THE BUBBLE OF NOTHING UNDER T-DUALITY

 S^1

Based on: [2312.09291] Matilda Delgado



 S_R^1

NOTHING IS ALWAYS AROUND THE CORNER

[McNamara, Vafa '19]

The cobordism conjecture implies that spacetime can always end in a theory of QG

There always exists a physical process that allows you to shrink the compact space to a point



NOTHING IS ALWAYS AROUND THE CORNER

[McNamara, Vafa '19]

The cobordism conjecture implies that spacetime can always end in a theory of QG

There always exists a physical process that allows you to shrink the compact space to a point

What is the simplest compactification?

The circle is a **consistent background** for all known string theories,

There must always exist a process that shrinks it to a point.











Cobordisms to nothing are finite action processes, but a decompactification takes and infinite amount of energy

This is too naive

1.



Q: Does the cobordism conjecture somehow imply that these blow-ups also exist?

This is too naive

- 1. Cobordisms to nothing are finite action processes, but a decompactification takes and infinite amount of energy
- T-duality exchanges KK and winding U(1) symmetries. The pinch-off breaks the winding U(1) symmetry: winding strings can slip off





Q: Does the cobordism conjecture somehow imply that these blow-ups also exist?

This is too naive

- 1. Cobordisms to nothing are finite action processes, but a decompactification takes and infinite amount of energy
- T-duality exchanges KK and winding U(1) symmetries. The pinch-off breaks the winding U(1) symmetry: winding strings can slip off



it must be that its dual breaks the momentum U(1) and depends non-trivially on the dual angle

[Gregory, Harvey, Moore, `97]



Q: Does the cobordism conjecture somehow imply that these blow-ups also exist?

This is too naive

- 1. Cobordisms to nothing are finite action processes, but a decompactification takes and infinite amount of energy
- T-duality exchanges KK and winding U(1) symmetries. The pinch-off breaks the winding U(1) symmetry: winding strings can slip off



it must be that its dual breaks the momentum U(1) and **depends non-trivially on the dual angle**

[Gregory, Harvey, Moore, `97]

Let's look at an example..

UNDER T-DUALITY

What about the simplest example, Witten's Bubble of Nothing? [Witten 82]



It is topologically allowed *in the absence of fermions* (so embedit in bosonic string theory)

$$ds^{2} = f(r)^{-1}dr^{2} + r^{2}\left[-dt^{2} + \cosh(t)^{2}d\Omega_{23}^{2}\right] + R_{0}^{2}f(r)d\theta^{2} \qquad f(r) = 1 - \left(\frac{r_{0}}{r}\right)^{23}$$

UNDER T-DUALITY

What about the simplest example, Witten's Bubble of Nothing? [Witten 82]



Witten's bubble of nothing is a smooth cobordism to nothing

It is topologically allowed *in the absence of fermions* (so embedit in bosonic string theory)

$$ds^{2} = f(r)^{-1}dr^{2} + r^{2}\left[-dt^{2} + \cosh(t)^{2}d\Omega_{23}^{2}\right] + R_{0}^{2}f(r)d\theta^{2} \qquad f(r) = 1 - \left(\frac{r_{0}}{r}\right)^{23}$$

constant dilaton 🔽 $\,\phi=\phi_0$

no curvature singularity 🔽

13

19/03/2024

UNDER T-DUALITY

What about the simplest example, Witten's Bubble of Nothing? [Witten 82]



It is topologically allowed in the absence of fermions (so embedit in bosonic string theory).

$$ds^{2} = f(r)^{-1}dr^{2} + r^{2}\left[-dt^{2} + \cosh(t)^{2}d\Omega_{23}^{2}\right] + R_{0}^{2}f(r)d\theta^{2} \qquad f(r) = 1 - \left(\frac{r_{0}}{r}\right)^{23}$$

no curvature singularity $\overline{oldsymbol{V}}$ constant dilaton $\overline{oldsymbol{V}}$ $\phi=\phi_0$

So how does it transform under T-duality?

UNDER T-DUALITY

What about the simplest example, Witten's Bubble of Nothing? [Witten 82]



Witten's bubble of nothing is a smooth cobordism to nothing

It is topologically allowed *in the absence of fermions* (so embedit in bosonic string theory)

$$ds^{2} = f(r)^{-1}dr^{2} + r^{2}\left[-dt^{2} + \cosh(t)^{2}d\Omega_{23}^{2}\right] + R_{0}^{2}f(r)d\theta^{2} \qquad f(r) = 1 - \left(\frac{r_{0}}{r}\right)^{23}$$

So how does it transform under T-duality?

TODAY: what is this mysterious dual object in (bosonic) String Theory?

T-duality trades a string theory on a circle of radius R for one on a radius $lpha' R^{+1}$

The Buscher rules dictate how the zero-modes of the metric and dilaton change

$$(\tilde{G}_s)_{\theta\theta} = \frac{1}{(G_s)_{\theta\theta}} \qquad (\tilde{G}_s)_{\mu\nu} = (G_s)_{\mu\nu} \qquad \tilde{\phi} = \phi - \frac{1}{2} \log |(G_s)_{\theta\theta}|$$

T-duality trades a string theory on a circle of radius R for one on a radius $lpha' R^{+1}$

The Buscher rules dictate how the zero-modes of the metric and dilaton change

$$(\tilde{G}_s)_{\theta\theta} = \frac{1}{(G_s)_{\theta\theta}} \qquad (\tilde{G}_s)_{\mu\nu} = (G_s)_{\mu\nu} \qquad \tilde{\phi} = \phi - \frac{1}{2}\log|(G_s)_{\theta\theta}|$$

Applying these rules to the BoN, we get: $d\tilde{s}^{2} = \left(e^{\frac{\phi_{0}}{6}}R_{0}^{2}f(r)\right)^{\frac{1}{12}}\left[f(r)^{-1}dr^{2} + (g_{bon})_{ij}dx^{i}dx^{j}\right] + e^{-\frac{23}{72}\phi_{0}}\left(R_{0}^{2}f(r)\right)^{-\frac{11}{12}}d\tilde{\theta}^{2}$ $\tilde{\phi} = \frac{11}{12}\phi_{0} - \frac{1}{2}\log\left[f(r)R_{0}^{2}\right].$ $\phi = \phi_{0}$ $\phi(r) \rightarrow \infty$

T-duality trades a string theory on a circle of radius R for one on a radius $lpha' R^{+1}$

The Buscher rules dictate how the zero-modes of the metric and dilaton change

$$(\tilde{G}_s)_{\theta\theta} = \frac{1}{(G_s)_{\theta\theta}} \qquad (\tilde{G}_s)_{\mu\nu} = (G_s)_{\mu\nu} \qquad \tilde{\phi} = \phi - \frac{1}{2} \log |(G_s)_{\theta\theta}|$$

Applying these rules to the BoN, we get: $d\tilde{s}^2 = \left(e^{\frac{\phi_0}{6}}R_0^2 f(r)\right)^{\frac{1}{12}} \left[f(r)^{-1}dr^2 + (g_{bon})_{ij}dx^i dx^j\right] + e^{-\frac{23}{72}\phi_0} \left(R_0^2 f(r)\right)^{-\frac{11}{12}} d\tilde{\theta}^2$ $\tilde{\phi} = \frac{11}{12}\phi_0 - \frac{1}{2}\log\left[f(r)R_0^2\right].$

Again, this is too naive!

 \Rightarrow Fitting the defect with a delta source: **the solution has finite tension**

T-duality trades a string theory on a circle of radius R for one on a radius $lpha' R^{+1}$

The Buscher rules dictate how the zero-modes of the metric and dilaton change

$$(\tilde{G}_s)_{\theta\theta} = \frac{1}{(G_s)_{\theta\theta}} \qquad (\tilde{G}_s)_{\mu\nu} = (G_s)_{\mu\nu} \qquad \tilde{\phi} = \phi - \frac{1}{2}\log|(G_s)_{\theta\theta}|$$

Applying these rules to the BoN, we get: $d\tilde{s}^{2} = \left(e^{\frac{\phi_{0}}{6}}R_{0}^{2}f(r)\right)^{\frac{1}{12}}\left[f(r)^{-1}dr^{2} + (g_{bon})_{ij}dx^{i}dx^{j}\right] + e^{-\frac{23}{72}\phi_{0}}\left(R_{0}^{2}f(r)\right)^{-\frac{11}{12}}d\tilde{\theta}^{2}$ $\tilde{\phi} = \frac{11}{12}\phi_{0} - \frac{1}{2}\log\left[f(r)R_{0}^{2}\right].$

Again, this is too naive!

 \Rightarrow Fitting the defect with a delta source: **the solution has finite tension**

 \Rightarrow The metric seems to preserve a momentum U(1): The dependence on the dual angle is smeared

T-duality trades a string theory on a circle of radius R for one on a radius $lpha' R^{+1}$

The Buscher rules dictate how the zero-modes of the metric and dilaton change

$$(\tilde{G}_s)_{\theta\theta} = \frac{1}{(G_s)_{\theta\theta}} \qquad (\tilde{G}_s)_{\mu\nu} = (G_s)_{\mu\nu} \qquad \tilde{\phi} = \phi - \frac{1}{2}\log|(G_s)_{\theta\theta}|$$

Applying these rules to the BoN, we get: $d\tilde{s}^2 = \left(e^{\frac{\phi_0}{6}}R_0^2f(r)\right)^{\frac{1}{12}}\left[f(r)^{-1}dr^2 + (g_{bon})_{ij}dx^i dx^j\right] + e^{-\frac{23}{72}\phi_0}\left(R_0^2f(r)\right)^{-\frac{11}{12}}d\tilde{\theta}^2$ $\tilde{\phi} = \frac{11}{12}\phi_0 - \frac{1}{2}\log\left[f(r)R_0^2\right]$.

So the dual is some sort of 24-dimensional object with finite tension ... What else can we say?

BEYOND GRAVITY

The T-dual of the BoN at large radius is dual to a **stringy sized** bubble of nothing.

Winding modes become light and break the U(1) momentum symmetry. $\Psi_n \sim \Psi(r) \cos(n ilde{ heta})$

BEYOND GRAVITY

The T-dual of the BoN at large radius is dual to a **stringy sized** bubble of nothing.

Winding modes become light and break the U(1) momentum symmetry. $\Psi_n \sim \Psi(r) \cos(n ilde{ heta})$

BEYOND GRAVITY

The T-dual of the BoN at large radius is dual to a **stringy sized** bubble of nothing.

Winding modes become light and break the U(1) momentum symmetry. $\Psi_n \sim \Psi(r)\cos(n heta)$

Only a full-fledged sigma model for the BoN can (maybe) help us.

12

19/03/2024

BEYOND GRAVITY

The T-dual of the BoN at large radius is dual to a **stringy sized** bubble of nothing.

Winding modes become light and break the U(1) momentum symmetry. $\Psi_n \sim \Psi(r) \cos(n heta)$

Only a **full-fledged sigma model** for the BoN can (maybe) help us.

This is difficult because:

1.

The BoN is **time-dependent**
$$ds^2 = f(r)^{-1}dr^2 + r^2[-dt^2 + \cosh(t)^2 d\Omega_{23}^2] + R_0^2 f(r) dt^2$$

24

19/03/2024

BEYOND GRAVITY

The T-dual of the BoN at large radius is dual to a **stringy sized** bubble of nothing.

Winding modes become light and break the U(1) momentum symmetry. $\Psi_n \sim \Psi(r) \cos(n heta)$

Only a **full-fledged sigma model** for the BoN can (maybe) help us.

This is difficult because:

$$ds^{2} = f(r)^{-1}dr^{2} + r^{2}\left[-dt^{2} + \cosh(t)^{2}d\Omega_{23}^{2}\right] + R_{0}^{2}f(r)d\theta^{2}$$

- 1. The BoN is **time-dependent**
- 2. The dual seems to be at strong coupling (SPT fails)

BEYOND GRAVITY

The T-dual of the BoN at large radius is dual to a **stringy sized** bubble of nothing.

Winding modes become light and break the U(1) momentum symmetry. $\Psi_n \sim \Psi(r) \cos(n heta)$

Only a **full-fledged sigma model** for the BoN can (maybe) help us.

This is difficult because:

1.

The BoN is *time-dependent*
$$ds^2 = f(r)^{-1}dr^2 + r^2[-dt^2 + \cosh(t)^2 d\Omega_{23}^2] + R_0^2 f(r)d\theta^2$$

2. The dual seems to be at strong coupling (SPT fails)

But there is still *****hope*****:

One could imagine stabilizing the BoN with new ingredients → & fall back on cigar CFT ?

BEYOND GRAVITY

The T-dual of the BoN at large radius is dual to a **stringy sized** bubble of nothing.

Winding modes become light and break the U(1) momentum symmetry. $\Psi_n \sim \Psi(r) \cos(n heta)$

Only a **full-fledged sigma model** for the BoN can (maybe) help us.

This is difficult because:

1.

The BoN is *time-dependent*
$$ds^2 = f(r)^{-1}dr^2 + r^2[-dt^2 + \cosh(t)^2 d\Omega_{23}^2] + R_0^2 f(r)d\theta^2$$

2. The dual seems to be at strong coupling (SPT fails)

But there is still *****hope*****:

One could imagine stabilizing the BoN with new ingredients → & fall back on cigar CFT ? Perhaps there is a limit where it looks static [WIP w/ M. Montero & A. Uranga]

BEYOND GRAVITY

The T-dual of the BoN at large radius is dual to a stringy sized bubble of nothing.

Winding modes become light and break the U(1) momentum symmetry. $\Psi_n \sim \Psi(r) \cos(n heta)$

Only a **full-fledged sigma model** for the BoN can (maybe) help us.

This is difficult because:

1.

The BoN is *time-dependent*
$$ds^2 = f(r)^{-1}dr^2 + r^2[-dt^2 + \cosh(t)^2 d\Omega_{23}^2] + R_0^2 f(r)d\theta^2$$

2. The dual seems to be at strong coupling (SPT fails)

But there is still *****hope*****:

One could imagine stabilizing the BoN with new ingredients → & fall back on cigar CFT ?

Perhaps there is a limit where it looks **static** [WIP w/ M. Montero & A. Uranga]

Other ideas?

Matilda Delgado	IFT UAM-CSIC	Geometry, Strings & the Swampland - Gong Show	19/03/2024	28
		τμληκςι		
		ΙΠΑΝΛΞ:		

Matilda Delgado	IFT UAM-CSIC	Geometry, Strings & the Swampland - Gong Show	19/03/2024	29
		EXTRA SLIDES		

Bon & CIGAR CFT

• The cigar CFT describes the worldsheet of the cigar geometry [Witten 91]

$$ds_{cigar}^2 = dr^2 + \tanh(r)^2 d\theta^2$$

It is based on the SI(2,R)/U(1) gauged WZW (coset) model

This also describes a circle that pinches-off

But it is not the same solution as the BoN:

Main difference: the cigar is not Ricci-flat \rightarrow you need to introduce a dilaton background for conformal invariance

 $\phi(r) = 2\ln\cosh r + \text{constant}$

PLUS: the BoN has a time-dependent sphere transverse to this plane..

UNDER T-DUALITY: EXAMPLES

• The Bordism defects of type IIA and IIB on a circle [McNamara, Vafa '19]



UNDER T-DUALITY: EXAMPLES

• The Bordism defects of type IIA and IIB on a circle [McNamara, Vafa '19]



• The T-duality between the Taub-NUT and the NS5-brane [Gregory, Harvey, Moore, `97] [Tong '02]



CLUES IN GRAVITY

After dimensional reduction, the dual object has **finite tension**

Perform the dimensional reduction along the dual circle.

Introduce a codimension 1 source and solve the equations of motion:

$$S_{25} \supset -T \int d^{25}x \sqrt{-g_{24}} e^{a\tilde{\phi}} e^{b\omega} \,\delta(r-r_0)$$

You get: $T \sim O(1) M_{p,25}^{24}$

The dual solution describes an object that is **smeared** along the circle

We know the momentum U(1) of the dual has to be broken

And yet the metric obtained by naive T-duality exhibits this symmetry:

$$d\tilde{s}^{2} = \left(e^{\frac{\phi_{0}}{6}}R_{0}^{2}f(r)\right)^{\frac{1}{12}}\left[f(r)^{-1}dr^{2} + (g_{bon})_{ij}dx^{i}dx^{j}\right] + e^{-\frac{23}{72}\phi_{0}}\left(R_{0}^{2}f(r)\right)^{-\frac{11}{12}}d\tilde{\theta}^{2}$$

$$\tilde{\phi} = \frac{11}{12}\phi_0 - \frac{1}{2}\log\left[f(r)R_0^2\right]$$
.