New Insights on Black Hole Evaporation and the Dark Matter Fraction Composed of Primordial Black Holes



Luis Anchordogui CUNY



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 $f = \sum \left[Z_{i} Z_{j} - \frac{2}{2} \left(X_{i} X_{j} + Y_{i} Y_{j} \right) \right] / \sqrt{3} + \frac{1}{2} X_{i}$ Le square h=0_ Z out of place L 5=01,00 sSG stable 1=1 · transport 4

arXiv:2401.09087 arXiv:2403.xxxx

1. -Hidden agenda

- .primordial black holes as dark matter .the dark dimension: a quick glance
- 2. Hawking's semiclassical approximation
 - .Schwarzschild black holes
 - .Near extremal black holes
- 3. Quantum effects r Memory burdem
- 4. Summary:
 - Limits on all-dark-matters interpretation of 5D primordial black holes





Primordial black holes as dark matter

- Scaling of cosmological energy density with time $= \rho \sim M_p^2/t^2$
- Required density for M_{BH} to collapse within its Schwarzschild radius = $ho~~M_{
 m p}~/M_{BH}^2$

$$M_{BH} \sim t M_p^2 \sim 10^{15} (t/10^{-23} s) s$$

PBHs would initially have around the cosmological horizon mass

Speculation BH could be formed from collapse of large amplitude fluctuations in early universe

Zeldovich 1967; Hawking 1971

What's MBH?





• $M_{BH} \sim M_p \sim 10^{-5}$ g r formed @ Planck time 10^{-43} s

- $M_{BH} \sim M_{\odot} \sim 10^{33}$ g r formed (a QCD epoch 10^{-5} s)
- $M_{BH} \sim 10^5 M_{\odot} \sim 10^{38} g$ r formed (a) t ~ 18



Behave like typical CDM particles on cosmological scales

PBHs could span an enormous mass range

Mass spectrum is yet to be shaped

A PBH all-dark-matter interpretation is severely constrained by observations







Green, Kavanagh 2020





Swampland distance conjeture . Infinite distance limits $\phi \rightarrow \infty$ in the field space of massless scalars are accompanied by an infinite tower of exponentially light states $\mathbf{m} \sim e^{-\alpha\phi} - \alpha \sim \mathcal{O}(1)$ distance and masses are measured in Planck units Ooguri, Vafa 2006

- Kaluza-Klein tower decompactification of extra dimensions $\phi = \ln R - m \sim 1/R$
- Smallness of some physical parameters might signal that we live in asymptotic corner of field space
- , Such parameters can be scales of dark energy and neutrino masses





The Dark Dimension: a quick glance 1.- AdS distance conjecture $\phi = -\ln |A|$ Lüst, Palti, Vafa 2019

- 2.- Extension to dS.
- could help elucidate radiative stability of cosmological hierarchy $\Lambda \sim 10^{-120} \mathrm{M}_{\mathrm{p}}^{4}$
- - because it connects size of compact space R_{\perp} to dark energy scale A $\mathbf{R}_{1} \sim \lambda \Lambda^{-\alpha}$
- Null results on deviations from Newton's law $rac{R} < 30 \ \mu m \Rightarrow \alpha = 1/4$
- Astrophysical constraints dark dimension
 - $\mathbf{R}_{\perp} \sim 1 \, \mathbf{to} \, 10 \, \mu \mathbf{n}$ $10^{-3} \leq \lambda \leq 10^{-2}$ Montero, Vafa, Valenzuela 2022
- $\alpha = 1/4$ consistent with string computations
 - LAA, Antoniadis, Lüst, Lüst 2023



1.- Species scale where gravity becomes strong $M_* \sim m^{1/3} M_p^{2/3} \sim 10^9 \text{ GeV}$

2.- Interesting dark matter candidates

5D PBHs

KK graviton tower

(explicit realization of dynamical dark matter)

Dvali 2007

LAA, Antoniadis, Lüst, 2022

Gonzalo, Montero, Obied, Vafa 2022

Dienes, Thomas 2012

- For point-like lenses.

 - and effects of wave optics suppress magnification,

Stunning coincidence

size of the dark dimension \sim wavelength of visible light

• Schwarzschild radius of 5D black holes - well below wavelength of light

this is precisely critical length where geometric optics breaks down

obstructing sensitivity to 5D PBH microlensing signals!







- 1.- Gravity propagates in bulk spacetime while SM fields lie on a brane
- 2.- Rely on probe brane approximation
- Ensures only effect of brane field is to bind black hole to brane
- Adequate approximation provided MBH »brane tension
 - presumably of order of but smaller than M_*
- 3.- Black hole can be treated as flat d-dimensional object
- Assumption valid for $r_s \sim (M_{BH} / M_*)^{1/(d-3)} M_*^{-1} \ll R_{\perp}$

Assumptions

Tangherlini 1963; Myers, Perry 1986







Schwarzschild black holes

$ds^2 = - f(r) dt^2 +$

• $f(r) = 1 - r_s/r$

. $d\Omega_{d-2} = interval of spherical solid angle in dimension d-2$

$$f^{-1}(r) dr^2 + d\Omega_{d-2}^2$$



Hawking's semiclassical approximation

BH emits thermal radiation as it were a blackbody





Particle emission rate



Hawking's lifetime



$$(M_{BH}/M_{*})^{-1/(d-3)}$$

$$T_{\rm H} \sim T_{\rm H} \sim M_* (M_{\rm BH}/M_*)^{-1/(d-3)}$$

 $T_{\rm H} \sim S_{\rm BH} r_{\rm s} \sim (M_{\rm BH}/M_*)^{(d-1)/(d-3)}$



Black holes radiate mainly on the brane

by geometric factor $(r_s/R_{\perp})^{d-4}$ relative to brane fields

that from single brane field

KK modes are excitations in full transverse space and so their overlap with small (higher dimensional) black holes is suppressed

This geometric suppression precisely compensates for the enormous number of modes and total KK contribution is only of same order as

Emparan, Horowitz, Myers 2000









Recent numerical analysis

Ireland, Profumo, Scharnhorst 2023



with initial total energy between (Q, Q + dQ)

$$\frac{d\dot{N}_i}{dQ} = \frac{\sigma_s}{8\pi^2} Q^2 \left[\exp\left(\frac{Q}{T_{\rm H}}\right) - (-1)^{2s} \right]^{-1}$$

. Average total emission rate for particle species i is then

$$\Gamma_i \equiv \dot{N}_i = f \frac{\Gamma_s}{32 \pi^3} \frac{(d-1)^{(d-1)/(d-3)}}{2^{2/(d-3)}} \frac{(d-3)}{\Gamma(3)} \zeta(3) T_{\rm H}$$

- . greybody factor taken as dimensionless contant $\Gamma_{s} = \sigma_{s} / A_{4}$
 - normalized to the horizon surface

seen by the SM fields $\Gamma_{s=0} = 1$, $\Gamma_{s=1/2} \approx 2/3$, $\Gamma_{s=1} \approx 1/4$

• f = 1 (f = 3/4) for bosons (fermions)

• Emission rate per degree of particle freedom i of particles of spin s

e area
$$A_4 = 4\pi \left(\frac{d-1}{2}\right)^{2/(d-3)} \frac{d-1}{d-3} r_s^2$$

LAA, Goldberg 2002





$T_{\rm H} \sim (M_{\rm BH} / 10^{16} g)^{-1} \, {\rm MeV}$

$T_{\rm H} \sim 13.8 \,(M_{\rm BH} \,/\,10^{14.7} \,g)^3 \,Gyr$



4D versus 5 D

$T_{\rm H} \sim (M_{\rm BH} / 10^{12} \text{ g})^{-1/2} \text{ MeV}$





LAA, Antoniadis, Lüst 2022





PBH all-dark-matter interpretation

• Difference in number density

Compensated by black hole mass

Same PBH density

$\frac{4D}{n_{PBH}(T_H)} < \frac{5D}{n_{PBH}(T_H)}$

$M_{BH}^{4D}(T_{H}) > M_{BH}^{5D}(T_{H})$

ho_{PBH}^{4D} (T_H) ~ ho_{PBH}^{5D} (T_H)













$C = (M_{RH}^2 - Q^2 M_{*}^{d-2})1/2$ $M_{BH}^2 - Q^2 M_*^{d-2} > 0 - two horizons$ $M_{RH}^2 - Q^2 M_{*}^{d-2} = 0 - extremal black hole$ $M_{BH} - Q^2 M_*^{d-2} < 0 - naked singularity$

Reissner-Nordström black holes $ds^2 = u(r) dt^2 - u^{-1}(r) dr^2 - d\Omega_{d-2}^2$ unit (d-2) sphere $u(r) = 1 - 2M_{BH} / (M_{*}^{d-2} r^{d-3}) + Q^{2} / (M_{*}^{d-2} r^{2d-6})$

Chamblin, Emparan, Johnson, Myers 1999



Near-extremal black holes

- For Q M^(d-2)/2/M_{BH} $\ll 1 C \sim M_{BH}$
 - leading to non-extremal relation between $C \sim M_* S_{RH}^{(d-3)} / (d-2)$
- For near-extremal case \blacksquare M_{BH} ~ Q M^(d-2)/2 expanding square root it follows that leading term cancels and sub-leading term gives

C ~

$$M_* \beta^{1/2} S_{BH}^{(d/2-2)} / (d-2)$$

- β r order-one parameter controling difference between M_{BH} and Q
 - Cribiori, Dierigl, Gnecchi, Lüst, Scalisi 2022; Basile, Cribiori, Lüst, Montella 2024





• Near-extremal temperature

• Particle emission rate

Evaporation rate of near-extremal black holes

would be suppressed by a factor of $(\beta / S_{BH})^{1/2}$

with respect to Schwarzschild black holes of same mass

$T_{ne} = T_{H} (\beta / S_{BH})^{1/2}$

$\Gamma_{ne} \sim T_{ne} = (\beta / S_{BH})^{1/2} \Gamma_{H}$

LAA, Antoniadis, Lüst 2024



- For $10^5 g = T_H \sim 4 \text{ GeV } BUT T_{ne} \sim 10^{-5} \beta^{1/2} \text{ eV}$
- For $10^5 g \tau_H \sim 4 \times 10^{-5} yr$
- Order of magnitude estimate:
 - If there were 5D primordial neBH in nature
 - - $10^5 < M_{BH}/g < 10^{21}$
- $\hat{\mathbf{C}} = \mathbf{C} / M_{BH} = (\beta / S)^{1/2}$

. For near-extremal black hole of the same mass $rac{}$ The ~ 20 / $\beta^{1/2}$ Gyr

a PBH all-DM interpretation would be possible in mass range

quantifies near-extremality is very small because of large entropy

LAA, Antoniadis, Lüst 2024





Self-similarity

- 1.- Semiclassical approximation relies on assumption of self-similarity. black hole gradually shrinks in size while maintaining the standard semi-classical relations between its parameters: rs, MBH, TH 2.- Would semiclassical approximation hold throughout entirety of the black hole lifetime? 3.- Assuming self-similarity $-S_{BH} \sim S_{emission}$ at τ_{Page} $M_{BH} \rightarrow M_{BH}/2$ $r_s \rightarrow r_s/2$ $S_{BH} \rightarrow S_{BH}/4$
 - for d=4

Page 2013





Quantum decay rate: effect of memory burden

- 1.- Self-similarity implies that at each stage radiation is thermal and information is maintained internally
- 2.- However remaining black hole has only 1/4 of its initial entropy and so much less information storage capacity
- 3.- Analysis of prototype models shows after τ_{Page} system gets effectively stabilized and gains much longer lifetime formula
- 4.- Assume evaporation is slowed down by further n powers of entropy S_{BH} $\tau_{\mathrm{H}}^{\{n\}} \sim \mathbf{S}_{\mathrm{BH}}^{1+n} \mathbf{r}_{\mathrm{s}}$
 - Dvali 2018; Dvali, Eisemann, Michel, Zell 2020; Alexandre, Dvali, Koutsangelas 2024







- Hawking decay rate of Schwarzschild black holes
- additional $1 / S_{BH}^n$ suppression
- Compared to evaporation of Schwarzschild black holes -
- Quantum decay rate of near-extremal black holes

$\Gamma_{\rm H}^{\{0\}} \sim \Gamma_{\rm H} \sim M_{*} \, {\bf S}_{\rm BH}^{-1/(d-2)} \sim M_{*} \, (M_{\rm BH}/M_{*})^{-1/(d-3)}$

• Compared to Hawking decay rate - quantum decay rate has

 $\Gamma_{\rm H}^{\{n\}} \sim T_{\rm H} / S_{\rm BH}^{n} \sim M_{*} S_{\rm BH}^{(-1)} = nd + 2n/(d-3) \sim M_{*} (M_{\rm BH} / M_{*}) (-1 - nd + 2n/(d-3))$ near-extremal black holes have additional $(\beta / S_{BH})^{1/2}$ suppression $\Gamma_{ne}^{\{0\}} \sim \beta^{1/2} T_{H} / S_{BH}^{1/2} \sim \beta^{1/2} M_{*} S_{BH}^{-d} / (2d - 4) \sim \beta^{1/2} M_{*} (M_{BH} / M_{*})^{-d} / (2d - 6)$

 $\Gamma_{ne}^{\{n\}} \sim \beta^{1/2} T_{H} / S_{BH}^{n+1/2} \sim \beta^{1/2} M_{*} S_{BH}^{(-d/2 - nd + 2n)(d-2)} \sim \beta^{1/2} M_{*} (M_{BH}/M_{*})^{(-d/2 - nd + 2n)(d-3)}$





4D versus 5 D





LAA, Antoniadis, Lüst to appear







Assuming only two phases in evaporation process

 $\tau_{\text{Page}} \sim 180 (M_{\text{BH}} / 10^4 \text{g})^2 \text{s}$ $\tau_{\text{H}}^{\{1\}} \sim 4 \times 10^{17} (M_{\text{BH}} / 5 \text{g})^2 \text{s}$

How would memory burden impact the mass range of PBH-DM interpretation?

LAA, Antoniadis, Lüst to appear









- 1.- 5D black holes are bigger, colder, and longer-lived than usual 4D black holes of same mass
- 2.- Adopting Hawking's semiclassical approximation PBH all dark matter interpretation would be possible if
 - Schwarzschild black holes

Near extremal black holes

3. Memory Burden - stay tuned

Canclusians

- $10^{15} < M_{BH}/g < 10^{21}$
- $10^5 < M_{BH}/g < 10^{21}$





Extra slides

4D versus 5D

$$\tau_H \sim 14 \left(\frac{M_{\rm BH}}{10^{14.5} \text{ g}}\right)^3 \left(\frac{4}{\sum_i c_i(T_H) \tilde{f} \Gamma_s}\right) \quad \text{Gyr} \quad \tau_H \sim 14 \left(\frac{M_{\rm BH}}{10^{12} \text{ g}}\right)^2 \left(\frac{4}{\sum_i c_i(T_H) \tilde{f} \Gamma_s}\right)$$

• c_i (T_H) - counts number of internal degrees of freedom of particle species i of mass m_i satisfying $m_i \ll T_H$ • f = 1 (f = 7/8) for bosons (fermions)

LAA, Antoniadis, Lüst 2022



