

ETW brane networks for Calabi-Yau moduli

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Geometry, Strings and the Swampland

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Based on: 2404.xxxxx

2312.16286 with [A. Uranga](#) and [A. Makridou](#)

2203.11240 with [A. Uranga](#), [J. Huertas](#), [M. Delgado](#) and [J. Calderon-Infante](#)

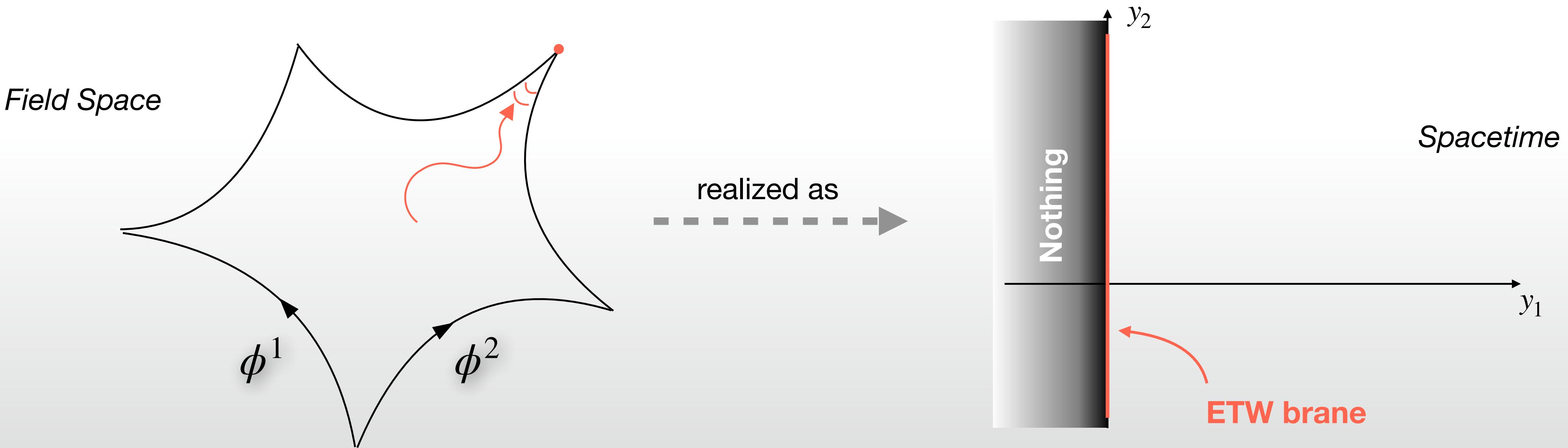
Dynamical Cobordisms to Nothing

They are realizations of spacetime boundaries through specific solutions of the effective action:

$$S = \int d^d x \sqrt{-g} \left\{ \frac{1}{2} R - \frac{1}{2} G_{ij} \partial \phi^i \partial \phi^j - V(\phi^1, \phi^2, \dots) \right\}$$

that:

- are **spacetime dependent**,
- show a **Ricci singularity** in the metric at **finite distance** in spacetime,
- explore **infinite distance** in the field space.

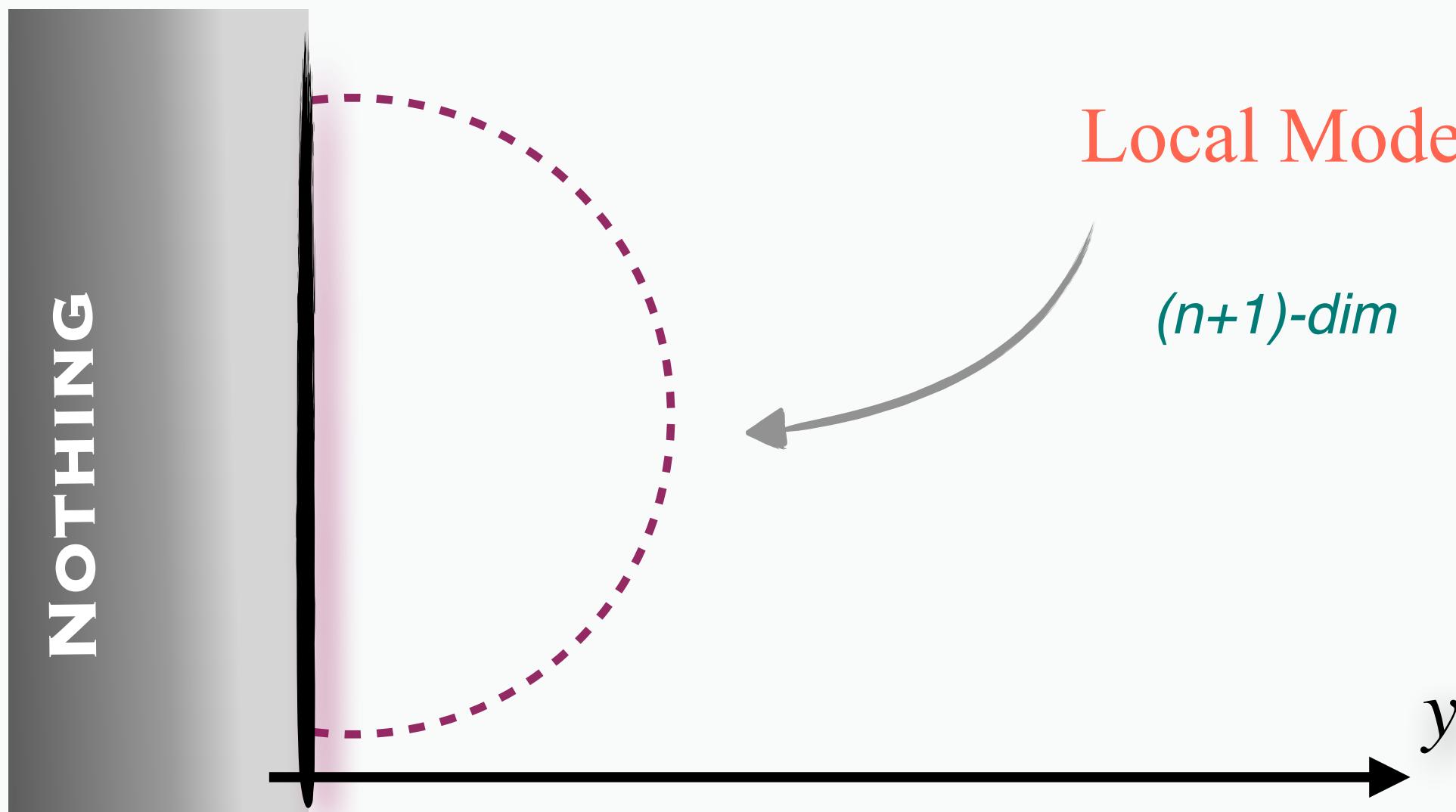


Codimension 1 ETW

[2203.11240-R.A., Calderon-Infante, Delgado, Huertas, Uranga]

We describe codimension 1 boundaries in $(n + 1)$ -dimensional spacetime associated to the action:

$$S = \int d^{n+1}x \sqrt{-g} \left\{ \frac{1}{2}R - \frac{1}{2}(\partial\varphi)^2 - V(\varphi) \right\}$$



Local Model:

$(n+1)\text{-dim}$

- $\varphi(y) \simeq -\frac{2}{\delta} \log y$
- $ds_{n+1}^2 = e^{-2\sigma(y)} ds_n^2 + dy^2$
with $\sigma(y) \simeq \pm \frac{4}{(n-1)\delta^2} \log y$

The class of potentials producing these solutions is: $V(\phi) \simeq -cae^{\delta\phi}$

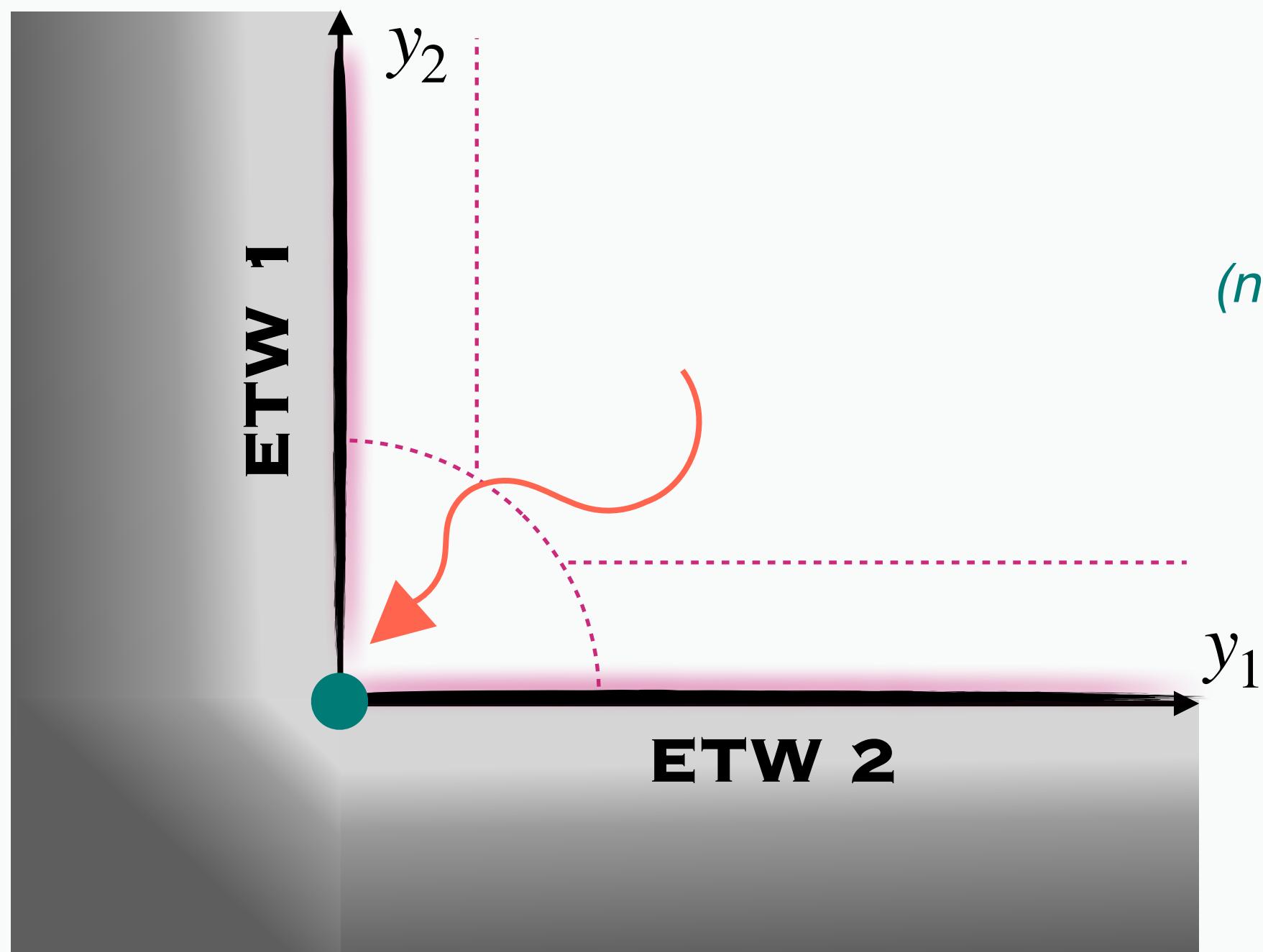
Intersecting ETWs

[2312.16286 - R.A., Makridou, Uranga]

Andriana's talk

We describe the intersection of multiple codimension 1 boundaries in spacetime associated to the action:

$$S = \int d^{n+2}x \sqrt{-g} \left\{ \frac{1}{2}R - \frac{1}{2}(\partial\varphi_1)^2 - \frac{1}{2}(\partial\varphi_2)^2 - \frac{\alpha}{2}\partial\varphi_1\partial\varphi_2 - V(\varphi_1, \varphi_2) \right\}$$



(n+2)-dim

- $\varphi_1(y_1) = -\frac{2}{\delta_1} \log y_1$
 - $\varphi_2(y_2) = -\frac{2}{\delta_2} \log y_2$
 - $ds_{n+2}^2 = e^{-2\sigma_1(y_1)-2\sigma_2(y_2)} ds_n^2 + e^{-2\sigma_2(y_2)} dy_1^2 + e^{-2\sigma_1(y_1)} dy_2^2$
- with $\sigma_1(y_1) = \pm \frac{4}{n\delta_1^2} \log y_1$ and $\sigma_2(y_2) = \pm \frac{4}{n\delta_2^2} \log y_2$

The class of potentials producing these solutions is: $V(\phi) \simeq -c_1 v_1 e^{\delta_1 \varphi_1} e^{a_2 \delta_2 \varphi_2} - c_2 v_2 e^{a_1 \delta_1 \varphi_1} e^{\delta_2 \varphi_2}$

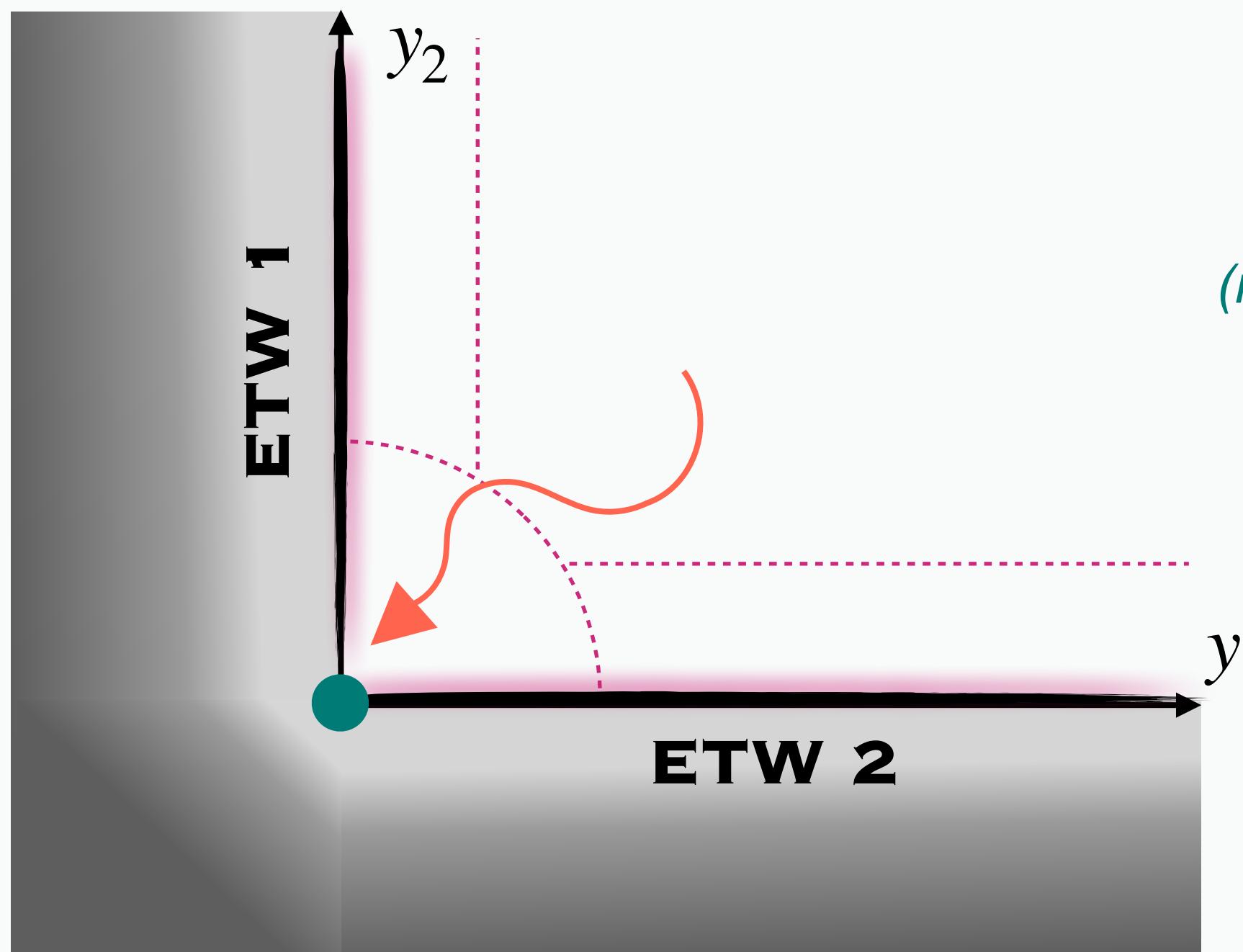
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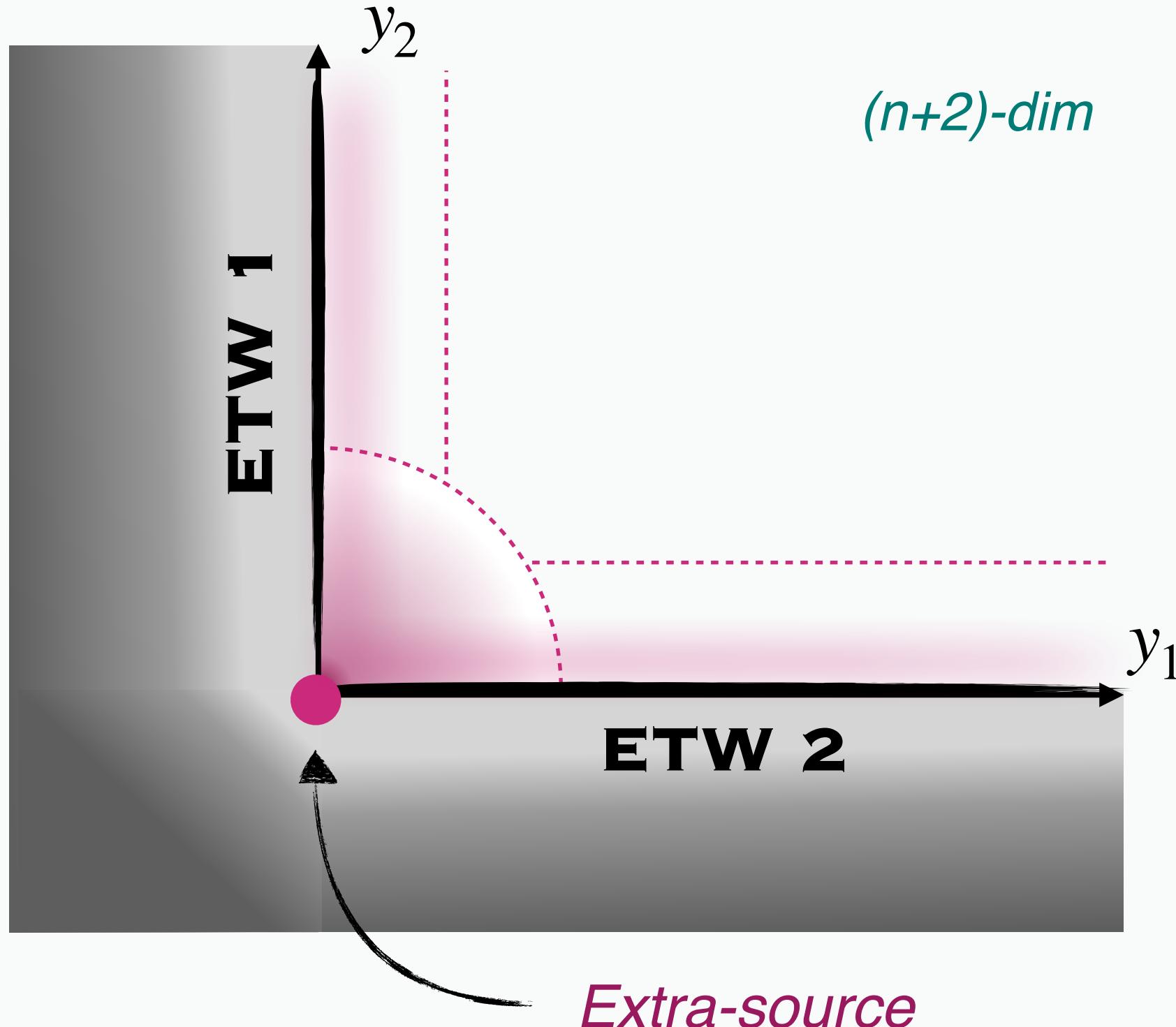
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Beyond the conformal flatness

[2404.xxxxx-R.A.]



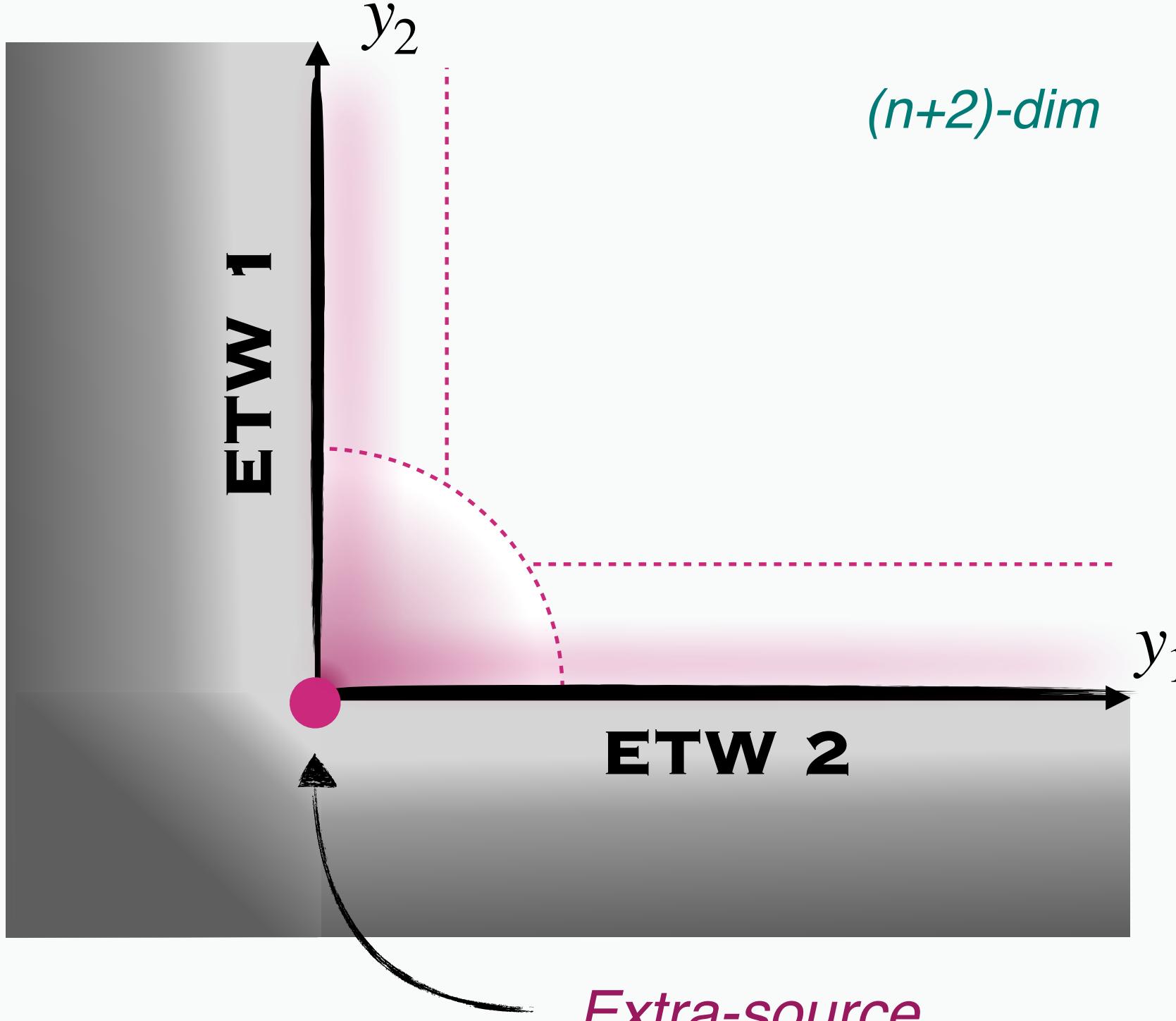
Solutions:

- $\varphi_1(y_1) = -\sqrt{a_1 n(2 - a_1 - a_1 n)} \log y_1 = -\frac{2(1 - a_1 n)}{\delta_1} \log y_1$
- $\varphi_2(y_2) = -\sqrt{a_2 n} \log y_2 = -\frac{2}{\delta_2} \log y_2$
- $ds_{n+2}^2 = y_1^{2(1-a_1n)} y_2^{2a_2} \left[y_1^{4a_1n-2} ds_n^2 + y_1^{-2(1-a_1n)} dy_1^2 - y_2^{-2a_2n} dy_2^2 \right]$

Potential class: $V(\phi) \simeq c \cdot v \cdot e^{\delta_1 \varphi_1} e^{\delta_2 \varphi_2}$

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Potential class: $V(\phi) \simeq c \cdot v \cdot e^{\delta_1 \varphi_1} e^{\delta_2 \varphi_2}$

$$\varphi_1 = \psi_1 - (1 - a_1 n) \sqrt{\frac{a_2}{a_1(2 - a_1 - a_1 n) - a_2(a_1 n - 1)^2}} \psi_2$$

$$\varphi_2 = \left(1 + \frac{a_2(a_1 n - 1)^2}{a_1(-2 + a_1 + a_1 n)} \right)^{-1/2} \psi_2$$

$$S = \int d^{n+2}x \sqrt{-g} \left\{ \frac{1}{2}R - \frac{1}{2} (\partial\psi_1)^2 - \frac{1}{2} (\partial\psi_2)^2 - V(\psi_1, \psi_2) \right\}$$

Complex Structure Moduli Space of CY₄

[Math: Cattani, Kaplan, Schmid '86; Deligne '91; Kerr, Pearlstein, Robles '19 ...]

[Phys: Grimm, Palti, Valenzuela '18; Grimm, Li, Palti '19; Grimm, Li, Valenzuela '20...]

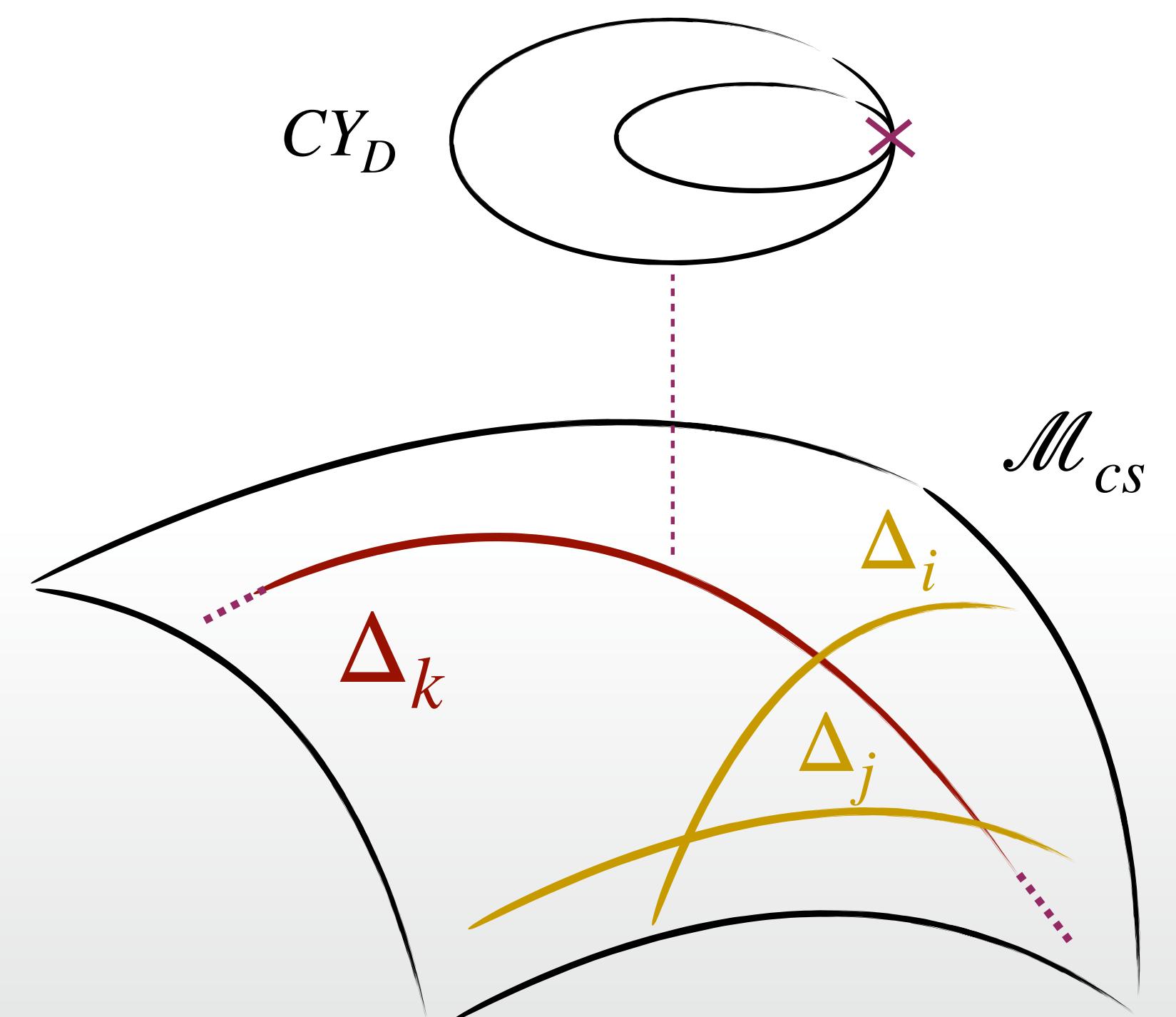
Let \mathcal{M}_{cs} be the complex structure moduli space of CY₄:

- It has complex dimension $h^{3,1} = 2$
- It is not smooth nor compact



Discriminant locus

$$\Delta = \bigcup_k \Delta_k$$



Complex Structure Moduli Space of CY_4

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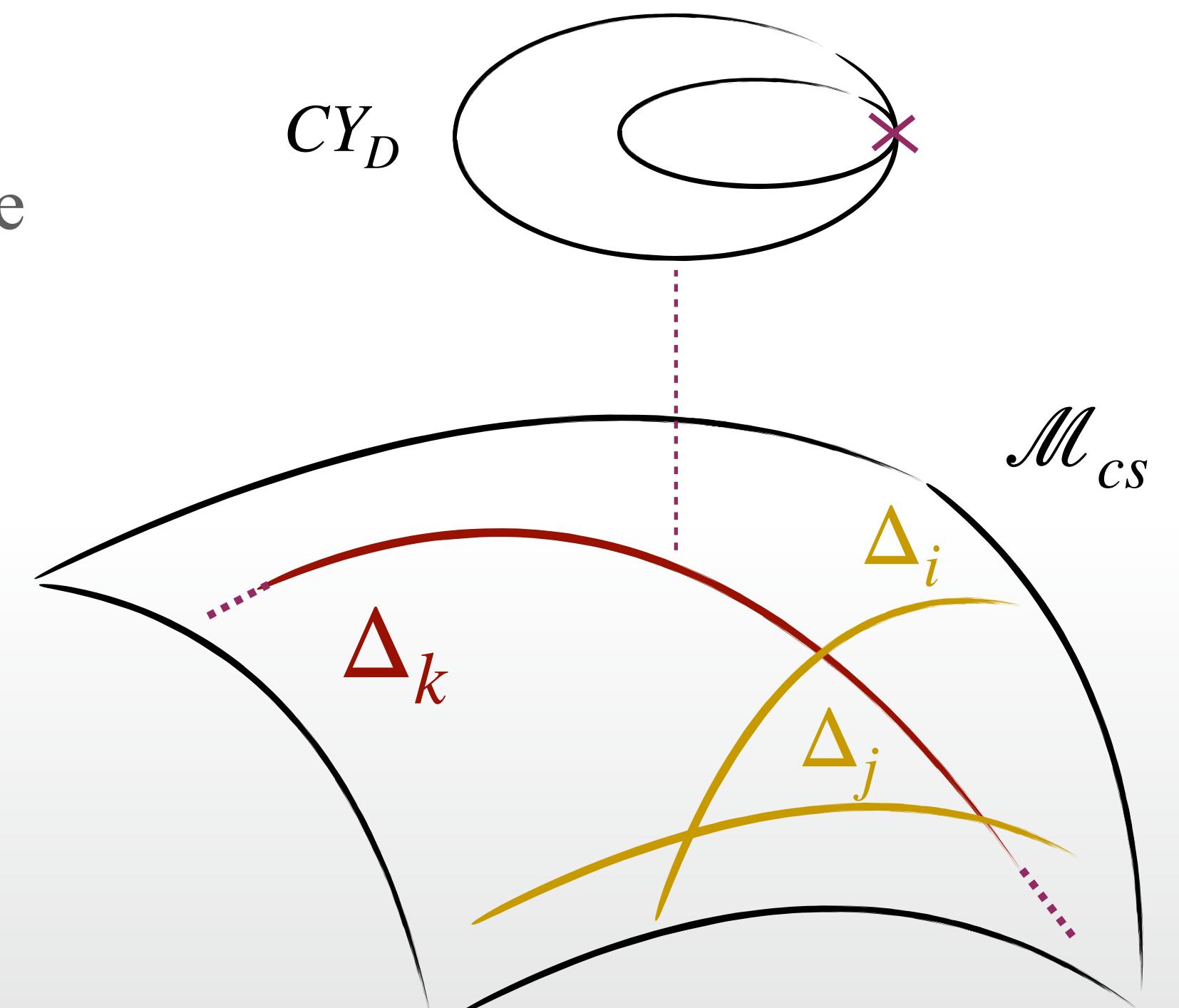
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$$\Delta = \bigcup_k \Delta_k$$

- Using the mixed Hodge-Deligne theory we get a classification of the singularities and their enhancements
- Singularities of type II, III, IV, V for CY_4 are located at infinite distance



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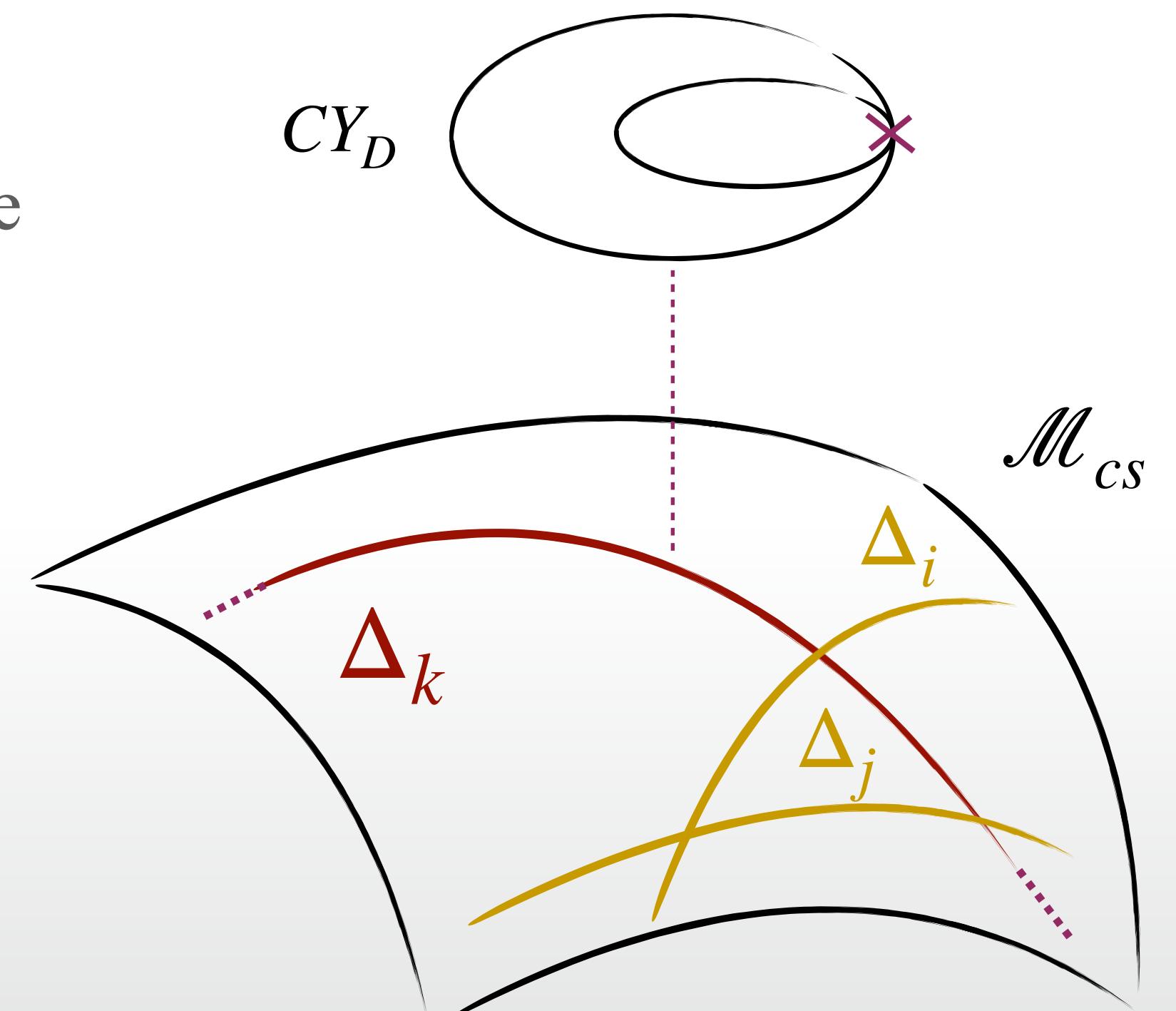


Discriminant locus

$$\Delta = \bigcup_k \Delta_k$$

- Using the mixed Hodge-Deligne theory we get a classification of the singularities and their enhancements
- Singularities of type II, III, IV, V for CY₄ are located at infinite distance
- Turning on a $G_4 \in H_p^4(Y_4, \mathbb{C})$ we generate the 3d potential:

$$V = \frac{1}{\mathcal{V}_4^3} \left(\int_{Y_4} G_4 \wedge \star \bar{G}_4 - \int_{Y_4} G_4 \wedge G_4 \right)$$



Complex Structure Moduli Space of CY₄

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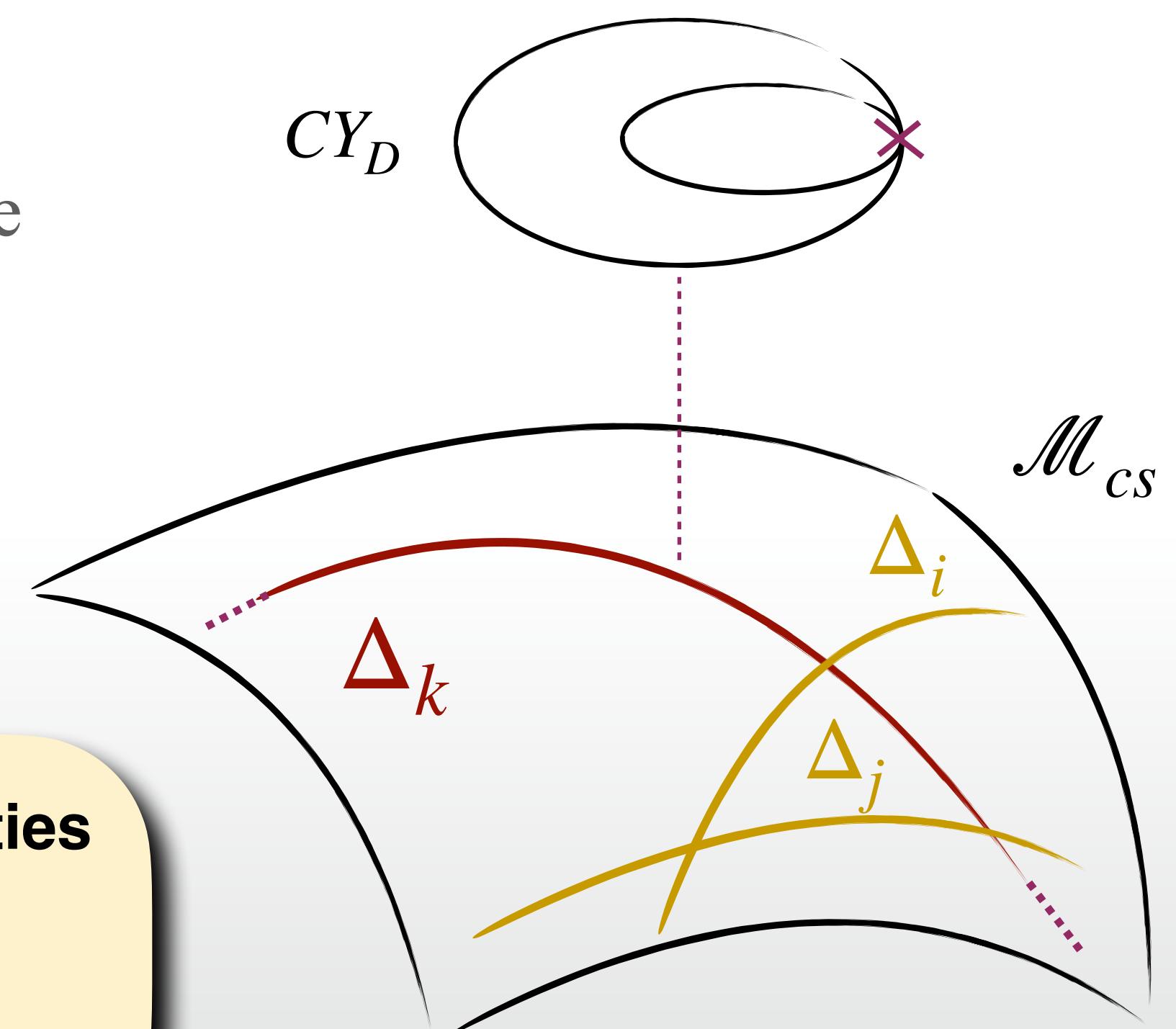
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Effective action for complex structure moduli near the network of singularities

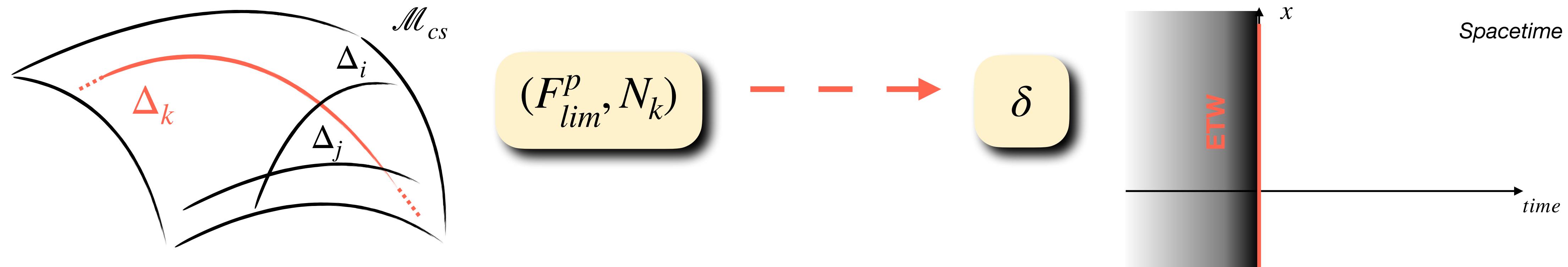
$$S = \int d^3x \sqrt{-g} \left\{ \frac{1}{2} R - G_{kk} \partial \phi^k \partial \phi^k + V(\phi^k) \right\}$$



ETW networks for Δ

Applying Dynamical Cobordism thecniques in the resulting three-dimensional effective action:

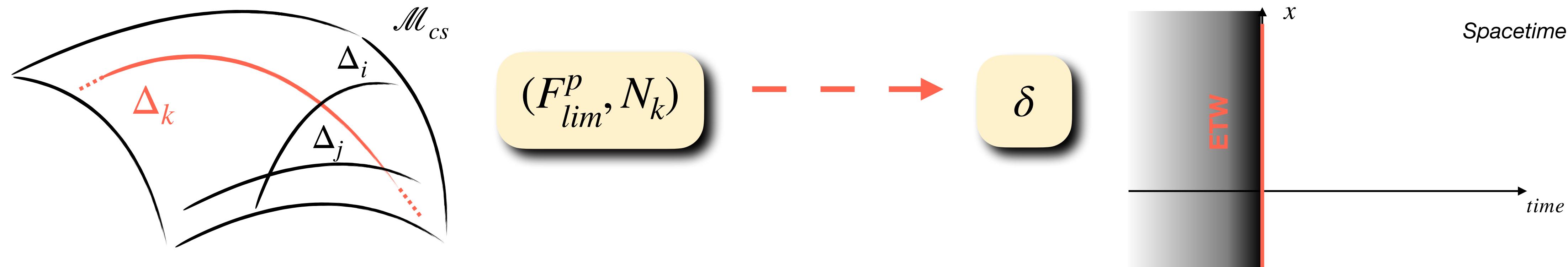
- We can associate to each singular divisor $\Delta_k \in \Delta$ a specific ETW brane in spacetime



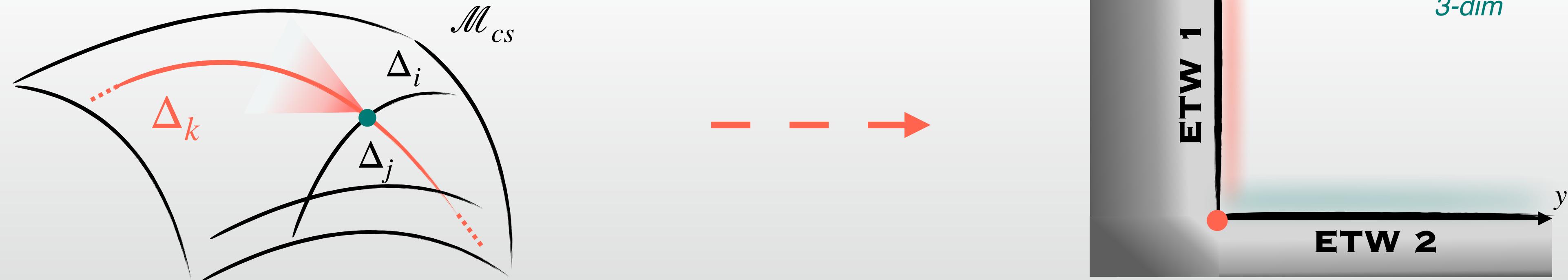
ETW networks for Δ

Applying Dynamical Cobordism thecniques in the resulting three-dimensional effective action:

- We can associate to each **singular divisor** $\Delta_k \in \Delta$ a specific ETW brane in spacetime



- We can provide a spacetime realization for the **Growth Sectors** associated to each enhancement as an **intersection of two ETW branes**: where the first brane represents the first singular divisor and the second brane represents the enhanced singularity



Thank you for your attention!