \right\rangle$$

Equal distribution of neutrinos of one flavor and anti-neutrinos of the other:

$$\omega_{\text{synch}} = 0$$

Oscillations completely suppressed

Flavor Equilibration: LMA and $\xi(\nu_\mu) = -\xi(\nu_\tau)$



Zero-Frequency Synchronized Oscillations

In a three-flavor system with dominant self-term, oscillations are suppressed (or infinitesimally slow) when

$$|\xi(\nu_e)| = |\xi(\nu_\mu)| = |\xi(\nu_\tau)|$$

Therefore, “equilibrium” in a strict sense not assured, but in that case number densities anyway equal between all flavors.

This situation requires a nonvanishing total lepton number.
If total lepton number is zero, oscillation suppression by self-term not possible.

Summary

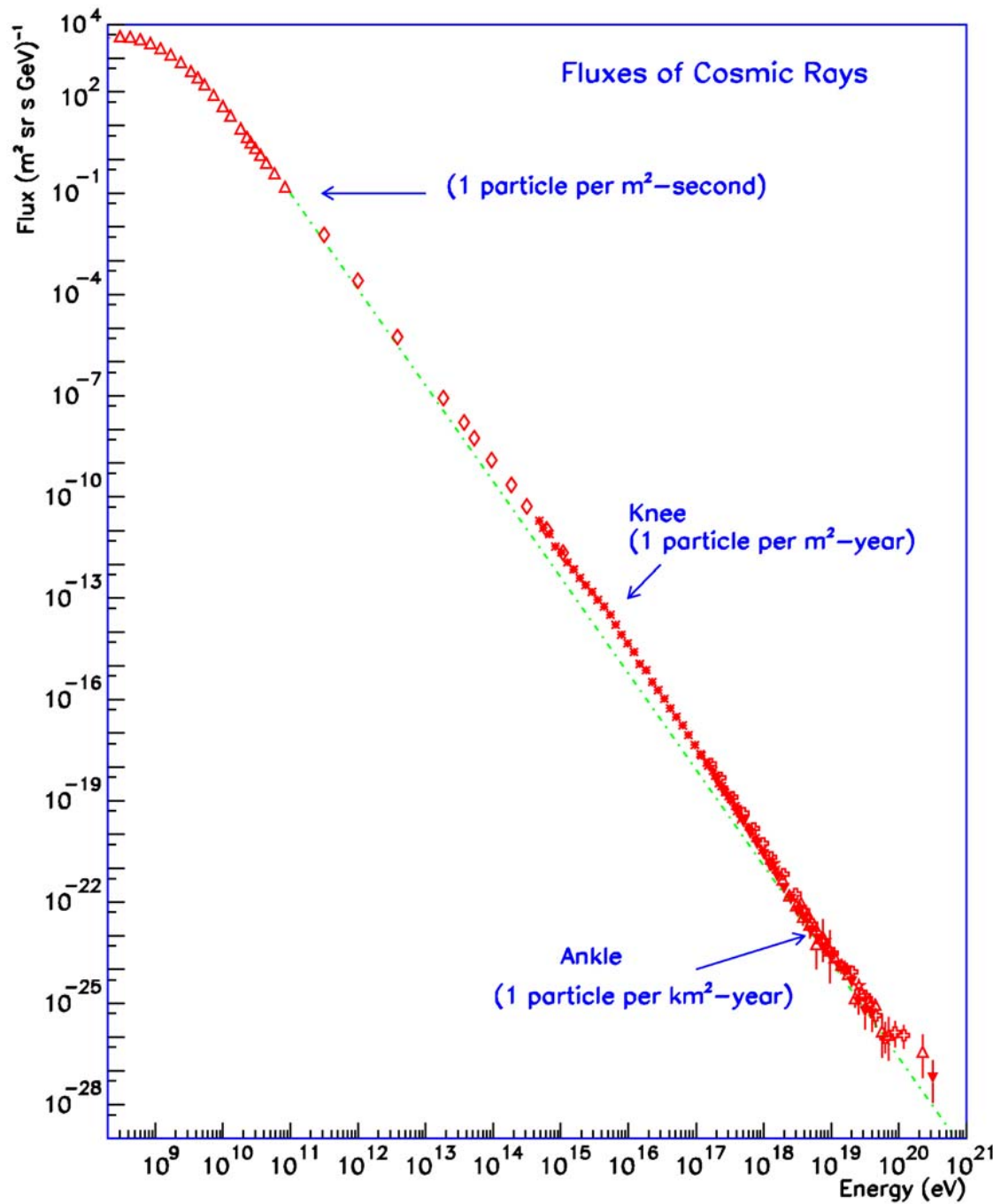
If KamLAND confirms LMA solar oscillation solution,
BBN implies
that the total number density of relic neutrinos agrees
with the standard value to better than 1%

Strict relation between cosmological limits on
hot dark matter component
and neutrino mass

Highest-Energy Cosmic Rays and the Cosmic Neutrino Sea



Global Cosmic Ray Spectrum



END TO THE COSMIC-RAY SPECTRUM?

Kenneth Greisen

Cornell University, Ithaca, New York

(Received 1 April 1966)

Greisen-Zatsepin-Kuzmin (GZK) Cutoff



PRL 16 (1966) 748

UPPER LIMIT OF THE SPECTRUM OF COSMIC RAYS

G. T. Zatsepin and V. A. Kuz'min

P. N. Lebedev Physics Institute, USSR Academy of Sciences

Submitted 26 May 1966

ZhETF Pis'ma 4, No. 3, 114-117, 1 August 1966

JETP Lett.
4 (1966) 78

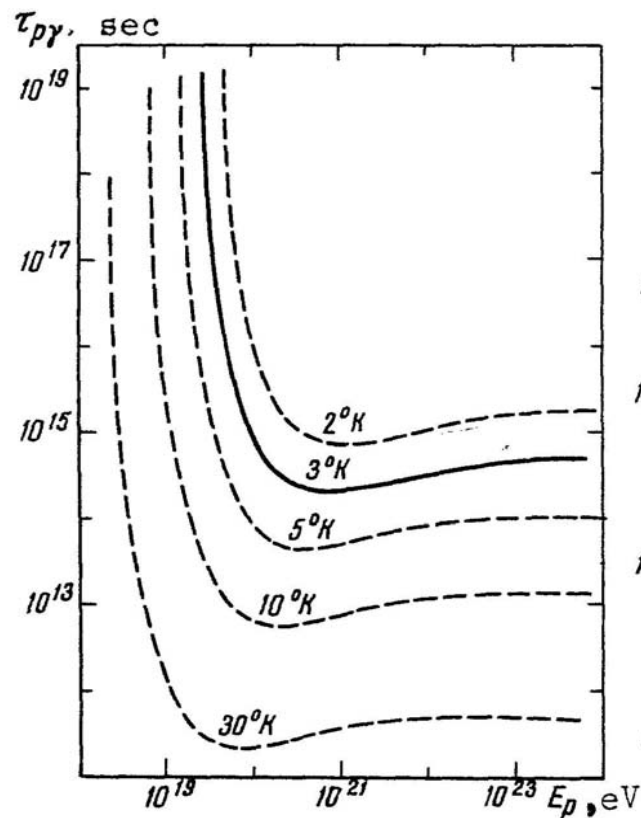


Fig. 1

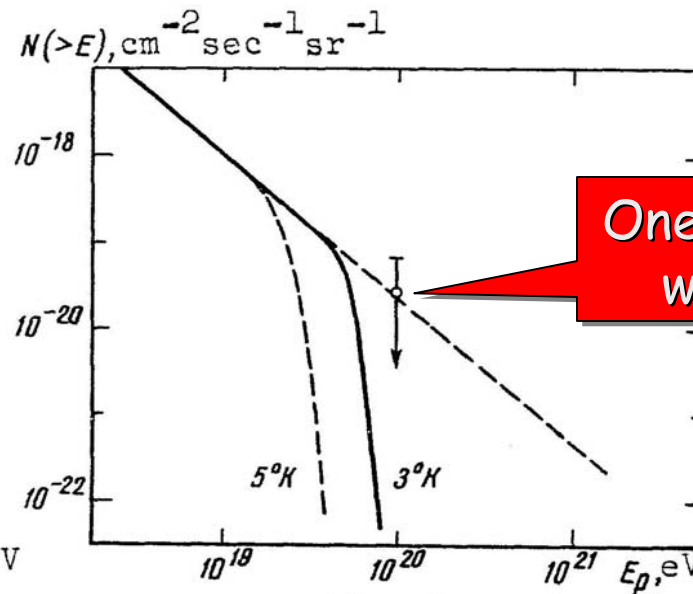
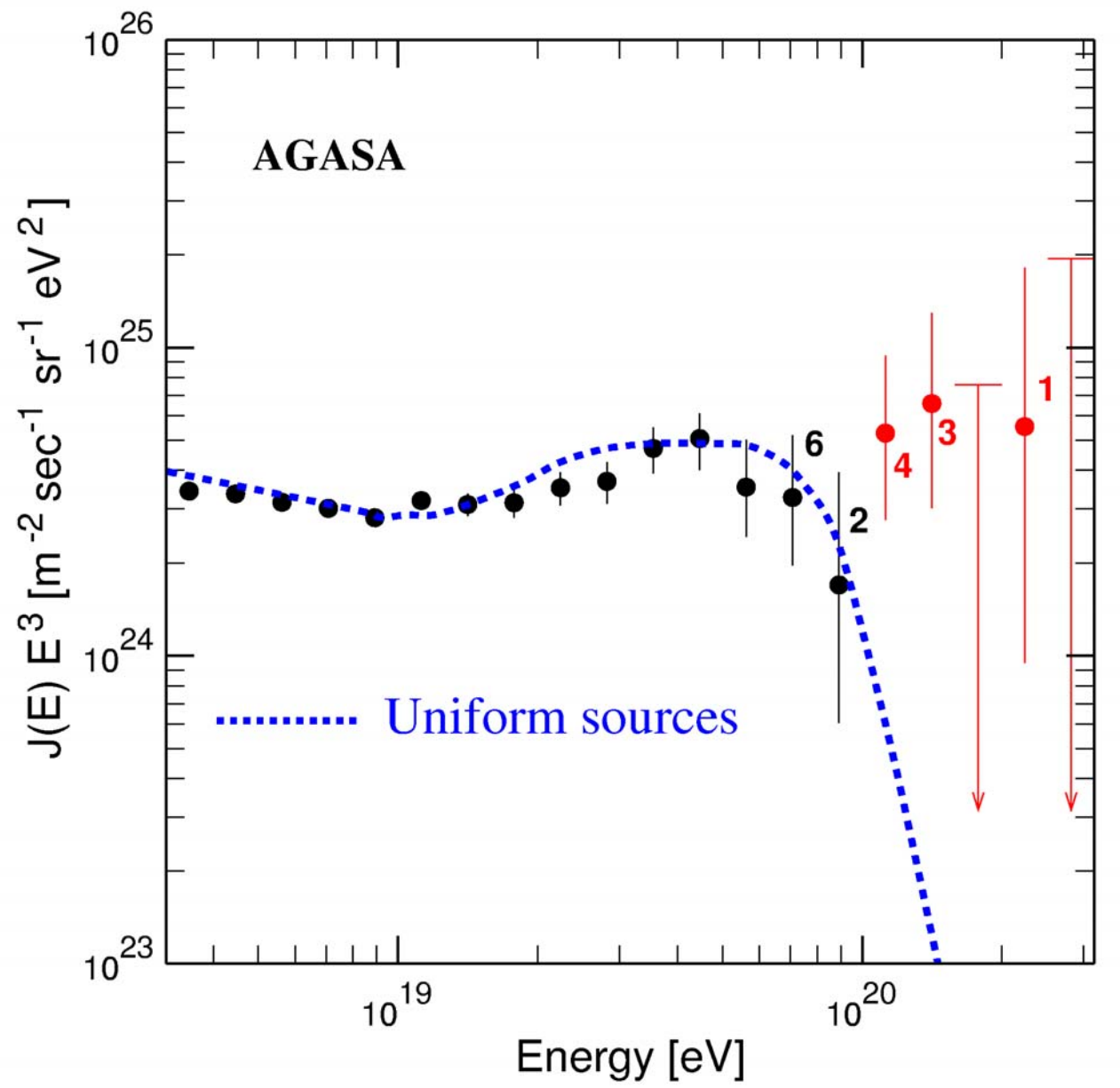


Fig. 2

One event (1962)
with 10^{20} eV

Spectrum of Highest-Energy AGASA Events

astro-ph/0008102



Z-Bursts and Highest-Energy Cosmic Rays

Neutrinos
 $E_\nu \sim 10^{21} - 10^{22} \text{ eV}$
from unknown sources

Resonant
Z-Boson
Production

Cosmic relic
neutrinos
 $m_\nu \sim 1 \text{ eV}$

Neutrino energy
on resonance

$$E_\nu = \frac{M_Z^2}{2m_\nu} = \frac{4.2 \times 10^{21} \text{ eV}}{m_{\text{eV}}}$$

Decay
(Z-Burst)

2 Nucleons

$10 \pi^0 \rightarrow 20 \gamma$

$17 \pi^\pm \rightarrow e^\pm \nu \bar{\nu}$

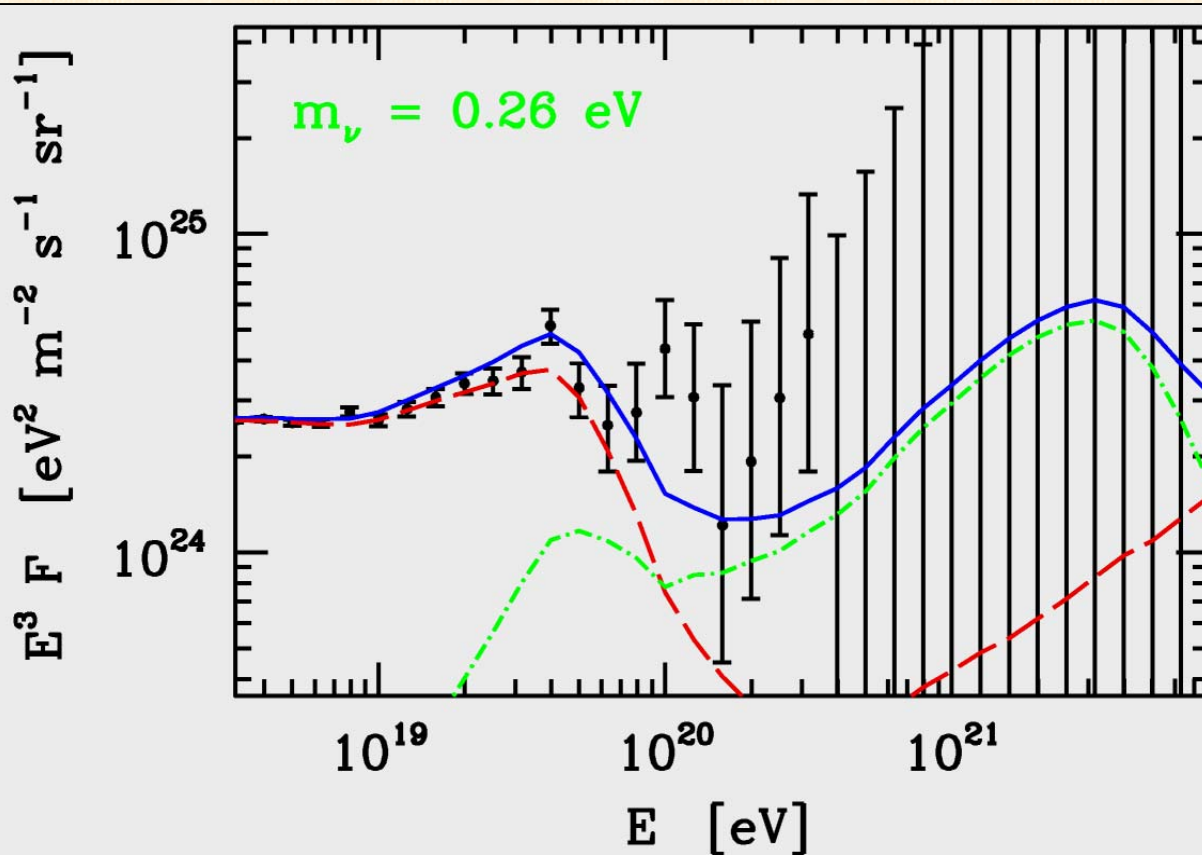
Measured cosmic rays

For example:

- Weiler
hep-ph/9710431;
- Fargion et al.
hep-ph/0112014
- Fodor, Katz, Ringwald
hep-ph/0203198;

Fitting the Cosmic Ray Spectrum with Z-Bursts

Cosmic ray spectrum near cutoff can be fit for a wide range of allowed neutrino masses



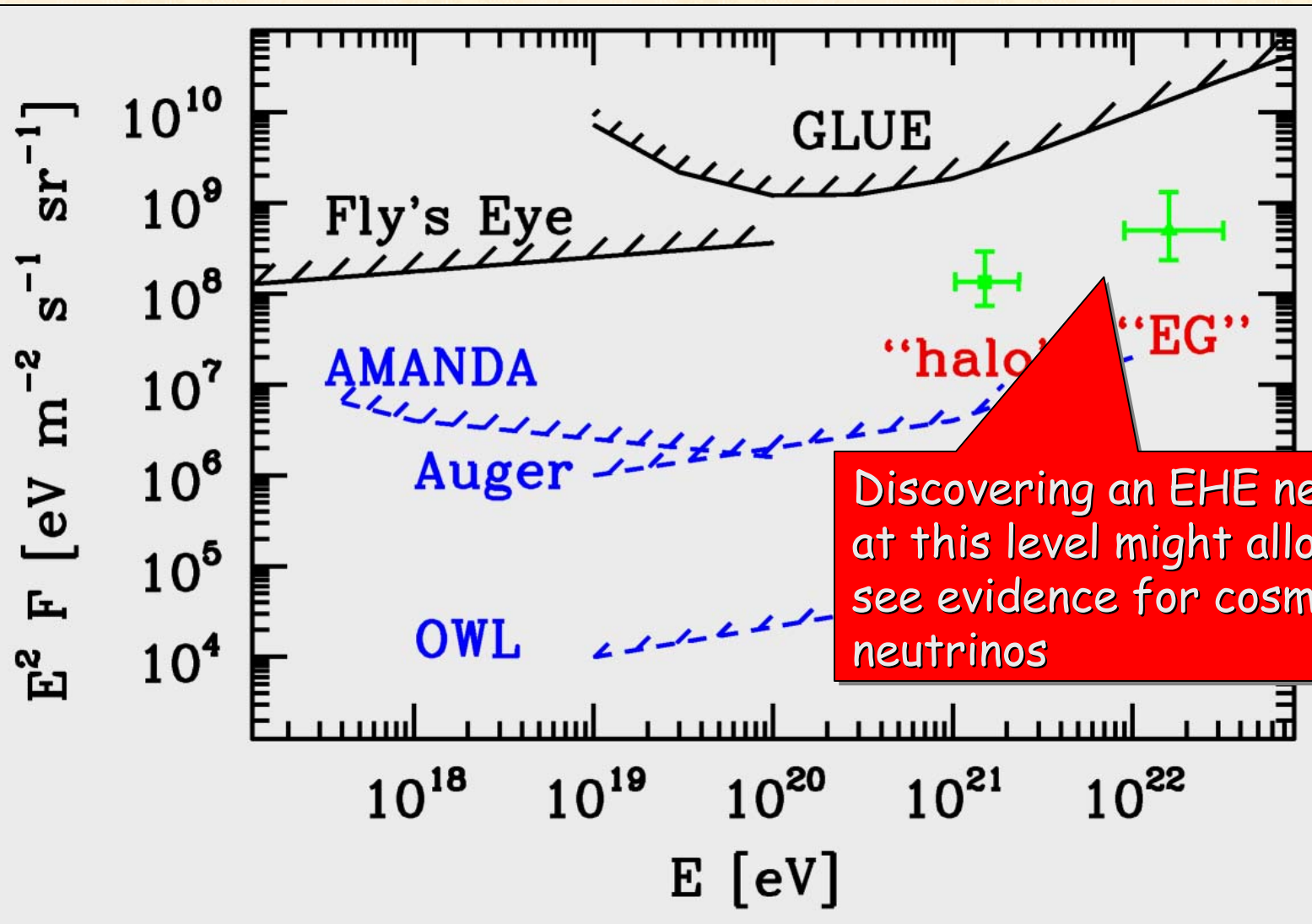
Main problems:

- Huge source flux of EHE neutrinos required
- No plausible sources known
- Accompanying photons must be perfectly obscured

Kalashev et al.
hep-ph/0112351

Fodor, Katz & Ringwald, hep-ph/0203198

Discovery Potential for Required Neutrino Fluxes



Fodor, Katz & Ringwald, hep-ph/0203198

Massive Neutrinos and the Cosmic Baryon Asymmetry



Baryogenesis in the Early Universe

Sakharov conditions for creating the Baryon Asymmetry of the Universe (BAU)

- C and CP violation
- Baryon number violation
- Deviation from thermal equilibrium

Particle-physics standard model

- Violates B and L by electroweak instanton effects
- Conserves $B - L$

In cosmological evolution

- Pre-existing $B+L$ erased at EW phase transition
- Creation of BAU at phase transition not possible, except for special parameters in SUSY models

Leptogenesis by Majorana Neutrino Decays

A classic paper:

Volume 174, number 1

PHYSICS LETTERS B

26 June 1986

BARYOGENESIS WITHOUT GRAND UNIFICATION

M. FUKUGITA

Research Institute for Fundamental Physics, Kyoto University, Kyoto 606, Japan

and

T. YANAGIDA

*Institute of Physics, College of General Education, Tohoku University, Sendai 980, Japan
and Deutsches Elektronen-Synchrotron DESY, D-2000 Hamburg, Fed. Rep. Germany*

Received 8 March 1986

A mechanism is pointed out to generate cosmological baryon number excess without resorting to grand unified theories. The lepton number excess originating from Majorana mass terms may transform into the baryon number excess through the unsuppressed baryon number violation of electroweak processes at high temperatures.

See-Saw Model for Neutrino Masses

Charged Leptons

Neutrinos

Dirac masses
from coupling
to standard
Higgs field ϕ



Heavy
Majorana
masses M_j

Lagrangian for
particle masses

$$L_{\text{mass}} = -\bar{\ell}_L \phi g_\ell \ell_R - \bar{\ell}_L \phi g_\nu N_R - \frac{1}{2} \bar{N}_R^c M N_R + \text{h.c.}$$

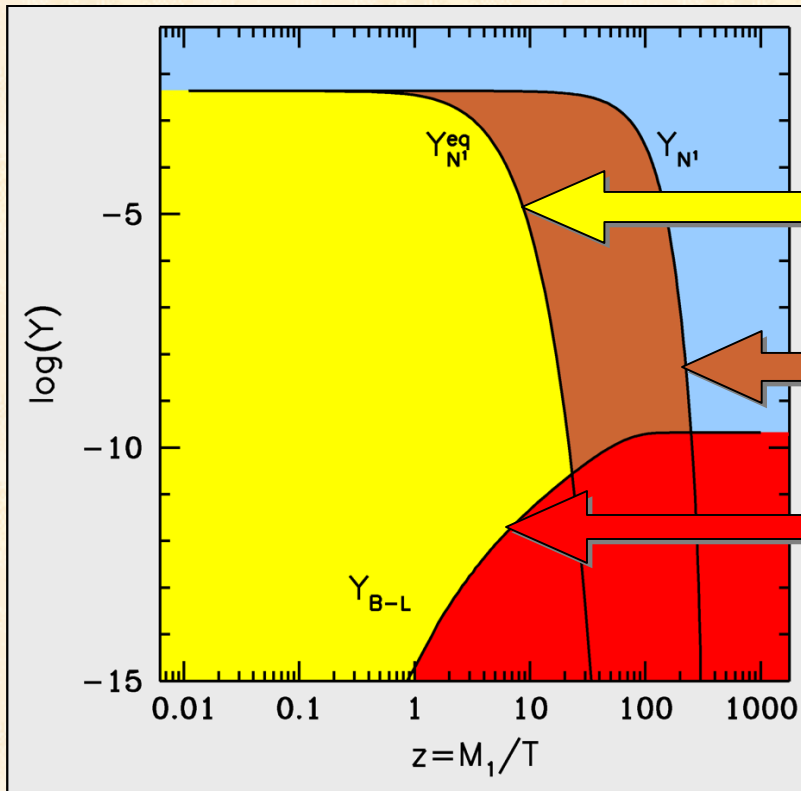
$$(\bar{\nu}_L \quad \bar{N}_R) \begin{pmatrix} 0 & g_\nu \langle \phi \rangle \\ g_\nu \langle \phi \rangle & M \end{pmatrix} \begin{pmatrix} \nu_L \\ N_R \end{pmatrix}$$

Diagonalize

Light Majorana mass

$$(\bar{\nu}_L \quad \bar{N}_R) \begin{pmatrix} \frac{g_\nu^2 \langle \phi \rangle^2}{M} & 0 \\ 0 & M \end{pmatrix} \begin{pmatrix} \nu_L \\ N_R \end{pmatrix}$$

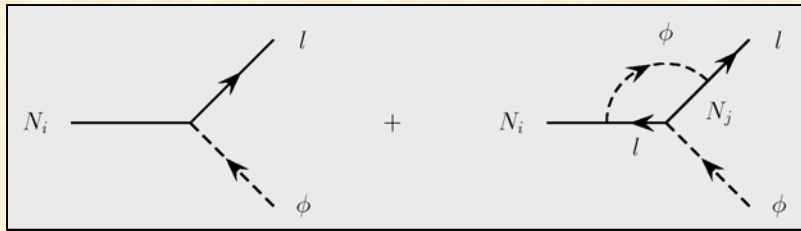
Leptogenesis by Out-of-Equilibrium Decay



Equilibrium abundance of heavy Majorana neutrinos

Real non-equilibrium abundance determined by decay rate

Lepton-number abundance created by CP-violating decays



CP-violating decays by interference of tree-level with one-loop diagram

$$\Gamma_{\text{Decay}} = g_v^2 \frac{M}{8\pi}$$

Connection to Neutrino Mass

$$\Gamma_{\text{Decay}} = g_v^2 \frac{M}{8\pi}$$

Decay rate of heavy Majorana neutrino

$$H \approx \sqrt{g_{\text{eff}}} \frac{T^2}{m_{\text{Pl}}}$$

Cosmic expansion rate

$$\Gamma_{\text{Decay}} < H|_{T=M}$$

Requirement for strong deviation from equilibrium ...

$$g_v^2 \frac{M}{8\pi} < \sqrt{g_{\text{eff}}} \frac{M^2}{m_{\text{Pl}}}$$

$$\frac{g_v^2}{M} < \frac{8\pi\sqrt{g_{\text{eff}}}}{m_{\text{Pl}}}$$

$$m_\nu = \frac{g_v^2 \langle \phi \rangle^2}{M} < \frac{8\pi\sqrt{g_{\text{eff}}}}{m_{\text{Pl}}} \langle \phi \rangle^2 \sim 10^{-3} \text{ eV}$$

... translates into a limit on the observable neutrino mass

Leptogenesis by Majorana Neutrino Decays

In see-saw models of neutrino masses, right-handed heavy Majorana neutrinos provide source for L-violation

Cosmological evolution:

- $B = L = 0$ early on
- Thermal freeze-out of heavy Majorana neutrinos
- Out-of-equilibrium decay creates net L
- Shift L excess into B at EW phase transition

Sufficient deviation from equilibrium distribution of heavy Majorana neutrinos at freeze-out

Limits on Yukawa couplings

Limits on *light* neutrino masses

Consistent with hierarchical masses below 0.1 eV

Leptogenesis - A Popular Research Topic

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AND MANY MORE ...

Conclusions



Future galactic SN signal is

- Sensitive to $\sim 1\text{-}3$ eV mass
- But not sub-eV range



Large-scale galaxy redshift surveys:

- Best limit of $m_\nu < 0.8$ eV,
- Future sensitivity ~ 0.3 eV



If solar LMA solution applies,
cosmic neutrino number density
precisely determined by BBN



If highest-E cosmic-ray neutrinos are
found, Z-bursts provide handle on m_ν



Majorana neutrino masses in the favored
range may indicate a leptogenesis scenario
for generating cosmic baryon asymmetry