

The tadpole conjecture at large complex-structure

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Geometry, Strings and the Swampland

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This talk is based on ::

- *The tadpole conjecture at large complex-structure*
E. Plauschinn
arXiv:2109.00029
- *Moduli Stabilization in Asymptotic Flux Compactifications*
T. Grimm, D. van de Heisteeg, E. Plauschinn
arXiv:2110.05511

Some motivation.

Compactifications of string theory give rise to an abundance of lower-dimensional theories — the **string theory landscape**.

Famous estimates for its **size** are 10^{500} , 10^{930} , 10^{1500} , 10^{272000} .

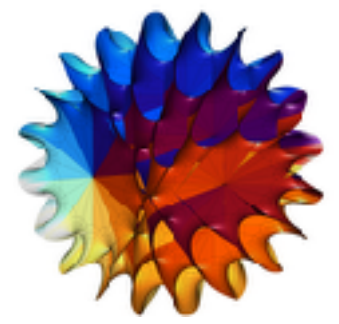
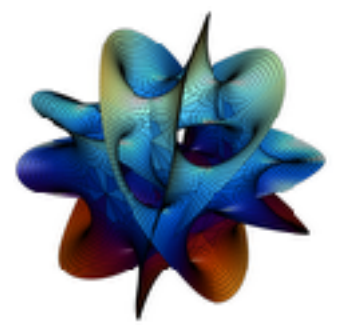
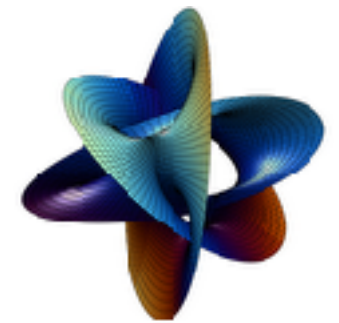
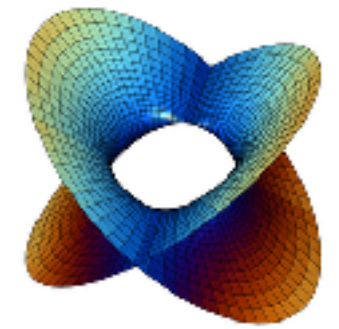
Bousso, Polchinski — 2000

Schellekens — 2016

Lerche, Lüst, Schellekens — 1987

Taylor, Wang — 2015

But one is often only interested in four-dimensional theories with no or few **massless scalar fields**.

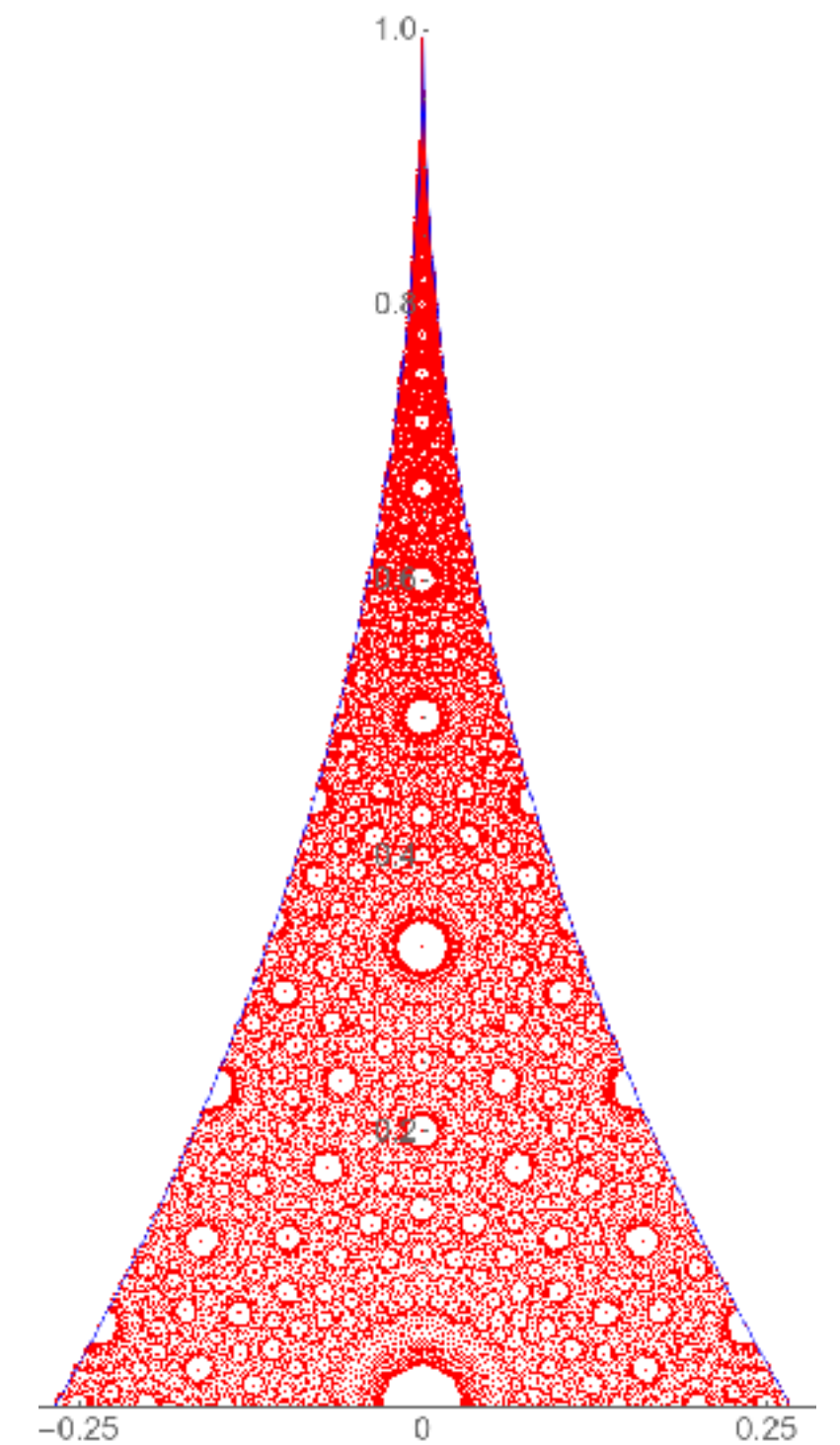


In type IIB orientifold compactifications on Calabi-Yau three-folds, **fluxes** generate a mass-term for the complex-structure and axio-dilaton moduli.

An underlying assumption of the KKLT and Large-Volume scenarios is that **all** of these moduli **can be stabilized** in a suitable regime.

This assumption can fail.

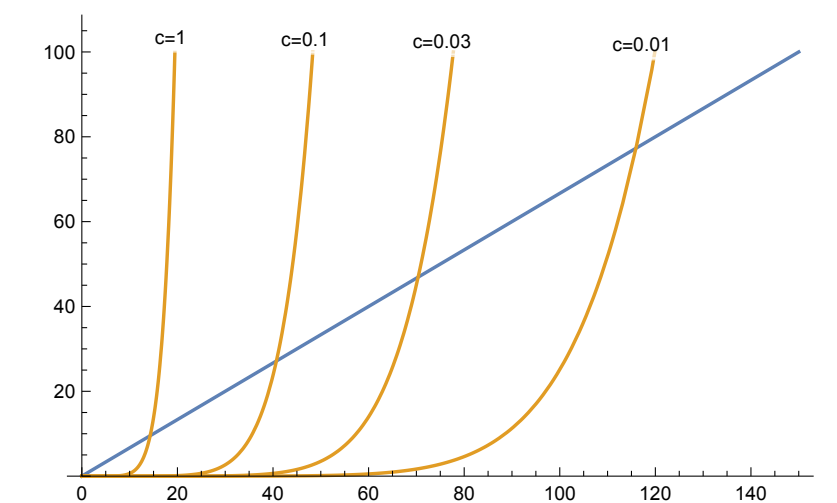
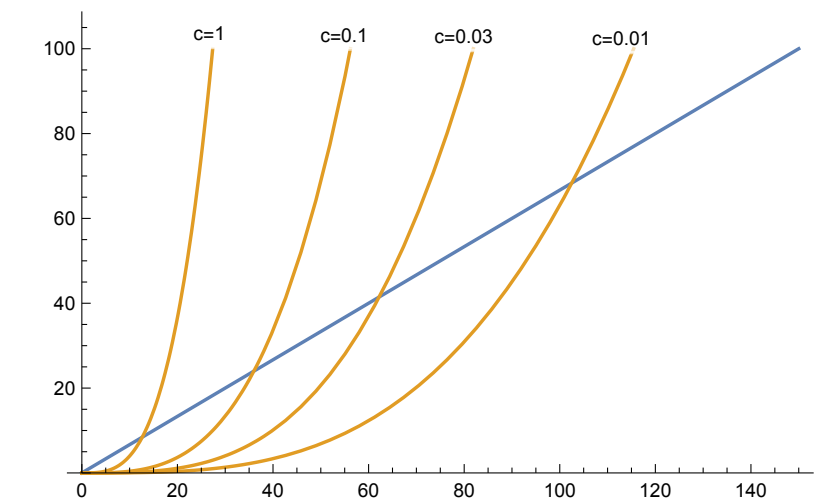
Bena, Dudas, Graña, Lüst — 2018
Betzler, EP — 2019
Braun, Valandro — 2020



The **tadpole conjecture** states that for a large number of complex-structure moduli, not all of them can be stabilized by fluxes.

Bena, Blåbäck, Graña, Lüst — 2020

- Goal** for this talk ::
- Explain why it is difficult to prove this conjecture,
 - show that in the large complex-structure limit generically the conjecture is satisfied, and
 - discuss moduli stabilization in asymptotic regions.



1. motivation
2. the tadpole conjecture
3. boundary behaviour
4. large complex-structure
5. summary
6. asymptotic hodge theory
7. summary

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the tadpole conjecture

Explanation of the tadpole conjecture.

D-branes and O-planes are charged under Ramond-Ramond gauge potentials and therefore contribute to Bianchi identities. The integrated expressions are the **tadpole cancellation conditions**.

In **F-theory**, the D3-brane tadpole equation reads (with χ the Euler number of the four-fold)

$$N_{\text{D3}} + \frac{N_{\text{flux}}}{2} = \frac{\chi}{24} .$$

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The **tadpole conjecture** states that in the large $h^{3,1}$ -limit, the flux number satisfies

$$\frac{N_{\text{flux}}}{2} > \frac{h^{3,1}}{3} .$$

Implications ::

- If true, the tadpole conjecture implies that the landscape of theories with no massless scalar fields is smaller than expected.

Comments ::

- The conjecture also applies to type IIB orientifold compactifications.
- In a number of examples the tadpole conjecture is satisfied, even for

$$\frac{N_{\text{flux}}}{2} > 0.44 \times n_{\text{mod}} .$$

Bena, Blåbäck, Graña, Lüst — 2020 & 2021

- A scenario which violates the tadpole conjecture has been proposed by Marchesano, Prieto & Wiesner.

Marchesano, Prieto, Wiesner — 2021

see also :: Lüst — 2021
Grimm, v.d. Heisteeg, EP — 2021

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Estimating the behaviour of N_{flux} when approaching a boundary.

Consider type IIB orientifold compactification on Calabi-Yau three-folds with NS-NS and R-R fluxes ::

- The global minimum of the scalar potential corresponds to the **self-duality condition**

$$\begin{aligned}\star G_3 &= i G_3, & G_3 &= F_3 - \tau H_3, \\ \tau &= c + i s.\end{aligned}$$

- The Hodge-star matrix M depends on the complex-structure moduli, and the above relation reads in **matrix notation**

$$F_3 = [\eta \mathcal{M} s + \mathbb{1} c] H_3, \quad \eta = \begin{pmatrix} 0 & +\mathbb{1} \\ -\mathbb{1} & 0 \end{pmatrix}.$$

- The **flux number** is defined and satisfies

$$N_{\text{flux}} = F_3^T \eta H_3 \geq 0.$$

The **Hodge-star matrix** $\mathcal{M}$ is a real, symmetric, symplectic matrix. It can therefore be decomposed as

$$\mathcal{M} = U^T \Sigma U ,$$

$$\Sigma = \begin{pmatrix} \lambda & 0 \\ 0 & \lambda^{-1} \end{pmatrix} ,$$

$$\lambda = \text{diag}(\text{eigenvalues } \mathcal{M}) ,$$

$$U \in \text{Sp}(2h^{2,1} + 2, \mathbb{R}) \cap \text{O}(2h^{2,1} + 2, \mathbb{R}) .$$

The **self-duality condition** and the **flux number** (at the minimum) can then be expressed as

$$\tilde{\mathbf{F}}_3 = \begin{pmatrix} \mathbb{1} c & \lambda^{-1} s \\ -\lambda s & \mathbb{1} c \end{pmatrix} \tilde{\mathbf{H}}_3 ,$$

$$\tilde{\mathbf{H}}_3 = U \mathbf{H}_3 , \quad \tilde{\mathbf{F}}_3 = U \mathbf{F}_3 ,$$

$$N_{\text{flux}} = s \sum_I \left[\lambda_I (\tilde{h}^I)^2 + \frac{1}{\lambda_I} (\tilde{h}_I)^2 \right] ,$$

$$\tilde{\mathbf{H}}_3 = \begin{pmatrix} \tilde{h}^I \\ \tilde{h}_I \end{pmatrix} .$$

The **flux number** (at the minimum) can be expressed in the following way

$$N_{\text{flux}} = s \sum_I \left[\lambda_I (\tilde{h}^I)^2 + \frac{1}{\lambda_I} (\tilde{h}_I)^2 \right] .$$

Boundary behaviour of the flux number (for $\tilde{h}^I \neq 0$) ::

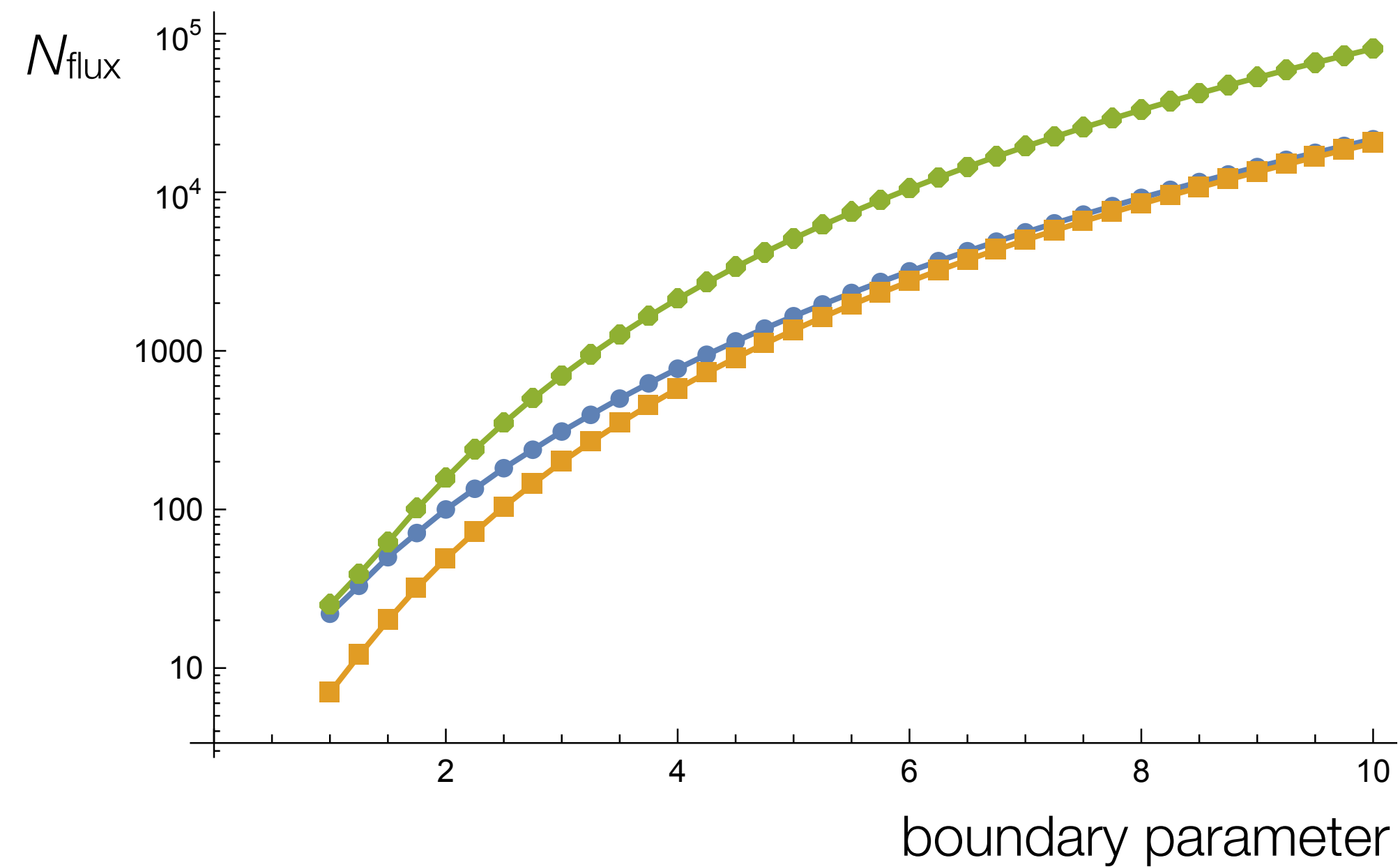
- At a boundary in **complex-structure** moduli space, the matrix M is expected to degenerate. This implies

$$\lambda_I \rightarrow \infty \quad \longrightarrow \quad N_{\text{flux}} \rightarrow \infty .$$

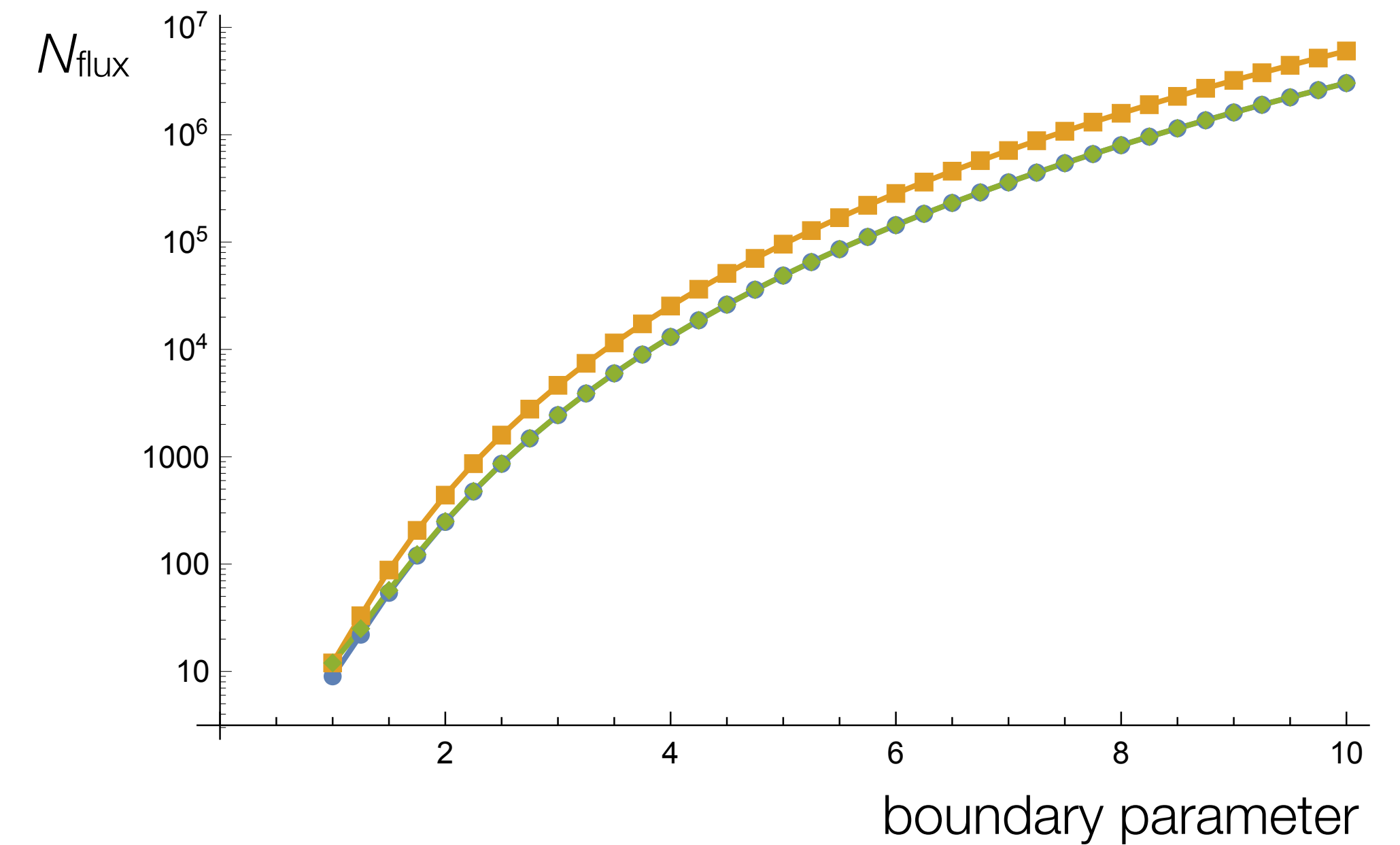
- For the boundary in **axio-dilaton** moduli space one finds

$$s \rightarrow \infty \quad \longrightarrow \quad N_{\text{flux}} \rightarrow \infty .$$

For simple examples and generic configurations, the **flux number** grows significantly near a boundary ::



$h^{2,1}=2$ (LCS)



$h^{2,1}=3$ (LCS)

Summary ::

- Near a boundary in moduli space, the flux number N_{flux} generically diverges.
(The behaviour in non-generic situations is not derived here.)

Implications ::

- The smallest values for N_{flux} will be found in the interior of moduli space.
- It can thus be challenging to prove the tadpole conjecture or to find computationally-controlled counter-examples.

Comment ::

- The same behaviour is found using asymptotic Hodge theory.

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The tadpole conjecture in the large complex-structure regime.

Consider the **large complex-structure regime**, for which the moduli space is described by the prepotential

$$\mathcal{F} = -\frac{1}{3!} \frac{\kappa_{ijk} X^i X^j X^k}{X^0},$$

$$z^i = X^i / X^0 = u^i + i v^i, \\ i, j, k = 1, \dots, h^{2,1}.$$

Using the self-duality condition, the **flux number** (at the minimum) can be expressed as

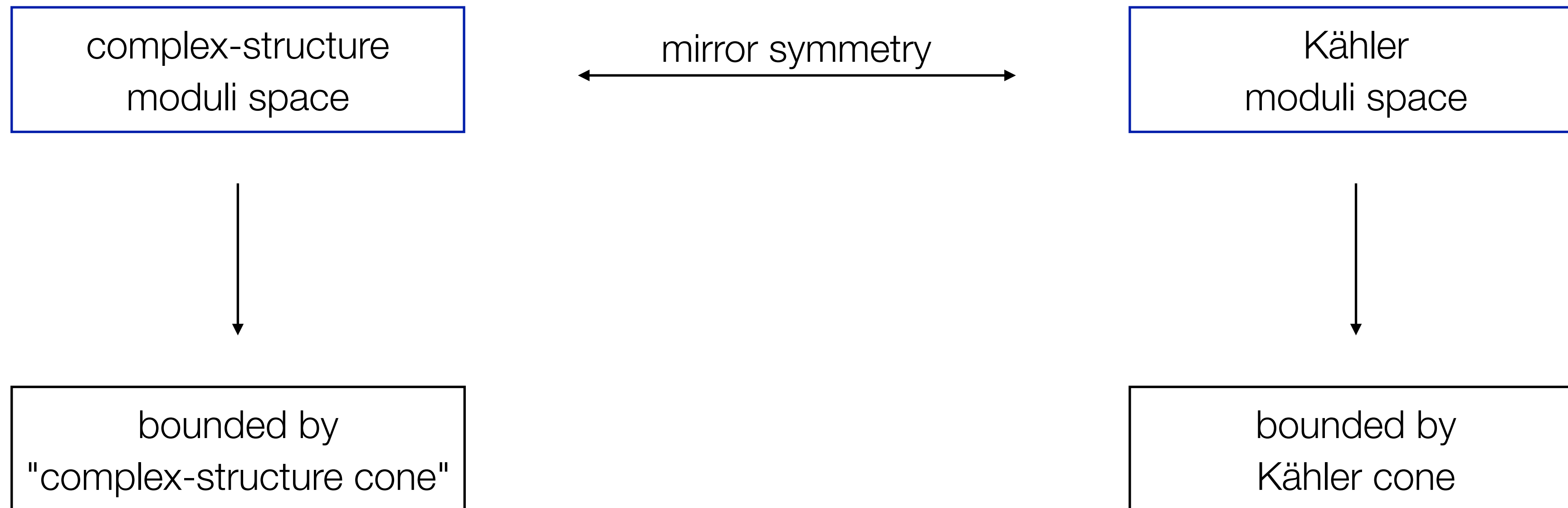
$$N_{\text{flux}} = s \left[\frac{\kappa}{6} (h^0)^2 + \frac{2}{3} \kappa (h^i - u^i h^0) G_{ij} (h^j - u^j h^0) \right] \\ + \frac{1}{s} \left[\frac{\kappa}{6} (f^0)^2 + \frac{2}{3} \kappa (f^i - u^i f^0) G_{ij} (f^j - u^j f^0) \right],$$

$$G_{i\bar{j}} = \text{Kähler metric}, \\ \kappa = \kappa_{ijk} v^i v^j v^K.$$

Using statistical results by Demirtas et al. for κ_{ijk} , the **scaling** of the flux number can be estimated as

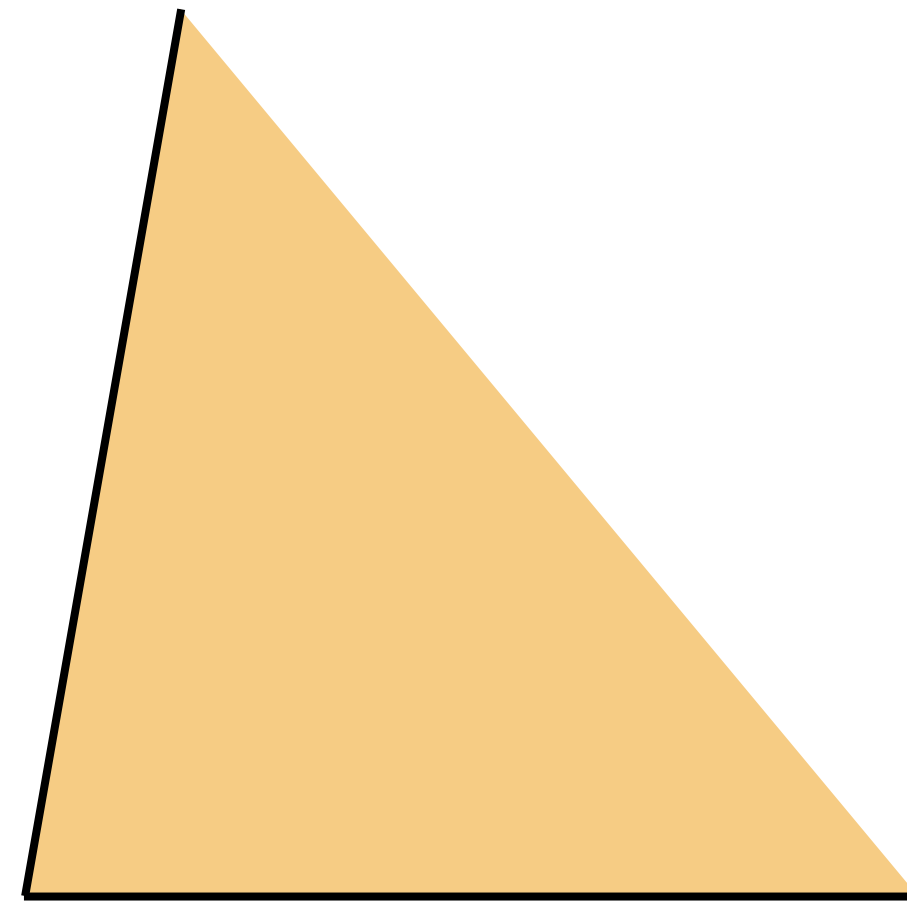
$$N_{\text{flux}} \gtrsim \begin{cases} (h^{2,1})^{+1/2} \|v\| & \text{for } h^0 = f^0 = 0, \\ (h^{2,1})^{-1/2} \|v\|^3 & \text{else.} \end{cases}$$

At large complex-structure, the complex-structure moduli space is **mirror-dual** to the Kähler moduli space



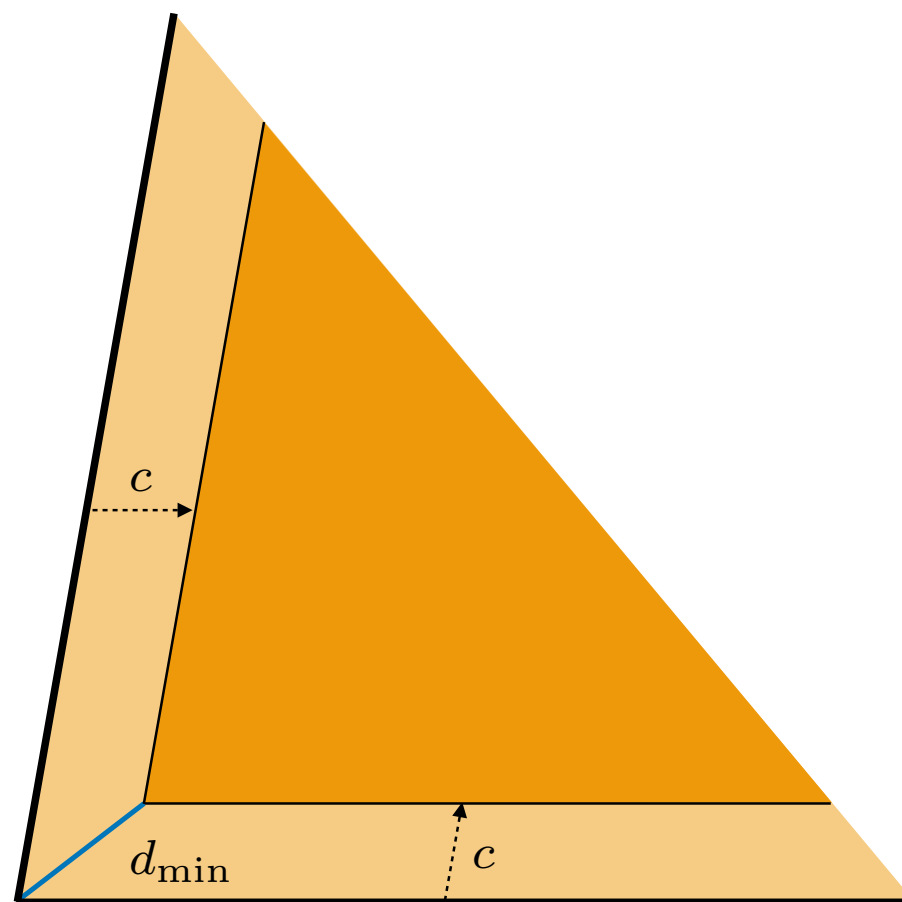
Statistical data on the Kähler cone has been determined by Demirtas et al. for the Kreuzer-Skarke list.

The **Kähler cone** contains all Kähler forms such that curves, divisors and the CY_3 itself have positive volume.

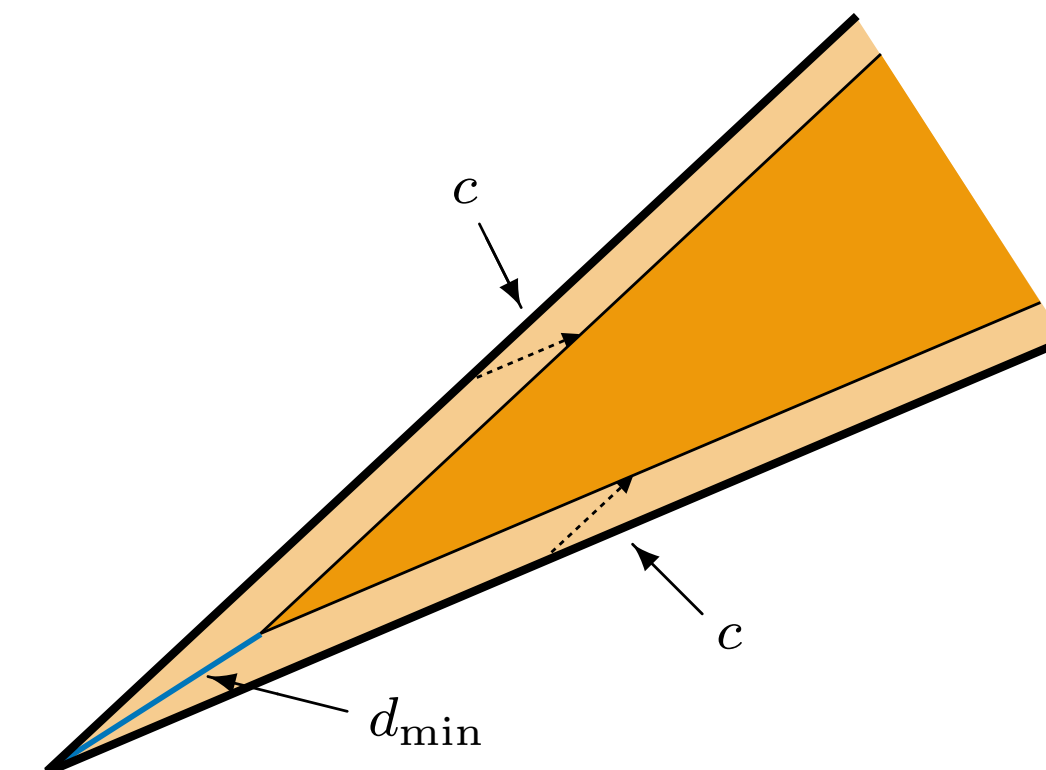


The **Kähler cone** contains all Kähler forms such that curves, divisors and the CY_3 itself have positive volume.

The **stretched Kähler cone** contains all Kähler forms such that all volumes are larger than a constant c .



wide Kähler cone

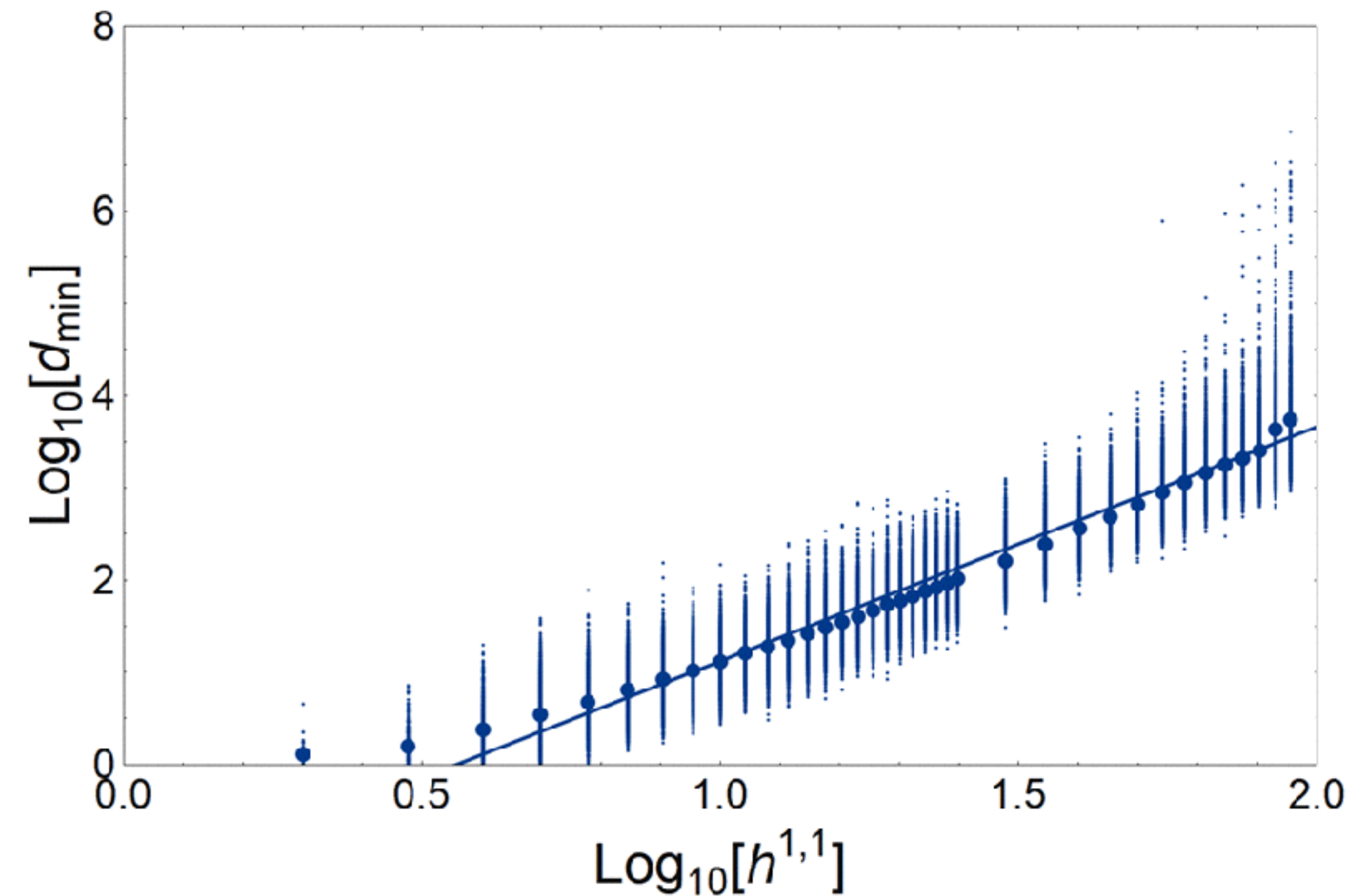


narrow Kähler cone

For the Kreuzer-Skarke list, one finds that the Kähler cone becomes narrower for larger dimensions $h^{1,1}$.

The **narrowness** can be encoded in the distance between the ordinary and stretched Kähler cones

$$d_{\min}[c] \simeq 10^{-1.4} (h^{1,1})^{2.5} c.$$



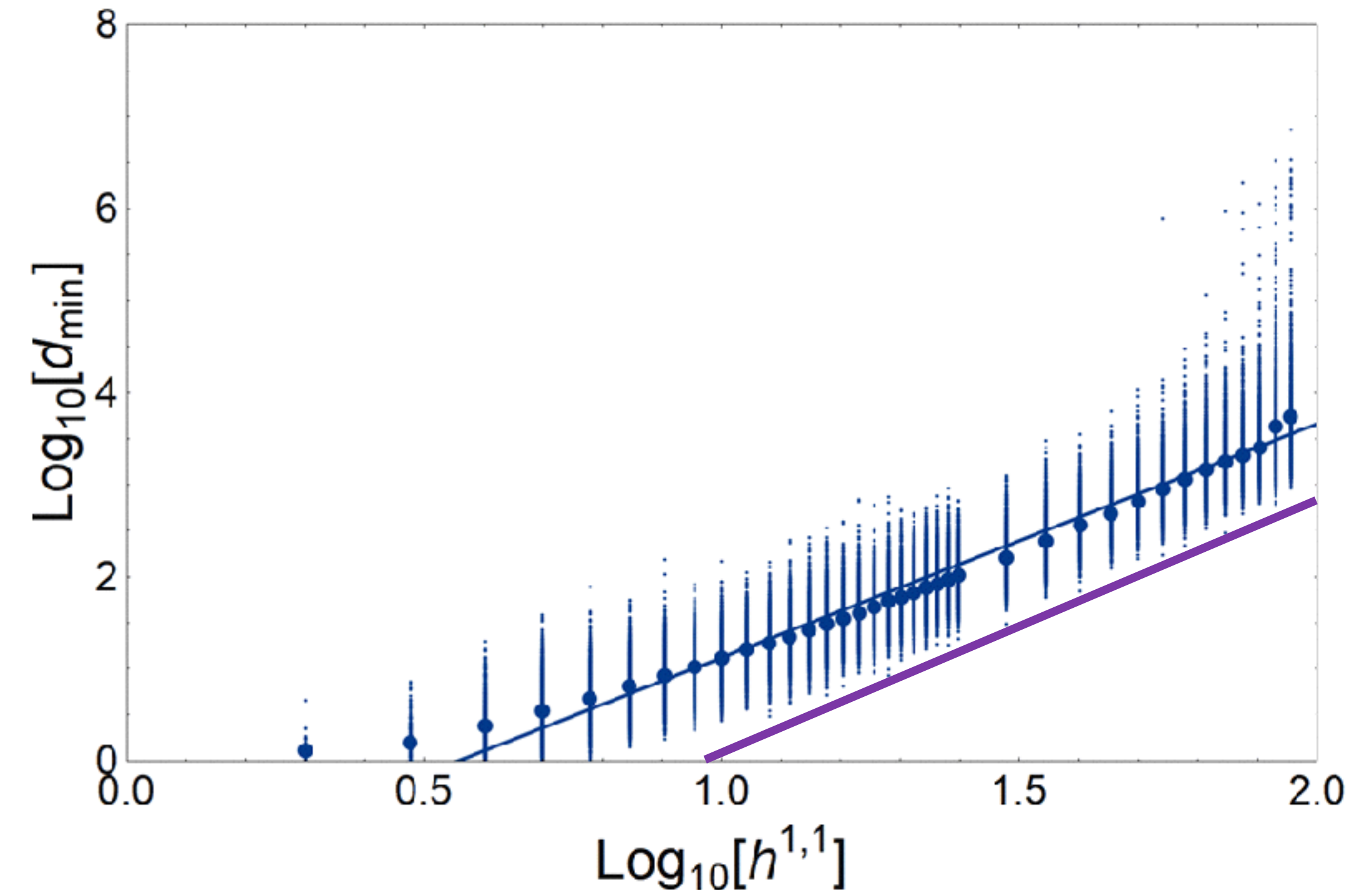
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$$d_{\min}[c] \simeq 10^{-1.4} (h^{1,1})^{2.5} c.$$

A **lower bound** can be estimated from the data as

$$d_{\min}[c] \gtrsim 10^{-2.6} (h^{1,1})^{2.7} c.$$



The result from above can now be combined in the following way (here $\gamma = 1, 3$) ::

$$N_{\text{flux}} \gtrsim (h^{2,1})^{1-\frac{\gamma}{2}} \|v\|^\gamma$$

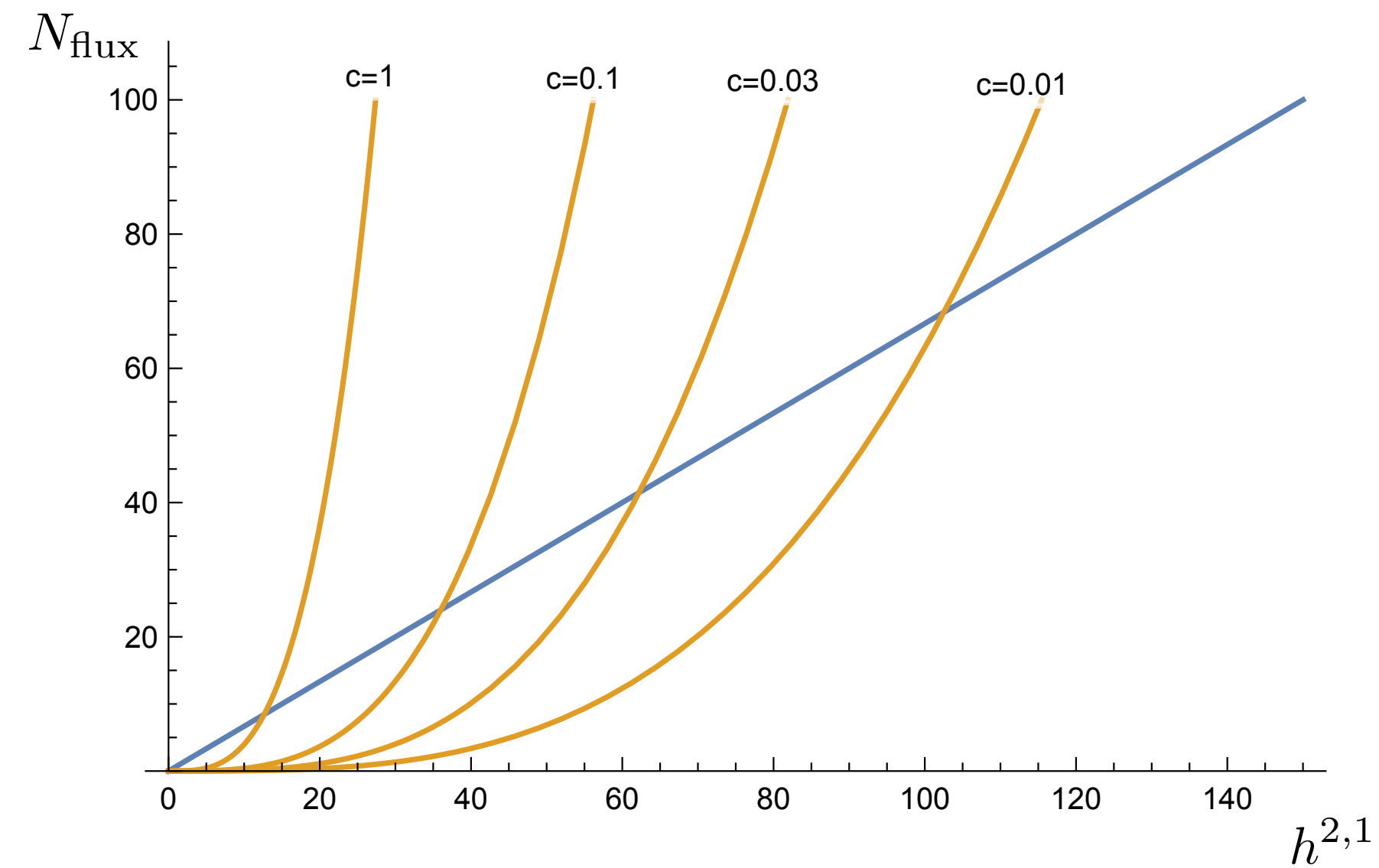
$$\|v\| \geq d_{\text{min}} \gtrsim 10^{-2.6} (h^{2,1})^{2.7} c$$


$$N_{\text{flux}}[c] \gtrsim (h^{2,1})^{1-\frac{\gamma}{2}} [10^{-2.6} (h^{2,1})^{2.7} c]^\gamma$$

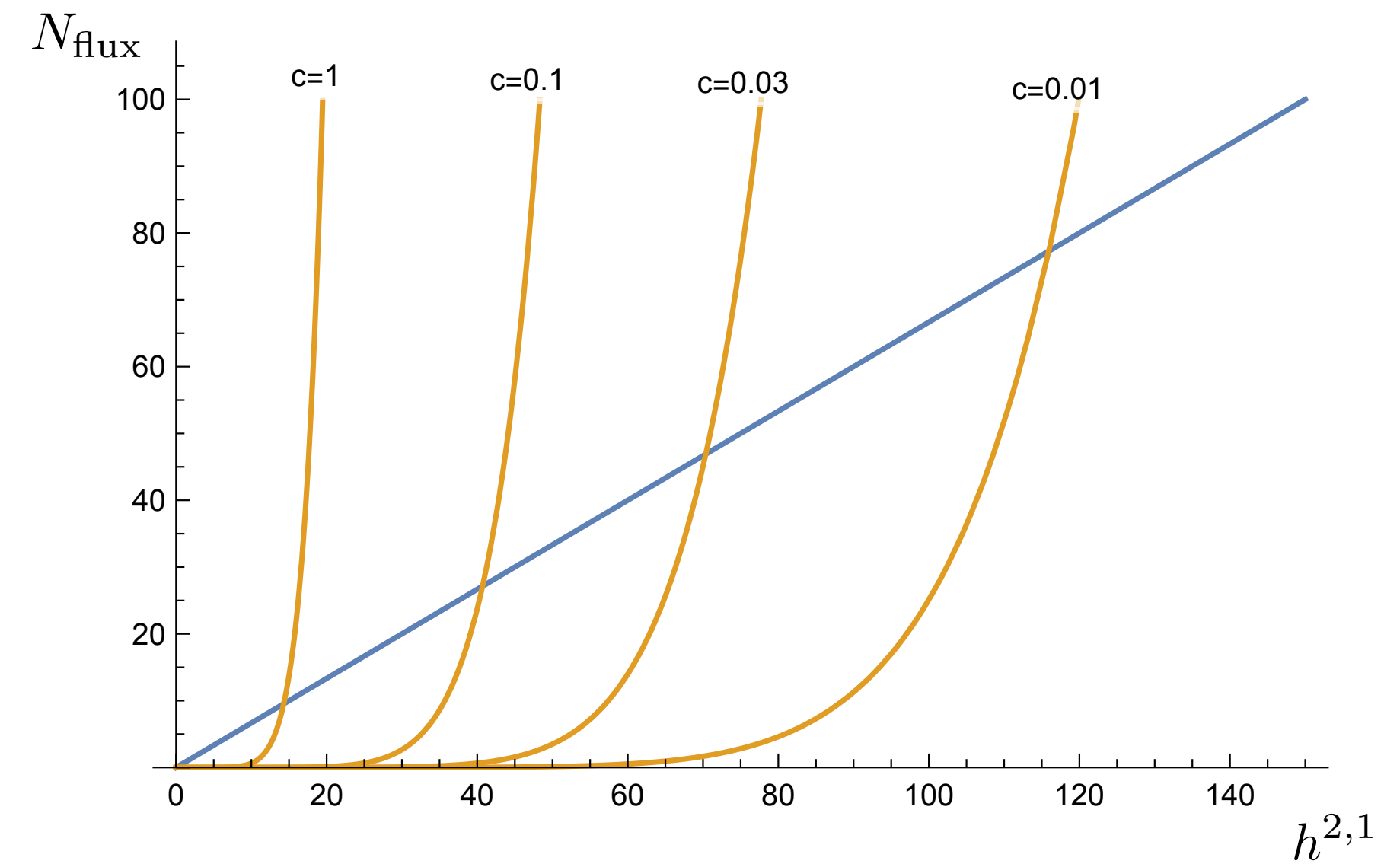
The **constant c** parametrizes minimal volumes of cycles on the mirror-dual side. For a well-controlled large complex-structure limit, c should not be much smaller than one.

large complex-structure — comparison

Even for small values of c , the **flux number** in the large complex-structure regime (orange curves) quickly exceeds the linear bound of the tadpole conjecture (blue curves).



$$\gamma = 1$$



$$\gamma = 3$$

Summary ::

- For generic flux choices and in the large complex-structure limit, the **tadpole conjecture is satisfied**.
- A smaller flux number, violating the conjecture, may be found outside of this regime or for **non-generic flux choices**.

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Summary ::

- The **tadpole conjecture** states that for large $h^{2,1}$, not all moduli can be stabilized.
- When approaching a **boundary** in moduli space, the flux number diverges.
- In the **large complex-structure** regime, the tadpole conjecture is satisfied for generic flux choices.

Implications ::

- If true, the tadpole conjecture implies that the **landscape** of theories with no or few massless scalar fields is smaller than expected.
- A common assumption in the **KKLT** and **Large Volume** scenarios may not be satisfied.